

Yield Stress of Mixed Suspension of Silica Particles and Lysozymes:  
The Effect of Zeta Potential and Adsorbed Amount

Atsushi Yamaguchi<sup>a</sup>, Motoyoshi Kobayashi<sup>b,\*</sup>, and Yasuhisa Adachi<sup>b</sup>

<sup>a</sup>Graduate School of Life & Environmental Sciences, University of Tsukuba  
1-1-1 Tennoudai, Tsukuba, Ibaraki 305-8572, Japan

<sup>b</sup>Faculty of Life & Environmental Sciences, University of Tsukuba  
1-1-1 Tennoudai, Tsukuba, Ibaraki 305-8572, Japan

\*Corresponding author

E-mail: kobayashi.moto.fp@u.tsukuba.ac.jp

Tel & Fax: +81-(0)29-853-5721

**Conflict of interest**

The authors declare that they have no conflict of interest associated with this article.

1 **Abstract**

2 To investigate the interactions between colloidal particles in the presence of oppositely charged  
3 proteins, the yield stress of a mixed suspension of silica particles and lysozymes was measured as a function  
4 of lysozyme dose and pH. Further, the corresponding surface properties of silica particles covered with  
5 lysozymes were determined by measuring the adsorbed lysozyme amount and zeta potential. The present  
6 results indicated that an increase in the adsorbed lysozyme amount increases the zeta potential of silica  
7 particles from negative to positive through an isoelectric point. As expected from the Derjaguin–Landau–  
8 Verwey–Overbeek (DLVO) theory, the maximum value of yield stress is obtained around the isoelectric  
9 point and the yield stresses decrease with an increase in the zeta potential magnitude. However, the  
10 maximum yield stress depends on the pH, and yield stresses at the same zeta potential magnitude are  
11 different depending on the zeta potential sign even if the pH values are similar. That is, the relationship  
12 between yield stress and zeta potential is asymmetric with respect to the isoelectric point. These asymmetric  
13 results of the yield stress indicate the existence of non-DLVO forces, such as patch-bridging attraction and  
14 lateral repulsion between adsorbed lysozymes, affected by the adsorbed amount of lysozyme. Our analysis  
15 suggested that the relative adsorbed amount, defined as the ratio of the adsorbed amount to the maximum  
16 adsorbed amount, can be used as a parameter to explain the asymmetric feature of the yield stress with  
17 respect to the zeta potential. In addition to the zeta potential, the surface occupation of adsorbed substances  
18 is critical in controlling particle–particle interactions.

19

20 *Key words:* Non-DLVO force; Surface coverage; Rheology; Proteins; Patch-bridging

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22

## 23 **1. Introduction**

24 Interactions between particles in a colloidal suspension govern the stability and rheological  
25 properties of the suspension. In a mixed suspension of colloidal particles and proteins, adsorbed proteins  
26 change particle–particle interactions by modifying the particle surface through adsorption. Hence, they are  
27 crucial in several industries, including pharmaceutical, cosmetic, food, and environmental industries, for  
28 understanding the mechanism of particle–particle interactions in the presence of proteins [1–3].

29 Particle–particle interactions are often discussed based on the Derjaguin–Landau–Verwey–  
30 Overbeek (DLVO) theory [4,5]. The DLVO theory suggests that the total interactions can be obtained by  
31 summing the electric double layer and the van der Waals forces. The validity of the DLVO theory is  
32 confirmed by direct force measurements [6,7] as well as indirect techniques, such as measuring the  
33 aggregation rate and yield stress [8–11]. Moreover, additional non-DLVO interactions, which are not  
34 included in the DLVO theory, are sometimes observed, especially in colloidal suspensions comprising more  
35 than two materials [12–21].

36 When one considers the interactions between colloidal particles in the presence of proteins, one  
37 can refer to studies on the mixtures of colloidal particles and other materials, such as polymers,  
38 polyelectrolytes, and surfactants. These materials are typically adsorbed on particles and cause non-DLVO  
39 interactions [12–21]. Well-known mechanisms of non-DLVO interactions include patch, bridging, steric,  
40 and hydrophobic interactions. Smellie and La Mer [22] introduced a simple consideration based on the  
41 bridging. They proposed that occupied and unoccupied surface fractions govern the probability of bridging  
42 formation. Thus, the probability of bridging is proportional to  $\theta(1 - \theta)$ , where  $\theta$  refers to the surface  
43 coverage. Zhao et al. [23] reported that the attraction between polystyrene microspheres coated with  
44 poly(N-isopropylacrylamide) microgels is the maximum when the ratio of the adsorbed amount to the  
45 maximum adsorbed amount under the same solution conditions was nearly 0.5. Although they did not  
46 examine the electric properties of the materials as they had focused on hydrophobic interactions, their  
47 results clearly indicated the dependence of particle–particle interactions on the adsorbed amount. According  
48 to the Smellie and La Mer concept, it is implicitly assumed that a surface can be completely covered by  
49 adsorbates. In the case of charged substances, however, the maximum surface coverage is often  
50 significantly less than the highest possible adsorption owing to electric repulsion between the adsorbed  
51 substances [24–27]. Thus, the applicability of Smellie and La Mer’s concept on charged substances may be  
52 limited.

53 The effects of the electric properties of materials on particle–particle interactions are often  
54 studied in the presence of oppositely charged polyelectrolytes [12,17–20]. The adsorption of oppositely  
55 charged polyelectrolytes onto colloidal particles induces charge neutralization and overcharging. The net  
56 charge of a particle is reduced or reversed by the adsorption of opposite charges from polyelectrolytes.  
57 Additionally, patch, bridging, and steric interactions may occur. These interactions should depend on the

58 amount of adsorbed substances. However, systematic and quantitative studies on the effect of the adsorbed  
59 amount of charged substances on particle–particle interactions are scarce.

60 The silica particle–lysozyme system is widely studied [28–32] because of the high degree of  
61 accessibility and usability of both materials for fundamental research and practical use. Lerche and Sobisch  
62 [30] used the sedimentation technique and demonstrated that suspensions are stabilized at a low pH  
63 although the absolute value of the zeta potential of lysozyme-coated silica nanoparticles is low. Meanwhile,  
64 suspensions at a higher pH were destabilized although the absolute values of the zeta potential were  
65 relatively high. Nevertheless, we demonstrated in our previous study [32] that a qualitative explanation for  
66 the aggregation rate of submicron silica particles in the presence of lysozymes can be partly achieved in  
67 terms of the DLVO theory by considering the heterogeneity of surface charges. However, further systematic  
68 and quantitative studies are necessary to understand the interactions between silica particles in the presence  
69 of lysozymes.

