High energy-resolved XANES of layered oxides for sodium-ion battery

Hideharu Niwa¹⁻³*, Kazuyuki Higashiyama^{1,2}, Kaoru Amaha¹, Wataru Kobayashi¹⁻³, Kenji Ishii⁴, Yutaka Moritomo¹⁻³*

¹Graduate School of Pure and Applied Sciences, University of Tsukuba, Tsukuba 305-8571, Japan

²Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba 305-8571, Japan
 ³Tsukuba Research Center for Energy Materials Science (TREMS), University of Tsukuba, Tsukuba 305-8571, Japan

⁴Synchrotron Radiation Research Center, National Institutes for Quantum and Radiological Science and Technology, Hyogo 679-5148, Japan

E-mail: niwa.hideharu.ga@u.tsukuba.ac.jp, moritomo.yutaka.gf@u.tsukuba.ac.jp

Abstract

The 3*d* transition metal (TM) oxides with layered structure, Na_xMO_2 (M = Mn, Co), are promising cathode materials of sodium ion secondary batteries (SIBs). Here, we investigated the electronic structure of *M* of four layered oxides ($Na_{0.91}CoO_2$, $Na_{0.66}CoO_2$, $Na_{1.00}MnO_2$, and $Na_{0.54}MnO_2$) by means of high energy resolution fluorescence detected X-ray absorption near edge structure (HERFD-XANES), which utilizes the 1*s* core-hole lifetime-broadening-reduction. The high energy-resolved spectroscopy revealed a shoulder structure in the pre-edge regions of the Co K-edge spectra in $Na_{0.91}CoO_2$. The structure is ascribed to the transition to the Co 3d/4p state via slight hybridization with the Na 3*s* state.

X-ray absorption fine structure (XAFS) is widely used to evaluate the local structure and electronic states of material sciences.¹⁻³⁾ In particular, X-ray absorption near edge structure (XANES) regions provides chemical information of materials such as oxidation states of metal compounds. Although XANES recorded with transmission mode or total fluorescence mode is very convenience for material science, the energy resolution is not necessarily sufficient to discuss the *d*-electron configuration. Recently, high energy resolution fluorescence detected XANES (HERFD-XANES) is receiving much attentions because it revealed detailed electronic structure⁴⁻⁶⁾ that cannot be detected by conventional XANES. For example, the energy resolution of the Pt L₃-edge XANES spectrum improved from 5.2 eV bv HERFD-XANES to 2.4 eV due to а core-hole lifetime-broadening-reduction. The improved energy resolution clarified fine structures in the Pt L₃-edge spectrum of the Pt nanoparticle catalysts in the polymer electrolyte fuel cells, which suggests that surface hydrated species increase the overpotential, and hence, decrease the energy conversion efficiency.⁴⁾

Figure 1(a) shows schematic diagram of X-ray absorption and fluorescence processes at the transition metal (TM) K-edge. The X-ray absorption and fluorescence processes contribute to the transmission XANES and HERFD-XANES, respectively. In the case of the transmission XANES, the lifetime of the absorption process from the 1*s*-core to the 4*s* (or 3*d*) levels is very short, and hence, causes the low energy resolution of the spectroscopy. For example, the spectral width (1*s*-core-hole lifetime width $\Gamma_{\rm K}$) for Mn and Co is 1.16 and 1.33 eV, respectively.⁷ On the other hand, the effective lifetime width of the HERFD-XANES (Γ) is given by

$$\Gamma = 1/\sqrt{(1/{\Gamma_{K}}^{2}) + (1/{\Gamma_{M}}^{2})},$$
(1)

where $\Gamma_{\rm M}$ is 3*p*-core-hole lifetime width.⁸⁾ The lifetime of the fluorescence process from the 3*p*-core to 1*s*-core levels is rather long, and hence, Γ is mainly determined by $\Gamma_{\rm M}$. In other words, the HERFD-XANES used $\Gamma_{\rm M}$ instead of $\Gamma_{\rm K}$ to improve the energy resolution of the spectra. In the actual measurement of the HERFD-XANES, the detection energy of the scattered X-ray is fixed at the $K_{\beta_{1,3}}$ fluorescence line while the energy of the incident X-ray is scanned over the absorption edge [Fig. 1(b)].

