

## A Thomas-Fermi Model

L. H. Thomas [58] and E. Fermi [59] obtained a simplified model of approximating the many-electron structure of atom as the free-electron gas around the nuclei. In free electron gas model, two electrons of opposed spin allow to be into a volume of wave number space equal to  $h^3$ , and in an atom, at a distance  $r$  from the nucleus, all momentum space up to  $p_F(r)$  is occupied. Then the number of electrons per unit volume,  $n_e(r)$ , is given by

$$n_e(r) = \frac{2}{h^3} \cdot \frac{4}{3}\pi p_F(r)^3 = \frac{8\pi}{3h^3} p_F(r)^3. \quad (38)$$

If  $V(r)$  is the electrostatic potential of the atom, the maximum kinetic energy at any distance  $r$  is  $-eV(r)$ , so that the maximum occupied energy level  $E_0$  is

$$E_0 = \frac{p_F(r)^2}{2m_e} - eV(r) \quad (39)$$

The assumption that all negative states are bound would make  $E_0$  zero. Using this equation,  $p_F(r)$  of Eqn. (38) is eliminated, then,

$$n_e(r) = \frac{8\pi}{3h^3} (2m_e e)^{3/2} V(r)^{3/2}. \quad (40)$$

The  $V(r)$  is satisfied with Poisson's equation since the potential is produced both by the nuclear charge and by the electron distribution:

$$\nabla^2 V(r) = -4\pi e n_e(r). \quad (41)$$

Eliminating  $n_e(r)$  between Eqns.(40) and (41), we have

$$\nabla^2 V(r) = -\mu V(r)^{3/2} \quad (42)$$

where

$$\mu = \frac{32\pi^2 e}{3h^3} (2m_e e)^{3/2}. \quad (43)$$

We formulate the potential function  $V(r)$  to

$$V(r) = -\frac{Ze}{r} \chi(x). \quad (44)$$

This means that we are screening the simple Coulomb field by the function  $\chi(x)$ . Suppose also that

$$x = \frac{r}{a_{TF}} \quad (45)$$

where

$$\begin{aligned} a_{TF} &= \left( \frac{3}{32\pi^2} \right)^{2/3} \frac{h^2}{2me^2 Z^{1/3}} \\ &\simeq \frac{0.88534a_0}{Z^{1/3}} \end{aligned} \quad (46)$$

and  $a_0$  is the Bohr radius:

$$a_0 = \frac{h^2}{4\pi^2 me^2} \quad (47)$$

From Eqns. (refequ:poi2) and (44), the screening function  $\chi(x)$  is obtained as a solution of a differential equation, Thomas-Fermi equation,

$$\frac{d^2\chi(x)}{dx^2} - \frac{\chi(x)^{3/2}}{\sqrt{x}} = 0. \quad (48)$$

The assumptions that the potential behaves like a simple Coulomb interaction in the extreme case as  $r$  approaches zero as  $r$  approaches infinity, yield the boundary conditions of the Thomas-Fermi function,

$$\chi(0) = 1 \quad (49)$$

and

$$\chi(\infty) = \chi'(\infty) = 0. \quad (50)$$