70 Yield stress measurement is a useful method to investigate particle–particle interactions [10,33].  
71 The yield stress of a suspension depends on the interparticle forces, solid volume fraction, particle size  
72 distribution, and particle networks [10,11,33–36]. By choosing and consolidating the sample conditions  
73 and preparation method, one can obtain information on interparticle forces directly from the yield stress.

74 We herein present the measured yield stresses of mixed suspensions of lysozymes and silica  
75 particles. Furthermore, the yield stresses are compared with experimentally obtained zeta potentials and  
76 adsorbed lysozyme amounts. Interestingly, the yield stress exhibits a DLVO-like interaction regime and a  
77 non-DLVO-like regime. The mechanisms of both interaction regimes are discussed in terms of the adsorbed  
78 amounts.

79

## 80 **2. Materials and Methods**

### 81 **2.1 Materials**

82 **Silica particles:** In this study, we used commercially available colloidal silica particles (KE-P30,  
83 Lot No. 3A15, Nippon Shokubai). The particles were synthesized using a process similar to that reported  
84 by Stöber [37]. Silica powder was heated at 800 °C for 24 h before preparing suspensions and storing them  
85 in a drying desiccator.

86 The number average diameter  $\langle d \rangle$  and the standard deviation of the silica particles were  
87 measured by transmission electron microscopy (TEM, H7650, Hitachi High-Technologies Corporation). To  
88 analyze the average diameter and standard deviation, areas of 800 projected circles of silica particles were  
89 measured from micrographs using the ImageJ software and converted into diameters. The average diameter  
90 was 302 nm, while the standard deviation was 20 nm.

91 The specific surface areas were obtained from nitrogen gas adsorption (BET) measurement  
92 performed on the SA 3100 instrument (Beckman Coulter) and from TEM micrographs using the  
93 relationship  $(6/\rho) \langle d^2 \rangle / \langle d^3 \rangle$ . Here,  $\rho$  is the density of the silica particles, which was reported to

94 be  $2.2 \times 10^3 \text{ kg/m}^3$  [37]. The specific surface area, as calculated from TEM, was  $9.0 \times 10^3 \text{ m}^2/\text{kg}$ , while  
95 that obtained from BET measurements was  $1.25 \times 10^4 \text{ m}^2/\text{kg}$ .

96 **Lysozymes:** Hen egg-white lysozyme (L6876, Sigma Aldrich) was used without further  
97 purification. It has been previously reported that a lysozyme is an oval sphere with dimensions of  $3 \text{ nm} \times 3$   
98  $\text{nm} \times 4.5 \text{ nm}$  [27,38,39]. In this study, we assumed the lysozymes to be spheres of 1.7 nm radius, as  
99 calculated from the cubic mean radius [31,40]. The molecular weight of the lysozyme was 14.3 kDa. All  
100 the lysozyme solutions were used within 2 weeks of preparation. The electrophoretic mobility of the  
101 lysozyme is reported elsewhere [39,40] and the isoelectric point (iep) of the lysozymes lies in the pH range  
102 of 9–11.

103 **Chemicals and Water:** The KCl (JIS special grade, Wako Pure Chemical Industries) solution  
104 was used as an ionic strength adjuster. KOH and HCl solutions were used to control the pH. Through the  
105 present study, deionized water (Elix Millipore) with an electrical conductivity of approximately  $0.07 \mu\text{S}/\text{cm}$   
106 was used to prepare all solutions and suspensions.

107

## 108 2.2 Yield stress

109 The yield stresses of mixed suspensions of silica particles and lysozymes were measured by the  
110 vane method. In the vane method, a vane-type spindle is inserted into a sample and rotated at an extremely  
111 low rotational speed. The relation of torque vs. time exhibits the maximum torque when the material yields.  
112 As yield stress  $\tau_y$  is the shear stress when a sample yields,  $\tau_y$  can be calculated from the maximum torque  
113  $T_m$  and spindle dimension by [41,42],

$$T_m = \frac{\pi D^3}{2} \left( \frac{H}{D} + \frac{1}{3} \right) \tau_y, \quad (1)$$

114 where  $D$  and  $H$  are the spindle diameter and height, respectively.

115 The concentration of silica was 50 wt.% and the mass ratio of lysozyme to silica ranged from 0–  
116  $0.012 \text{ g/g}$ . The KCl concentration was maintained at 10 mM. First, the silica particles were placed in the  
117 sample bottle after satisfying the requirements for the vane method, as detailed by Dzuy and Boger [41].  
118 Deionized water was added to the sample, subsequently it was shaken vigorously and stored overnight to  
119 wet the silica particles. Next, the KCl solution, lysozyme solution, and HCl or KOH were added to the silica  
120 suspension after it was sonicated. The total ionic strength of the silica suspension was adjusted to 10 mM  
121 with KCl, HCl, and KOH solutions. The sample pH was measured with a combination electrode  
122 (6.0234.100, Metrohm). Small amounts of HCl and/or KOH solutions were added to the sample to adjust  
123 the pH within an error margin of  $\pm 0.1$  with respect to the target pH. The sample was subsequently shaken  
124 gently for 24 h.

125 Yield stress measurements were performed with viscometers (LV DV-II+Pro, HADV-II+Pro,  
126 BROOKFIELD) with the V-74 vane spindle. To minimize the effect of the sample preparation process, the

127 following operations were performed before each measurement. The sample suspension was shaken  
128 vigorously for 20 s with a vortex mixer (MFG No.6087546, TAITEC) to maintain a uniform floc  
129 morphology [43] and left to stand for 2 min. Subsequently, the spindle was inserted into the sample. After  
130 3 min, the measurement was started. The rotational speed of the spindle was 0.1 rpm, and the torque was  
131 measured after every 1 s. The yield stress measurements were performed at least five times for each sample  
132 at intervals of 10 min. The measurements were performed at 20 °C.