Sodium ion secondary batteries (SIBs) have attracted much attention as promising candidates for next-generation batteries beyond lithium ion secondary batteries.⁹⁻¹²⁾ Among

several types of cathode materials for SIBs, layered transition metal oxides $Na_xMO_2(M = Co, Mn)$ are the most typical and promising cathode materials.^{13,14)} They consist of the MO_2 sheet of edge-sharing MO_6 octahedra. In the charge process, Na^+ is extracted between the neighboring MO_2 sheets and an electron is removed from the MO_2 sheet. In the discharge process, Na^+ is inserted between the neighboring MO_2 sheets and an electron is removed from the MO_2 sheet. In the discharge process, Na^+ is inserted between the neighboring MO_2 sheets and an electron is added to the MO_2 sheet. In order to develop high performance materials for SIB, deeper comprehension of the electronic state of TM is indispensable.

In this paper, we investigated the electronic structure of TM of four layered oxides $(Na_{0.91}CoO_2, Na_{0.66}CoO_2, Na_{1.00}MnO_2, and Na_{0.54}MnO_2)$ by means of the HERFD-XANES with high energy resolution. As far as we know, there exists no report on application of the HERFD-XANES to TM in the battery materials. The high energy-resolved spectroscopy revealed a shoulder structure in the pre-edge regions of the Co K-edge spectra in Na_{0.91}CoO_2. The structure is ascribed to the transition to Co 3d/4p state via slight hybridization with the Na 3s state.

Layered oxides were prepared by solid state reaction. For Na_{0.91}CoO₂, Na₂O₂ and Co₃O₄ were mixed in a 1.25: 1 atomic ratio and calcined at 823 K for 16 h in O₂. Then, the product was finely ground, and again calcined in the same condition. For Na_{0.66}CoO₂, Na₂CO₃ and Co₃O₄ were mixed in a 0.7:1 atomic ratio and calcined at 1073 K for 12 h in air. For Na_{1.00}MnO₂, Na₂CO₃ and Mn₂O₃ were mixed in a 1: 1 atomic ratio and calcined at 943 K for 24 h in Ar. For Na_{0.54}MnO₂, Na₂CO₃ and MnCO₃ were mixed in a 0.7:1 atomic ratio and calcined at 1273 K for 12 h in air. The Na concentrations were determined by the Rietveld analyses of the synchrotron-radiation X-ray powder diffraction (XRD) patterns, as described below.

Synchrotron-radiation XRD measurements were performed at the BL-8A beamline of the Photon Factory, KEK. The samples were finely ground and placed in ϕ 0.3 mm glass capillaries. The capillaries were sealed and mounted on the Debye-Scherrer camera. The powder diffraction patterns were detected with an imaging plate. The exposure time was 5 minutes. The wavelength (= 0.68903 Å) of the X-ray was calibrated by the lattice constant of standard CeO₂ powders. The XRD patterns of O3-Na_{0.91}CoO₂ was analyzed by the Rietveld method (Rietan-FP)¹⁵⁾ with a trigonal model ($R\bar{3}m$; Z = 3, hexagonal setting). The XRD pattern of the Mn-based O'3-Na_{1.00}MnO₂ was analyzed by the Rietveld method

with a monoclinic model (C2/m; Z = 2). The XRD patterns of the P2-Na_{0.66}CoO₂ and P2-Na_{0.54}MnO₂ were analyzed by the Rietveld method with a hexagonal model ($P6_3/mmc$; Z = 2). X-ray diffraction patterns and Rietveld refinement profiles are shown in Figs. S1 - S4. In all compounds, no traces of impurities nor secondary phases were observed. The obtained structural parameters are listed in Table S1 – S4.

The Co and Mn K-edge XANES spectra were measured in both the transmission and HERFD modes. The measurements of the transmission XANES were conducted at BL-9C beamline of the Photon Factory, KEK. The powder was finely ground, mixed with BN, and pressed into pellets with 5 mm in diameter. The transmission XANES were recorded with a Si (111) double-crystal monochromator at 300 K. The energy resolution ($\Delta E/E$) was $\sim 2 \times 10^{-4}$ and the photon flux at sample position was $\sim 1 \times 10^{11}$ phs/s. The measurements of the HERFD-XANES were carried out at the contract undulator beamline BL11XU of SPring-8.^{16,17}) The incident X-ray beam was monochromatized with a Si (111) double-crystal monochromator followed by a 2-bounce Si (400) channel-cut monochromator to further increase the incident energy resolution (263 meV at 7670 eV and $\Delta E = 288$ meV at 6460 eV). The pellet sample was placed at normal incident angle (ϕ = 85°) and grazing exit angle (θ = 5°) to the surface in order to suppress the self-absorption.¹⁸⁾ The estimated contributions of self-absorption for the pre-edge regions at Co and Mn K-edges are less than 0.5% and therefore it is negligible. The estimated contributions of self-absorption for the main absorption peak at Co and Mn K-edges are less than 10%. The details of the estimation is described in Supplementary Data. The beam size at sample position was about 0.05 mm (horizontal) \times 1 mm (vertical) and the photon flux was ~ 8×10^{12} phs/s. The emitted X-rays in the horizontal plane were analyzed by a Rowland-mount type emission spectrometer. To select the Co $K_{\beta_{1,3}}$ fluorescence line (~7650 eV), a Ge (444) spherically bent analyzer was used. The total energy resolution was ~1.2 eV at 7650 eV. To select the Mn $K_{\beta_{1,3}}$ fluorescence line (~6490 eV), a Si (333) spherically bent analyzer was used. The total energy resolution was ~1.2 eV at 6490 eV. In the analyses, the background subtraction and normalization were performed using the ATHENA program.¹⁹⁾

The partial density of states (pDOSs) were calculated for O3-NaCoO₂ and O'3-NaMnO₂ based on the density functional theory with use of the plane-wave self-consistent field (PWscf) package.²⁰⁾ A plane wave basis set with a cutoff energy of 952 eV was chosen,

and the projector-augmented wave (PAW) potentials were used. We adopted the exchange-correlation functional of Perdew-Burke-Ernzerhof type within generalized gradient approximation. Electron spin was taken into account in the level of local spin density approximation. In the calculation of O3-NaCoO₂, we used the lattice constant (*a* and *c*) and atomic coordinates of O3-Na_{0.91}CoO₂ obtained by the Rietvelt refinement (Table S1). In the calculation of O'3-NaMnO₂, we used the lattice constant (*a*, *b* and *c*) and the atomic coordinates of O3-Na_{1.00}MnO₂ obtained by the Rietvelt refinement (Table S3). Brillouin-zone sampling for the pDOS calculations was made using the Monkhorst-Pack method with $36 \times 36 \times 24$ and $24 \times 36 \times 24$ *k* points for O3-NaCoO₂ and O'3-NaMnO₂, respectively.