133

### 134 **2.3 Zeta potential**

135 The zeta potential of the silica particles covered with lysozymes was measured by the  
136 electroacoustic method (DT-1202, Dispersion Technology Inc.). Silica suspensions of 50 g/L (0–0.024 g/g  
137 mass ratio of lysozymes to silica particles) were prepared with 10 mM KCl. The pH of the suspensions was  
138 adjusted to 5 or 7 with KOH and/or HCl within an error margin of  $\pm 0.1$ . The suspensions were shaken  
139 gently for more than 24 h to equilibrate the adsorption of lysozymes. Each measurement was started  
140 immediately after gently mixing the suspension to prevent the sedimentation of silica particles. The  
141 measurements were performed thrice on each sample; the zeta potential measurements were performed at  
142 25 °C. It is noteworthy that we can compare the results of the zeta potential measurements with those of  
143 the yield stress measurements using the mass ratio of lysozymes to silica particles because the number of  
144 lysozymes to one silica particle is constant.

145

### 146 **2.4 Adsorbed amount**

147 The adsorbed amounts of lysozymes on silica particles were taken from a previous study [32].  
148 In our previous study, the mixed suspension of silica particles and lysozyme at pH 5 and 7 with 10 mM  
149 KCl was prepared similarly as that for the zeta potential measurements. The silica particles covered with  
150 lysozymes were removed by centrifugation with 20630 g for 5 min at room temperature (20 °C) using a  
151 centrifuge equipment (3520, Kubota). The lysozyme concentrations in the supernatants were determined  
152 by ultraviolet (UV) absorbance measurements at a wavelength of 280 nm using a spectrophotometer (UV-  
153 1650PC, Shimadzu) and the standard calibration curve of absorbance as a function of lysozyme  
154 concentration. The adsorbed lysozyme amount was obtained as the difference between lysozyme dose and  
155 supernatant concentration.

156 The surface coverage was calculated from the adsorbed amount by the relation  $\theta =$   
157  $1000N_A A_{LSZ}\Gamma/(M_w \cdot a_{silica}^{TEM})$ . Here,  $N_A$  is the Avogadro constant;  $A_{LSZ}$  is the projected area of the  
158 lysozyme, i.e.,  $9.3 \times 10^{-18}$  m<sup>2</sup>, when the lysozyme is regarded as a sphere with a radius of 1.7 nm;  $\Gamma$  is  
159 the adsorbed amount of lysozyme in g/g;  $M_w$  is the molecular weight of lysozyme;  $a_{silica}^{TEM}$  is the TEM  
160 specific surface area of the silica particle. It is noteworthy that  $A_{LSZ}$  in this study is the value between  
161 projected areas at side-on and end-on adsorptions [27].

162

163 **3. Results and Discussion**

164 **3.1 Surface properties of silica particles covered with lysozymes**

165 The values of zeta potential of the lysozyme-coated silica particles in a 10 mM KCl solution at  
166 pH 5 and 7 are plotted against the adsorbed amount and the corresponding surface coverage in Figure 1.  
167 The symbols denote the experimental values. The adsorption of lysozyme is almost saturated in this  
168 experimental range. That is, a further increase in lysozyme dose does not change the zeta potential and the  
169 adsorbed amount significantly. Meanwhile, under both pH conditions, the lysozymes are completely  
170 adsorbed on the silica particles until the adsorbed amount reaches almost the maximum value.

171 Figure 1 shows that the bare silica particles possess negative zeta potentials at both pH values.  
172 An increase in the adsorbed lysozyme amount results in a reduced zeta potential magnitude. A further  
173 increase in the adsorbed amount results in charge reversal and the zeta potentials of the lysozyme-coated  
174 silica particles becoming positive. Finally, the zeta potentials become almost constant at adequately high  
175 lysozyme doses and the adsorbed lysozyme amount becomes saturated (see Figure S1 in Supporting  
176 Information). The iep at pH 5 is found at a lower lysozyme dose than that at pH 7 because silica particles  
177 have lower negative charge density, while lysozymes exhibit a strong positive charge at pH 5.

178 The maximum adsorbed amount at pH 7 is 4.2 times larger than that at pH 5. These trends have  
179 been reported elsewhere as well [27–29,31]. We consider that the dependence of the adsorbed amount on  
180 pH can be understood in terms of the three-body random sequential adsorption (RSA) model that considers  
181 the electrostatic forces between lysozymes and those between a lysozyme and silica surface [24,25].  
182 According to the three-body RSA model, at pH 5, the strong positive charge of the lysozymes and the weak  
183 negative charge of the silica particles result in a strong lateral repulsion between the lysozymes and, hence,  
184 decreases the maximum adsorbed amount. Further, at pH 7, the weak positive charges on the lysozymes  
185 and the strong negative charges on the silica particles result in a weak lateral repulsion between the  
186 lysozymes and in a large maximum adsorbed amount.

187

188 **3.2 Results of yield stress measurement**

189 The yield stress is plotted against the adsorbed amount and corresponding surface coverage in  
190 Figure 2. The symbols represent experimental values in 10 mM KCl at pH values of 5 and 7. The lines  
191 correspond to the yield stresses of silica suspensions without lysozyme in 0.5 M KCl, where electrostatic  
192 repulsions should be ineffective [32]. Here, the adsorbed amounts are interpolated from the experimental  
193 values. Under both pH conditions, an increase in the adsorbed amount results in an increase in the yield  
194 stress until it reaches the maximum value. Subsequently, the yield stresses decrease with an increase in the  
195 adsorbed amount. The maximum yield stress at pH 5 is located at a lower adsorbed amount of lysozymes  
196 than that at pH 7. These results agree well with the zeta potential measurements shown in Figure 1.

197 The yield stresses of the suspensions are plotted against the zeta potentials of silica particles  
198 covered with lysozymes, as shown in Figure 3. The horizontal line is the yield stress of the bare silica

199 particle suspension with 0.5 M KCl. The vertical line is the zeta potential, where the relative adsorbed  
200 amount (which is discussed in the fifth section at section 3.3) is equal to 0.5. The pictures above are a  
201 schematic representation of the gap between silica particles. Here, the zeta potentials are interpolated from  
202 the experimental values (see Figure S1). As expected from the DLVO theory and previous studies, a strong  
203 correlation exists between the experimental results of yield stress and zeta potential. In other words, Figure  
204 3 shows that, under both pH conditions, the maximum yield stresses are located around the iep, and  
205 increasing the magnitude of the zeta potential decreases the yield stress. These trends are explained in terms  
206 of the classical DLVO theory. In simple terms, large yield stresses around the iep are caused by the van der  
207 Waals attractive forces and a weak electric repulsion between silica particles. Moreover, increasing the  
208 magnitudes of the zeta potentials increase the electric repulsion between the silica particles, and a large  
209 electric repulsion results in a lower yield stress.