Figure 2 (a) shows the Co K-edge XANES spectra of Na_{0.91}CoO₂ recorded in the HERFD (red circle) and transmission (solid black line) modes. The main absorption peak observed at 7725 eV is ascribed to the electric dipole transition from Co 1s to Co 4p orbital. Although the peak-top intensity of Co K-edge HERFD-XANES spectrum is slightly higher than that of transmission spectra, the peak width is no significant spectral difference between the two modes, indicating the main peak has intrinsic line width wider than $\Gamma_{\rm K}$ (=1.33 eV). Importantly, we observed significant spectral difference in the pre-edge region [inset in Fig. 2(a)]. The HERFD-XANES spectrum shows a shoulder structure at around 7708 eV while no trace of the structure is discernible in the conventional transmission spectrum. Figure 2 (b) shows the Mn K-edge XANES spectra of Na_{1.00}MnO₂ recorded in the HERFD (red circle) and transmission (solid black line) modes. The main absorption peak observed at 6555 eV is ascribed to the electric dipole transition from Mn 1s to Mn 4p orbital. Although the peak-top intensity of Mn K-edge HERFD-XANES spectrum is slightly higher than that of transmission spectra, the peak width is no significant spectral difference between the two modes, indicating the main peak has intrinsic line width wider than $\Gamma_{\rm K}$ (=1.16 eV). In the pre-edge region (inset in Fig. 2(b)), however, we observed significant spectral difference between the two modes, reflecting higher energy resolution in the HERFD mode. The HERFD-XANES spectrum shows well-resolved two sharp absorptions at 6535 and 6537 eV. This makes a sharp contrast with the much broader structure in the transmission mode. Similarly, we observed extra fine pre-edge structures in the Co K-edge XANES spectrum of Na_{0.66}CoO₂ and Mn K-edge XANES spectrum of Na_{0.54}MnO₂ recorded in HERFD mode (Figure 5S).

In Fig. 3(a), we compare the Co K-edge HERFD-XANES spectra in the pre-edge region between Na_{0.91}CoO₂ and Na_{0.66}CoO₂. With increase in the formal valence of Co from +3.09 (Na_{0.91}CoO₂) to +3.34 (Na_{0.66}CoO₂), the intensity of the band at 7706 eV increases. The behavior is reasonable because Na_{0.66}CoO₂ has higher density of states in the unoccupied 3*d* orbitals. In Fig. 3(b), we compare the Mn K-edge HERFD-XANES spectra in the pre-edge region between Na_{1.00}MnO₂ and Na_{0.54}MnO₂. In both the spectra, well-resolved two absorption bands are observed at 6535eV and 6537eV. With increase in the formal valence of Mn from +3.00 (Na_{1.00}MnO₂) to +3.46 (Na_{0.54}MnO₂), the intensity of the higher-lying band at 6537eV increases. The behavior is reasonable because Na_{0.54}MnO₂ has higher density of states in the unoccupied 3*d* orbitals.

Let us compare the K-edge HERFD-XANES spectra in the pre-edge region with the first-principles calculations. Figure 4(a) shows Co (blue curves), O (red curves), and Na (green curves) pDOSs of O3-NaCoO₂. The system is a non-magnetic insulator with the gap of ~ 1 eV.²⁰⁾ The O3-type structures consist of the CoO₂ sheet of the edge-sharing CoO₆ octahedra, causing crystal field splitting of the Co 3d orbital into the lower-lying t_{2g} and upper-lying e_g ones. The t_{2g} and e_g states correspond to the bands at around -0.8 and 1.6 eV, respectively. The t_{2g} band is fully occupied while the e_g band is unoccupied. We observed a strong O 2p – Co 3d hybridization in both the t_{2g} and e_g bands. The 7706 eV peak observed in the Co K-edge spectrum in the pre-edge region [Fig.3(a)] is assigned to the transition from the Co 1s to the e_g band. We note that rather broad Na 3s state is observed at higher energy region above 4eV. We calculated the integrated pDOdoSs in the Na 3s band from 3.5 to 6.0 eV (see Table 1). Even though the Na 3s band is dominated by the Na 3s and O 2p state, the band contains the Co 4s (l = 0), Co 3d (l = 1) and Co 4p (l = 2) states. This suggests that the Co 3d/4s/4p states is hybridized with Na 3s and O 2p state. We further confirmed that the errors of pDOSs were within 6.5% by ab initio calculation under three different conditions (Table S5). Since the Co 1s – Co 4s transition is forbidden, we ascribed the shoulder structure above 7708 eV to the transition to the Co 3d/4p states via slight hybridization with the Na 3s state. The shoulder structure seems to disappear in Na_{0.66}CoO₂. This is probably due to the energy difference between the up and down spin e_g bands, and resultant broadening of the XANES spectrum. Actually, ab initio calculation of P2-Na_{1/2}CoO₂²¹⁾ shows the energy difference of ~ 1 eV between the up and down spin e_g bands.

Figure 4(b) shows Mn (blue curves), O (red curves), and Na (green curves) pDOSs of O'3-NaMnO₂. The system is an antiferromagnetic insulator with the gap of ~ 1 eV. The total energy of the antiferromagnetic state is lower by 65 ± 2 meV/FU than that of the competing ferromagnetic state. The magnitude and direction of the Mn spin is $3.09 \text{ m}_{\text{B}}$ and perpendicular to the MnO₂ layer, respectively. The electronic structure of NaMnO₂ is similar to that of isostructural LiMnO₂.²²⁾ The O'3-type structures consist of the MnO₂ sheet of the edge-sharing MnO₆ octahedra with Jahn-Teller distortion, causing crystal field splitting of the Mn 3d orbital into the lower-lying t_{2g} and upper-lying e_g ones. The t_{2g} and e_{g} bands correspond to the bands around 1.0 and 2.4 eV, respectively. The fine structure in the t_{2g} bands is probably ascribed to the Jahn-Teller distortion. We observed strong O 2p – Co 3d hybridization in both the t_{2g} and e_g bands. The 6535eV and 6537eV peaks observed in the Mn K-edge spectrum in the pre-edge region [Fig.3(b)] are assigned to the transition from the Mn 1s to the t_{2g} and e_g bands, respectively. Similarly to the case of Na_{0.91}CoO₂, we may expect a shoulder structure due to the transition to the Mn 3d/4p states (see Table 2 and Table S6). Nevertheless, no trace of the structure is observed in Na_{1.00}MnO₂ and Na_{0.54}MnO₂. This is probably because lower-lying tail of the strong main absorption at 6555 eV obscures the structure.

Now, let us discuss the interrelation between the intrinsic linewidth and core-hole lifetime width ($\Gamma_{\rm K}$ and $\Gamma_{\rm M}$). In both the Co and Mn compounds, the peak widths of the main absorption peaks that are dominated by the 1s - 4p transition show no significant difference between the two modes (Fig. 2, Table S7). In addition to the small contribution of self-absorption for the main absorption peak (less than 10%), this observation indicates that the intrinsic linewidths of the main peaks are wider than $\Gamma_{\rm K}$. Probably, the main peak is overlapped by the transition to the continuum states such as 5p, 6p, and 7p to cause the wide linewidth. On the other hand, the spectral profiles due to the 1s - 3d transition in the pre-edge region show significant difference between the two modes. This observation indicates that the intrinsic linewidths of the 1s - 3d transitions are narrower than $\Gamma_{\rm K}$. As explained in the introduction, the lifetime width of the HERFD-XANES (Γ) is mainly determined by Γ_M instead of Γ_K . Considering that 2*p*- core hole lifetime width Γ_L of Mn is ~0.5 eV,⁵⁾ the value of $\Gamma_M(<\Gamma_L)$ is considered to be much smaller than 0.5 eV. For further quantitative argument, we should properly consider the instrumental resolution. In the HERFD-XANES at BL11XU beamline of SPring-8, the instrumental resolution (Γ_{inst}), which mainly originated in linewidth of the incident X-ray and the crystal analyzer, was

~1.2 eV at the Co and Mn $K_{\beta_{1,3}}$ fluorescence line. We note that Γ_{inst} is comparable to Γ_K of Co (= 1.33 eV) and Mn (= 1.16 eV). Then, the actual energy resolution of the HERFD-XANES is mainly governed by Γ_{inst} and is slightly larger than 1.2 eV. A future improvement of the spectrometer as well as the X-ray source would enhance the energy resolution.