210           Some trends in Figure 3 can be explained in terms of the DLVO theory. Nevertheless, other trends  
211 cannot be explained unless non-DLVO forces are considered. First, the maximum yield stress depends on  
212 the pH. Next, under both pH conditions, the maximum yield stresses of the mixed suspensions of lysozymes  
213 and silica particles are larger than the yield stresses of bare silica suspensions in 0.5 M KCl, where electric  
214 repulsions do not come into play. Furthermore, the graph is not symmetrical with respect to the iep  
215 especially at pH 7. That is, the magnitude of the yield stress depends on the pH and zeta potential sign. The  
216 yield stresses are different even at the same magnitude of zeta potential if the sign is different. This indicates  
217 the importance of the surface coverage effects that are not reflected in the zeta potential.

218

### 219 **3.3 Discussion on the mechanisms of non-DLVO forces**

220           When one considers the mechanisms governing the maximum yield stress at each pH, the specific  
221 feature of the flocculation of silica particles must be considered. Figure 4 shows the yield stress of silica  
222 suspension without lysozymes in 0.5 M KCl. Low yield stresses are observed under conditions of low pH;  
223 this observation is similar to previously reported results [43]. The low yield stress at low pH values is  
224 considered to be due to repulsive forces between the hairy or gel-like structures on silica surfaces [37,44],  
225 or the reduction in additional attractions by ion–ion correlation [43]. This specific feature of silica particles  
226 should be the contributing factor to the pH dependence of the maximum yield stress of mixed suspensions  
227 of silica particles and lysozymes. Furthermore, large errors at pH 6 and 7 are observed. We infer that this is  
228 caused by the instability of the suspension. The specific hindering of aggregation of bare silica particles  
229 occurs at a low pH and the boundary between the aggregation and dispersion regimes is located at pH 6–7  
230 [44]. Therefore, we consider that the suspensions around pH 6–7 are unstable and affected easily by the  
231 small difference in pre-treatment prior to each measurement.

232           Kobayashi et al. [37,44] demonstrated that the aggregation rate of silica particles decelerated at  
233 low pH values or when the particle size was small. Such specific hindering of aggregation can be suppressed  
234 by heating. Furthermore, in our previous study, we used the same particles with heating as those used in

235 this study, and demonstrated that the aggregation rate in the fast aggregation regime was independent of pH  
236 [32]. Additionally, in the present study, a decrease in the yield stress is observed at low pH values although  
237 the silica particles are heated before experimentation. We theorize that this contradiction between the yield  
238 stress and aggregation rate is owing to the difference in characteristic distances that affect each phenomenon.  
239 That is, the characteristic distance for aggregation may be larger than that for yielding because aggregation  
240 is an approaching process, while yielding is the broken retracting process of floc [15]. Thus, the aggregation  
241 rate may be less affected by surface structures.

242 Another important feature of the maximum yield stress is that the maximum yield stresses of  
243 silica–lysozyme suspensions in 10 mM KCl are larger than the yield stresses of bare silica suspensions in  
244 0.5 M KCl, although the KCl concentration is low (see Figure 3). Needless to say, we could not detect the  
245 yield stress of suspensions of bare silica particles in 10 mM KCl at pH 5 and 7. According to the DLVO  
246 theory, electric repulsion between silica particles decreases with an increase in the ionic strength owing to  
247 the compression of the electric double layer. Previously, Franks [43] demonstrated that the yield stress of  
248 silica suspensions either increased or remained with ionic strength increase. Therefore, our observation of  
249 higher maximum yield stresses of mixed suspensions of lysozymes and silica particles compared to the  
250 yields stress of bare silica particle suspension with 0.5 M KCl cannot be explained in terms of the specific  
251 features of bare silica particles reported by Franks or the DLVO theory. Thus, additional attractive forces  
252 must exist, such as bridging and patch attraction. Furthermore, these additional attractive forces induce the  
253 pH dependence of the maximum yield stress as the surface properties of silica particles and lysozymes  
254 depend on pH.

255 According to the DLVO theory, yield stress depends only on the van der Waals force and electric  
256 double layer force, and is characterized by the zeta potential. Thus, if one plots yield stresses against zeta  
257 potentials, the graphs become symmetrical with respect to the iep. However, in Figure 3, the graphs are  
258 asymmetrical with respect to the iep. At pH 5, the difference between the negative and positive zeta potential  
259 is insignificant. At pH 7, the range of the detected yield stresses is wider when the zeta potential is negative  
260 as compared to its counterpart when the zeta potential is positive. These results indicate the existence of  
261 additional forces that depend on the adsorbed amount of lysozymes and pH. We consider that these  
262 additional forces include patch attraction, bridging, and lateral repulsion between adsorbed lysozymes.

263 To examine the effect of the adsorbed amount of lysozymes on particle–particle interactions, we  
264 introduce the relative adsorbed amount,  $\theta_{rel}$ , which is the ratio of the adsorbed amount  $\Gamma$  to the maximum  
265 adsorbed amount  $\Gamma_{max}$  under the same solution conditions. If a linear relation between the adsorbed  
266 amount and surface coverage is assumed,  $\theta_{rel}$  can be obtained by  $\theta_{rel} = \Gamma/\Gamma_{max} = \theta/\theta_{max}$  with the  
267 maximum surface coverage  $\theta_{max}$ . Here,  $\theta_{rel}$  may represent the fraction of occupied spaces, including the  
268 region where electrostatic repulsion occurs, and the probability of occurrence of patch–bridging attraction  
269 should be proportional to  $\theta_{rel}(1 - \theta_{rel})$ . In other words, at  $\theta_{rel} = 0.5$ , the patch–bridging attraction can  
270 be considered maximal, and when  $\theta_{rel}$  is far from 0.5, the attraction may decrease owing to the small

271 probability of patch-bridging attraction and/or overlap of adsorbed lysozymes. The consideration on  $\theta_{rel}$   
272 is based on Smellie and La Mer's concept of the probability of bridging formation. However, in the present  
273 study,  $\theta_{rel}$  is based on  $\Gamma_{max}$  to consider the electric repulsion between lysozymes. In Figure 5, the yield  
274 stresses are plotted against the relative adsorbed amount,  $\theta_{rel}$ . The lines represent the relative adsorbed  
275 amounts corresponding to the iep. As shown in Figure 5, the maximum yield stresses at pH 5 and 7 are  
276 located around the isoelectric point instead of around  $\theta_{rel} = 0.5$ . This implies that, in mixed suspensions  
277 of oppositely charged lysozymes and silica particles, the relative adsorbed amount exhibits a smaller  
278 correlation with the yield stress than the zeta potential. However, the concept of relative adsorbed amount  
279 may explain the asymmetry between the yield stresses and zeta potential by the following consideration.