In conclusions, we investigated the electronic structure of TM of four layered oxides $(Na_{0.91}CoO_2, Na_{0.66}CoO_2, Na_{1.00}MnO_2, and Na_{0.54}MnO_2)$ by means of HERFD-XANES with high energy resolution. The high energy-resolved spectroscopy revealed a shoulder structure in the pre-edge regions of the Co K-edge spectra in Na_{0.91}CoO_2. The structure is ascribed to the transition to the Co 3d/4p state via slight hybridization with the Na 3s state. Our experiment demonstrated that the HERFD-XANES can clarify the fine structures in the pre-edge regions of the Co and Mn K-edge spectra. We believe that application of HERFD-XANES to the other TM battery materials will contribute to deeper comprehension of the variation of the electronic structure with the redox processes. We note that the HERFD-XANES, which uses hard X-ray with high transmittance in materials, is suitable for *in situ/operando* measurements under electrochemical control.

Acknowledgments

This work was supported by JSPS KAKENHI (Grant Number JP16K20940 and JP17H0113). A part of this work was performed under the Shared Use Program of the National Institutes for Quantum and Radiological Science and Technology (QST) Facilities (Proposal No. 2016B-H03 and 2017A-H04) supported by the QST Advanced Characterization Nanotechnology Platform as a program of "Nano-technology Platform" (project No. A-16-QS-0018 and No. A-17-QS-0004) of the Ministry of Education, Culture, Sports, Science and Technology (MEXT), Japan. The synchrotron radiation experiments of HERFD-XAFS were performed at the BL11XU of SPring-8 with the approval of Japan Synchrotron Radiation Research Institute (Proposal No. 2016B3563 and No. 2017A3584). The transmission mode XAFS measurements were performed at BL-9C under the approval of the Photon Factory Program Advisory Committee (Proposal No. 2016G043 and 2017G002) and the X-ray powder diffraction experiments were performed at BL-8A under the approval of the Photon Factory Program Advisory Committee (Proposal No. 2014G507). Part of the *ab initio* calculation was performed using the supercomputing resources at Cyberscience Center, Tohoku University.

References

1) M. Newville, Rev. Mineral. Geochemistry. 78, 33 (2014).

- 2) Y Iwasawa, K Asakura, M Tada, *XAFS Techniques for Catalysts, Nanomaterials, and Surfaces* (Springer, Switzerland, 2017).
- 3) A.E. Russell, A. Rose, Chem. Rev. 104, 4613 (2004).
- 4) Y.-T. Cui, Y. Harada, H. Niwa, T. Hatanaka, N. Nakamura, M. Ando, T. Yoshida, K. Ishii,
- D. Matsumura, H. Oji, H. Ofuchi, and M. Oshima, Sci. Rep. 7, 1482 (2017).
- 5) P. Glatzel and U. Bergmann, Coord. Chem. Rev. 249, 65 (2005).
- 6) M. Bauer, Phys. Chem. Chem. Phys. 16, 13827 (2014).
- 7) M. O. Krause and J. H. Oliver, Natural widths of atomic K and L levels, J. Phys. Chem. Ref. Data **8**, 329 (1979).
- 8) F.M.F. de Groot, M.H. Krisch, J. Vogel, Phys. Rev. B 66, 195112 (2002).
- 9) N. Yabuuchi, K. Kubota, M. Dahbi, S. Komaba, Chem. Rev. 114, 11636 (2014).
- 10) J.W. Choi, D. Aurbach, Nat. Rev. Mater 1, 16013 (2016).
- 11) D. Kundu, E. Talaie, V. Duffort, L.F. Nazar, Angew. Chem. Int. Ed. 54, 3432 (2015).
- 12) J.-Y. Hwang, S.-T. Myung, Y.-K. Sun, Chem. Soc. Rev. 46, 3529 (2017).
- 13) Y. Lei, X. Li, L. Liu, G. Ceder, Chem. Mater. 26, 5288 (2014).
- 14) M.H. Han, E. Gonzalo, G. Singh, T. Rojo, Energy Environ. Sci. 8, 81 (2015).
- 15) F. Izumi and K. Momma, J. Solid State Phemom. 130, 15 (2007).
- 16) K. Ishii, T. Tohyama and J. Mizuki, J. Phys. Soc. Jpn. 82, 021015 (2013).