280 Figure 3 illustrates the relationship between yield stresses and zeta potentials at pH 5 and 7 and  
281 the schematic representations of the gap between silica particles under the corresponding conditions. The  
282 vertical lines represent the zeta potentials when the relative adsorbed amount  $\theta_{rel}$  is 0.5. At pH 5 (Figure  
283 3 (a)), the zeta potential is +8 mV at  $\theta_{rel} = 0.5$ . Thus, attractive forces occur between lysozyme-coated  
284 silica particles regardless of the zeta potential sign. Furthermore, the adsorbed amount of lysozyme is much  
285 smaller than that at pH 7. These adsorption profiles result in a symmetric DLVO-like or small asymmetric  
286 feature of yield stress with respect to the zeta potential at pH 5.

287 At pH 7 (Figure 3 (b)),  $\theta_{rel} = 0.5$  is located near the iep. Thus, attractive forces between  
288 adsorbed lysozymes on a silica and the other silica surfaces exist when the zeta potential is negative.  
289 Meanwhile, when the zeta potential is positive, the adsorbed lysozymes may form a homogeneous layer on  
290 the silica surface. In this region, the adsorption density of lysozymes in the gap between silica particles  
291 reaches the maximum adsorbed amount. Therefore, repulsive forces exist between lysozymes on a particle  
292 and lysozymes on another particle. At pH 7, the additional attraction between silica particles covered with  
293 small amount of lysozymes ( $\theta_{rel} < 0.5$ ) works in the range of negative zeta potential, while the repulsion  
294 between silica particles covered with larger amount of lysozymes ( $\theta_{rel} > 0.5$ ) works at positive zeta  
295 potential regime. Therefore, the yield stresses are detectable in a wider zeta potential range when the zeta  
296 potential is negative.

297 It is known that the magnitude of yield stress is proportional to the square of the zeta potential.  
298 This linearity was observed in theoretical models [45,46], in which interparticle forces are expressed in  
299 terms of the DLVO theory, and has been confirmed experimentally [11,17,33]. At low lysozyme doses and  
300 pH 7, the plots of yield stresses against the square of the zeta potential are curved (see supporting  
301 information Figure S2). This observation supports the existence of additional non-DLVO forces such as  
302 patch-bridging attraction. Meanwhile, the linear relationship between the yield stresses and the square of  
303 the zeta potential at a high dose of lysozyme at pH 7 implies that the interaction between silica particles is  
304 DLVO like because of a homogeneous lysozyme layer on the silica surface.

#### 305 306 **4. Conclusions**

307 To investigate the interactions between colloidal particles in the presence of proteins, the yield  
308 stresses of mixed suspensions of lysozymes and silica particles were measured as a function of lysozyme  
309 dose under two different pH conditions. The measured yield stresses were compared to surface properties  
310 that were characterized by the adsorbed amounts and zeta potentials. Our results indicated that the zeta  
311 potential was an important parameter in estimating the interactions, as expected from the DLVO theory.  
312 However, the zeta potential alone could not explain the several trends followed by the yield stresses, such  
313 as the pH dependence of the maximum yield stress and the different yield stresses at the same magnitude  
314 of zeta potential. These results confirmed the existence of non-DLVO forces, such as the patch-bridging  
315 attraction and lateral repulsion between adsorbed lysozymes. The relative adsorbed amount of lysozymes  
316 could explain the asymmetric feature of yield stresses with respect to the iep. The relationship between  
317 yield stresses and squares of zeta potentials suggested that a transition between DLVO-like and non-DLVO-  
318 like interactions occurred by the increase in adsorbed lysozyme amounts. Furthermore, the relative  
319 adsorbed amount could be used as a parameter to analyze the transition of interaction mechanisms and the  
320 existence of patch-bridging attraction and lateral repulsion between adsorbed lysozymes.

321

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#### 329 **References**

- 330 [1] R. Cukalevski, S.A. Ferreira, C.J. Dunning, T. Berggård, T. Cedervall, IgG and fibrinogen  
331 driven nanoparticle aggregation, *Nano Res.* 8 (2015) 2733–2743. doi:10.1007/s12274-  
332 015-0780-4.
- 333 [2] U. Gassenschmidt, K.D. Jany, B. Tauscher, H. Niebergall, *Biochimica et Biophysica*  
334 *&ta Isolation and characterization of a flocculating protein from Moringa oleifera*  
335 *Lam, Biochim. Biophys. Acta.* 1243 (1995) 477–481. doi:10.1016/0304-4165(94)00176-  
336 X.
- 337 [3] S. Saallah, I.W. Lenggoro, *Nanoparticles carrying biological molecules: Recent advances*  
338 *and applications, KONA Powder Part. J.* 2018 (2018) 89–111.  
339 doi:10.14356/kona.2018015.
- 340 [4] B. V Derjaguin, L. Landau, *Theory of stability of highly charged lyophobic sols and*  
341 *adhesion of highly charged particles in solutions of electrolytes, Acta Physico-Chimica*  
342 *USSR.* 14 (1941) 58.

- 343 [5] E.J. Willem, J.T.G. Overbeek, J. Theodoor, G. Overbeek, Theory of the stability of  
344 lyophobic colloids, Courier Corporation, 1999.
- 345 [6] W.A. Ducker, T.J. Senden, R.M. Pashley, Direct measurement of colloidal forces using an  
346 atomic force microscope, *Nature*. 353 (1991) 239–241. doi:10.1038/353239a0.
- 347 [7] N. Helfricht, A. Mark, L. Dorwling-Carter, T. Zambelli, G. Papastavrou, Extending the  
348 limits of direct force measurements: colloidal probes from sub-micron particles,  
349 *Nanoscale*. 9 (2017) 9491–9501. doi:10.1039/C7NR02226C.
- 350 [8] S.H. Behrens, D.I. Christl, R. Emmerzael, P. Schurtenberger, M. Borkovec, Charging and  
351 Aggregation Properties of Carboxyl Latex Particles : Experiments versus DLVO Theory,  
352 *Langmuir*. 16 (2000) 2566–2575.
- 353 [9] M. Schudel, S.H. Behrens, H. Holthoff, R. Kretzschmar, M. Borkovec, Absolute  
354 Aggregation Rate Constants of Hematite Particles in Aqueous Suspensions: A  
355 Comparison of Two Different Surface Morphologies., *J. Colloid Interface Sci.* 196 (1997)  
356 241–253. doi:10.1006/jcis.1997.5207.
- 357 [10] S.B. Johnson, G. V. Franks, P.J. Scales, D. V. Boger, T.W. Healy, Surface chemistry-  
358 rheology relationships in concentrated mineral suspensions, *Int. J. Miner. Process.* 58  
359 (2000) 267–304. doi:10.1016/S0301-7516(99)00041-1.
- 360 [11] S.B. Johnson, A.S. Russell, P.J. Scales, Volume fraction effects in shear rheology and  
361 electroacoustic studies of concentrated alumina and kaolin suspensions, *Colloids Surfaces*  
362 *A Physicochem. Eng. Asp.* 141 (1998) 119–130. doi:10.1016/S0927-7757(98)00208-8.
- 363 [12] J. Gregory, Rates of flocculation of latex particles by cationic polymers, *J. Colloid*  
364 *Interface Sci.* 42 (1973) 448–456. doi:10.1016/0021-9797(73)90311-1.
- 365 [13] L. Alison, P.A. Rühls, E. Tervoort, A. Teleki, M. Zanini, L. Isa, A.R. Studart, Pickering  
366 and network stabilization of biocompatible emulsions using chitosan-modified silica  
367 nanoparticles, *Langmuir*. (2016) acs.langmuir.6b03439.  
368 doi:10.1021/acs.langmuir.6b03439.
- 369 [14] Y.-K. Leong, Inter-particle forces arising from adsorbed bolaform surfactants in colloidal  
370 suspensions, *J. Chem. Soc. Faraday Trans.* 93 (1997) 105–109. doi:10.1039/a604953b.
- 371 [15] M. Kobayashi, S. Yuki, Y. Adachi, Effect of anionic surfactants on the stability ratio and  
372 electrophoretic mobility of colloidal hematite particles, *Colloids Surfaces A Physicochem.*  
373 *Eng. Asp.* 510 (2016) 190–197. doi:10.1016/j.colsurfa.2016.07.063.
- 374 [16] Y.K. Leong, D. V. Boger, P.J. Scales, T.W. Healy, Interparticle forces arising from  
375 adsorbed surfactants in colloidal suspensions: An additional attractive force, *J. Colloid*  
376 *Interface Sci.* 181 (1996) 605–612. doi:10.1006/jcis.1996.0418.
- 377 [17] B.C. Ong, Y.K. Leong, S.B. Chen, Interparticle forces in spherical monodispersed silica  
378 dispersions: effects of branched polyethylenimine and molecular weight., *J. Colloid*

- 379 Interface Sci. 337 (2009) 24–31. doi:10.1016/j.jcis.2009.05.018.
- 380 [18] W. Lin, P. Galletto, M. Borkovec, Charging and aggregation of latex particles by  
381 oppositely charged dendrimers., *Langmuir*. 20 (2004) 7465–73. doi:10.1021/la049006i.
- 382 [19] L. Feng, Y. Adachi, A. Kobayashi, Kinetics of Brownian flocculation of polystyrene latex  
383 by cationic polyelectrolyte as a function of ionic strength, *Colloids Surfaces A*  
384 *Physicochem. Eng. Asp.* 440 (2014) 155–160. doi:10.1016/j.colsurfa.2012.09.023.
- 385 [20] Y. Adachi, L. Feng, M. Kobayashi, Kinetics of flocculation of polystyrene latex particles in  
386 the mixing flow induced with high charge density polycation near the isoelectric point,  
387 *Colloids Surfaces A Physicochem. Eng. Asp.* 471 (2015) 38–44.  
388 doi:10.1016/j.colsurfa.2015.02.011.
- 389 [21] Y. Adachi, A. Kobayashi, M. Kobayashi, Structure of colloidal flocs in relation to the  
390 dynamic properties of unstable suspension, *Int. J. Polym. Sci.* 2012 (2012).  
391 doi:10.1155/2012/574878.
- 392 [22] R.H.J. Smellie, V.K. La Mer, A Quantitative theory of filtration of flocculated  
393 suspensions, *J. Colloid Sci.* 23 (1958) 589.
- 394 [23] C. Zhao, G. Yuan, D. Jia, C.C. Han, Macrogel induced by microgel: bridging and  
395 depletion mechanisms, *Soft Matter*. 8 (2012) 7036–7043. doi:10.1039/c2sm25409c.
- 396 [24] R. Pericet-Camara, B.P. Cahill, G. Papastavrou, M. Borkovec, Nano-patterning of solid  
397 substrates by adsorbed dendrimers., *Chem. Commun. (Camb)*. (2007) 266–8.  
398 doi:10.1039/b612249c.
- 399 [25] B.P. Cahill, G. Papastavrou, G.J.M. Koper, M. Borkovec, Adsorption of poly(amido  
400 amine) (PAMAM) dendrimers on silica: importance of electrostatic three-body  
401 attraction., *Langmuir*. 24 (2008) 465–73. doi:10.1021/la7021352.
- 402 [26] K. Sofin'ska, Z. Adamczyk, M. Kujda, M. Nattich-Rak, Recombinant Albumin  
403 Monolayers on Latex Particles, *Langmuir*. 30 (2014) 250–258. doi:10.1021/la403715s.
- 404 [27] K. Kubiak-Ossowska, M. Cwieka, A. Kaczynska, B. Jachimska, P.A. Mulheran, Lysozyme  
405 adsorption at a silica surface using simulation and experiment: effects of pH on protein  
406 layer structure, *Phys. Chem. Chem. Phys.* 17 (2015) 24070–7. doi:10.1039/C5CP03910J.
- 407 [28] B. Bharti, J. Meissner, G.H. Findenegg, Aggregation of silica nanoparticles directed by  
408 adsorption of lysozyme., *Langmuir*. 27 (2011) 9823–33. doi:10.1021/la201898v.
- 409 [29] J. Meissner, A. Prause, B. Bharti, G.H. Findenegg, Characterization of protein adsorption  
410 onto silica nanoparticles: influence of pH and ionic strength, *Colloid Polym. Sci.* 293  
411 (2015) 3381–3391. doi:10.1007/s00396-015-3754-x.
- 412 [30] D. Lerche, T. Sobisch, Evaluation of particle interactions by in situ visualization of  
413 separation behaviour, *Colloids Surfaces A Physicochem. Eng. Asp.* 440 (2014) 122–130.  
414 doi:10.1016/j.colsurfa.2012.10.015.