17) K. Ishii, I. Jarrige, M. Yoshida, K. Ikeuchi, T. Inami, Y. Murakami, and J. Mizuki, J. Electron Spectros. Relat. Phenomena **188**, 127 (2013).

18) P. Pfalzer, J.P. Urbach, M. Klemm, S. Horn, M.L. Den Boer, A.I. Frenkel, and J.P. Kirkland, Phys. Rev. B 60, 9335 (1999).

- 19) B. Ravel and M. Newwille, J. Synchrotron. Rad. 12, 537 (2005).
- 20) P. Giannozzi, S. Baroni, N. Bonini, M. Calandra, R. Car, C. Cavazzoni, et al., J. Phys. Condens. Matter. **21**, 395502 (2009).
- 21) H. Niwa, K. Higashiyama, K. Amaha, W. Kobayashi, and Y. Moritomo, J. Power Sources **384**, 156 (2018).
- 22) D.J. Singh, Phys. Rev. B 55, 309 (1997).

Table 1. Integrated pDOS in the Na 3s band from 3.5 to 6.0 eV for O3-NaCoO₂. *l* is the azimuthal quantum number. Errors are evaluated by *ab initio* calculations under four different conditions. (See Table S5).

element	Na	0	Co (<i>l</i> =0)	Co (<i>l</i> =1)	Co (<i>l</i> =2)
pDOS (state/FU)	0.3421	0.1749	0.0167	0.0005	0.0085
Error (%)	1.5	0.6	6.5	0.1	5.4

Table2. Integrated pDOS in the Na 3s band from 3.0 to 5.0 eV for O'3-NaMnO₂. l is the azimuthal quantum number. Errors are evaluated by *ab initio* calculations under four different conditions. (See Table S6).

element	Na	0	Mn (<i>l</i> =0)	Mn (<i>l</i> =1)	Mn (<i>l</i> =2)
pDOS (state/FU)	0.2448	0.1552	0.0392	0.0003	0.0340
Error					
(%)	3.4	0.9	7.1	1.7	2.8

Figure captions

Figure 1.

(a) X-ray absorption and fluorescence processes. The former and latter processes were used in the transmission XANES and HERFD-XANES, respectively. $\Gamma_{\rm K}$ and $\Gamma_{\rm M}$ are 1*s*-core hole and 3*p*-core hole lifetime widths, respectively. K_{β_1} is the transition from M₃- to K-shell. K_{β_3} is the transition from M₂- to K-shell. (b) Experimental setup for HERFD-XAFS. The detection photon energy of the scattered X-ray is fixed at an energy of K_{$\beta_{1,3}$} fluorescence line, while the energy of incident X-ray is scanned over the absorption edge.

Figure 2

(a) Co K-edge XANES spectra of $Na_{0.91}CoO_2$ recorded in the HERFD (red circle) and transmission (black solid line) modes. The two spectra were normalized at 7749 eV. (b) Mn K-edge XANES spectra of $Na_{1.00}MnO_2$ recorded in the HERFD (red circle) and transmission (black solid line) modes. The two spectra were normalized at 6579 eV. The insets show magnified spectra around the pre-edge regions.

Figure 3

(a) Pre-edge region of Co K-edge XANES spectra for $Na_{0.91}CoO_2$ and $Na_{0.66}CoO_2$ recorded in the HERFD mode. (b) Pre-edge region of Mn K-edge XANES spectra for $Na_{1.00}MnO_2$ and $Na_{0.54}MnO_2$ recorded in the HERFD mode.