- 415 [31] S. Kumar, V.K. Aswal, P. Callow, PH-dependent interaction and resultant structures of  
416 silica nanoparticles and lysozyme protein, *Langmuir*. 30 (2014) 1588–1598.  
417 doi:10.1021/la403896h.
- 418 [32] Y. Huang, A. Yamaguchi, T.D. Pham, M. Kobayashi, Charging and aggregation behavior  
419 of silica particles in the presence of lysozymes, *Colloid Polym. Sci.* 296 (2018) 145–155.  
420 doi:10.1007/s00396-017-4226-2.
- 421 [33] A. Otsuki, Coupling colloidal forces with yield stress of charged inorganic particle  
422 suspension: A review, *Electrophoresis*. (2018) 1–12. doi:10.1002/elps.201700314.
- 423 [34] M. Kobayashi, S. Ooi, Y. Adachi, On the Yield Stress of Sheared Coagulated Suspensions,  
424 *Proc. Hydraul. Eng.* 46 (2002) 637–640. doi:10.2208/prohe.46.637.
- 425 [35] G. V Franks, S.B. Johnson, P.J. Scales, D. V Boger, T.W. Healy, Ion-Specific Strength of  
426 Attractive Particle Networks, *Langmuir*. 15 (1999) 4411–4420. doi:10.1021/la9815345.
- 427 [36] M. Kobayashi, Y. Adachi, S. Ooi, On the Steady Shear Viscosity of Coagulated  
428 Suspensions., *Nihon Reoroji Gakkaishi*. 28 (2000) 143–144.  
429 doi:10.1678/rheology.28.143.
- 430 [37] M. Kobayashi, M. Skarba, P. Galletto, D. Cakara, M. Borkovec, Effects of heat treatment  
431 on the aggregation and charging of Stober-type silica, *J. Colloid Interface Sci.* 292 (2005)  
432 139–147. doi:10.1016/j.jcis.2005.05.093.
- 433 [38] W.F. Tan, L.K. Koopal, L.P. Weng, W.H. van Riemsdijk, W. Norde, Humic acid protein  
434 complexation, *Geochim. Cosmochim. Acta*. 72 (2008) 2090–2099.  
435 doi:10.1016/j.gca.2008.02.009.
- 436 [39] B. Jachimska, A. Kozłowska, A. Pajor-Swierzy, Protonation of lysozymes and its  
437 consequences for the adsorption onto a mica surface., *Langmuir*. 28 (2012) 11502–10.  
438 doi:10.1021/la301558u.
- 439 [40] A. Yamaguchi, M. Kobayashi, Quantitative evaluation of shift of slipping plane and  
440 counterion binding to lysozyme by electrophoresis method, *Colloid Polym. Sci.* 294  
441 (2016) 1019–1026. doi:10.1007/s00396-016-3852-4.
- 442 [41] N.Q. Dzuy, D. V. Boger, Direct Yield Stress Measurement with the Vane Method, *J.*  
443 *Rheol. (N. Y. N. Y.)*. 29 (1985) 335–347. doi:10.1122/1.549794.
- 444 [42] J. Mewis, N. Wagner, *Colloidal Suspension Rheology*, Cambridge University Press, 2011.  
445 doi:10.1017/CBO9780511977978.
- 446 [43] G. V. Franks, Zeta potentials and yield stresses of silica suspensions in concentrated  
447 monovalent electrolytes: isoelectric point shift and additional attraction., *J. Colloid*  
448 *Interface Sci.* 249 (2002) 44–51. doi:10.1006/jcis.2002.8250.
- 449 [44] M. Kobayashi, F. Juillerat, P. Galletto, P. Bowen, M. Borkovec, Aggregation and charging  
450 of colloidal silica particles: effect of particle size, *Langmuir*. 21 (2005) 5761–5769.

- 451           doi:10.1021/la046829z.
- 452    [45]   R.J. Hunter, S.K. Nicol, The dependence of plastic flow behavior of clay suspensions on  
453           surface properties, *J. Colloid Interface Sci.* 28 (1967) 250–259.
- 454    [46]   P.J. Scales, S.B. Johnson, T.W. Healy, P.C. Kapur, Shear Yield Stress of Partially  
455           Flocculated Colloidal Suspensions, *AIChE J.* 44 (1998) 538–544.  
456           doi:10.1002/aic.690440305.
- 457



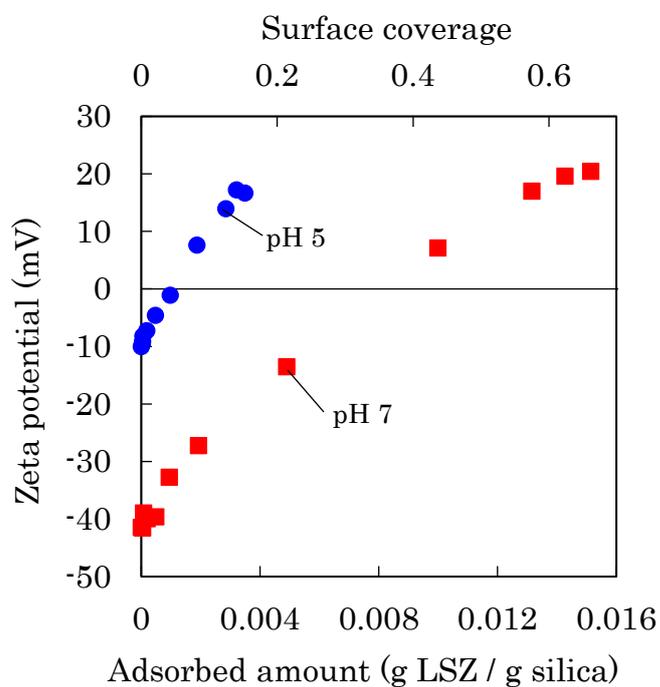


Figure 1 Zeta potential of lysozyme coated silica particles in 10 mM KCl at pH 5 and 7 are plotted against adsorbed amount (bottom) and corresponding surface coverage (upper).