Figure 4.

Calculated partial density of states (pDOSs) of (a) O3-NaCoO₂ and (b) O'3-NaMnO₂. Insets show schematic structures of O3-NaCoO₂ and O'3-NaMnO₂. E_f is the Fermi level.

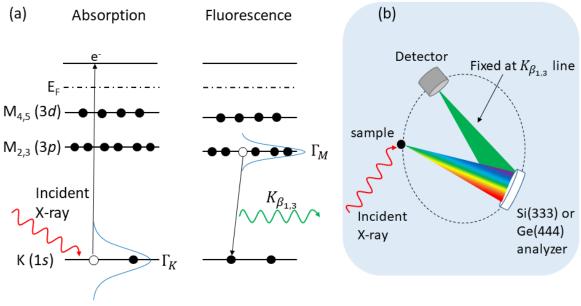


Figure 1

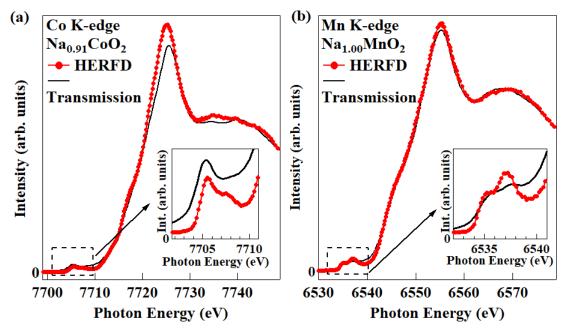


Figure 2

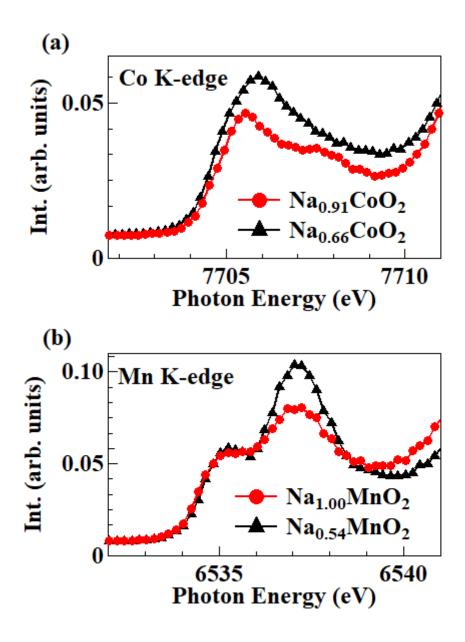


Figure 3

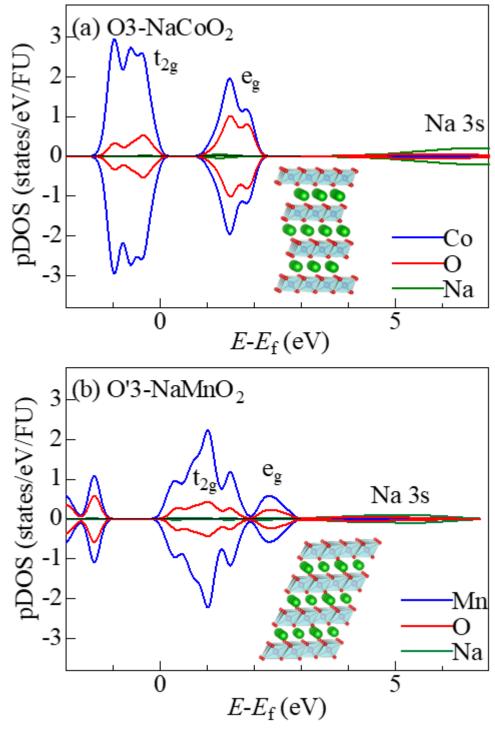


Figure4