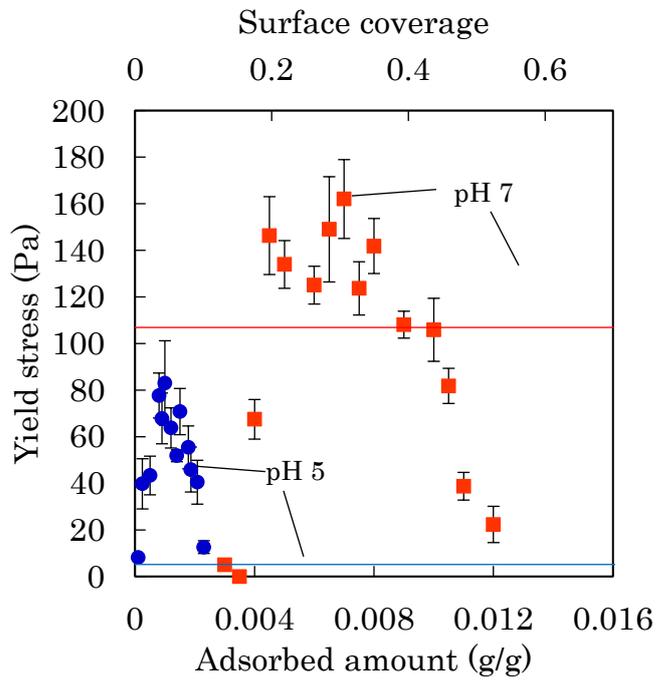


Figure 2 Yield stress of mixed suspension of lysozyme and silica particles at 10 mM KCl vs. adsorbed amount (bottom) and corresponding surface coverage (upper). The lines denote yield stress of silica suspension at 0.5 M KCl at pH 5 and at pH 7.

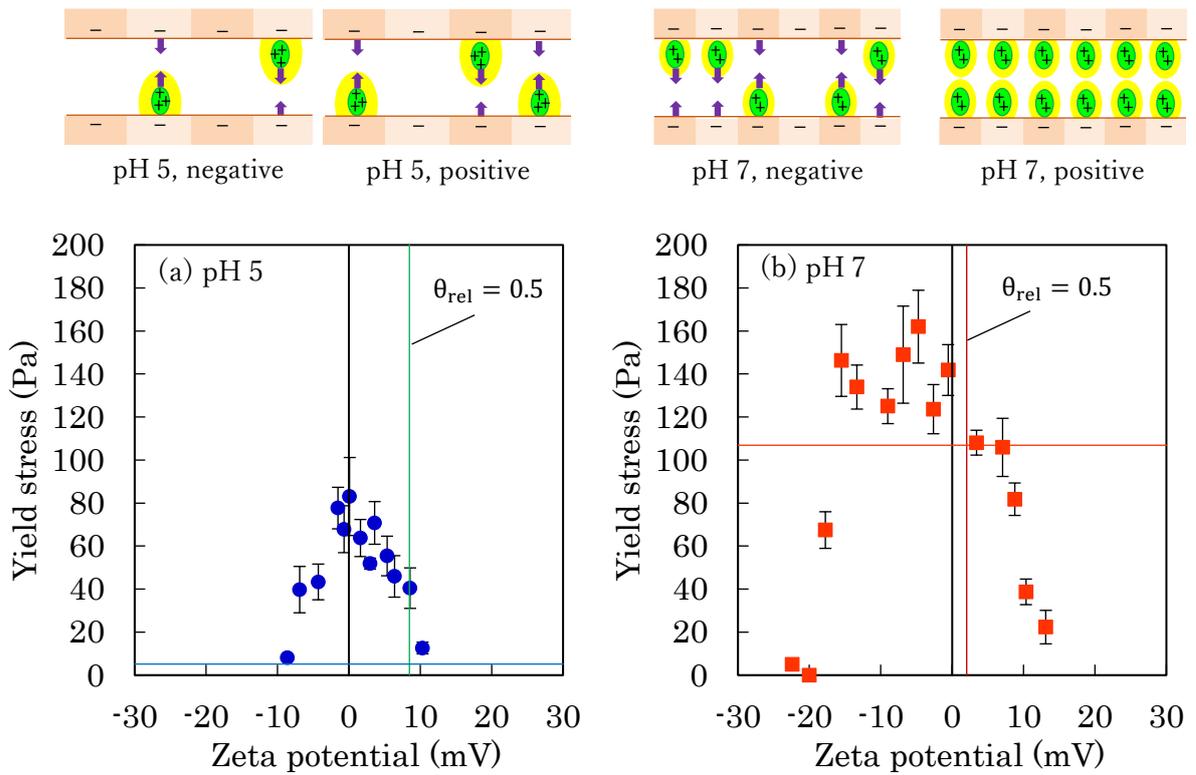


Figure 3 Yield stress of mixed suspension of silica and lysozyme at 10 mM KCl vs. zeta potential (a) at pH 5 and (b) at pH 7. The horizontal line is the yield stress of bare silica particle suspension with 0.5 M KCl. The vertical line is the zeta potential where the relative adsorbed amount equals to 0.5. The above pictures are schematic representation of gap between silica particles.

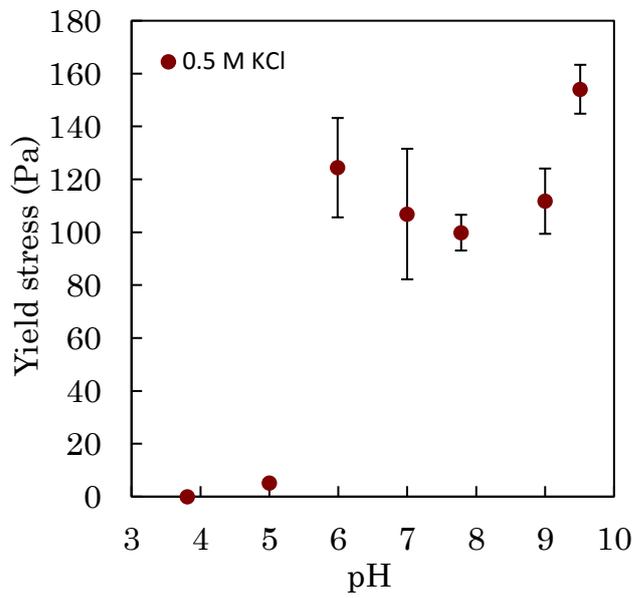


Figure 4 Yield stress of the suspension of bare silica particles at 0.5 M KCl vs. pH.

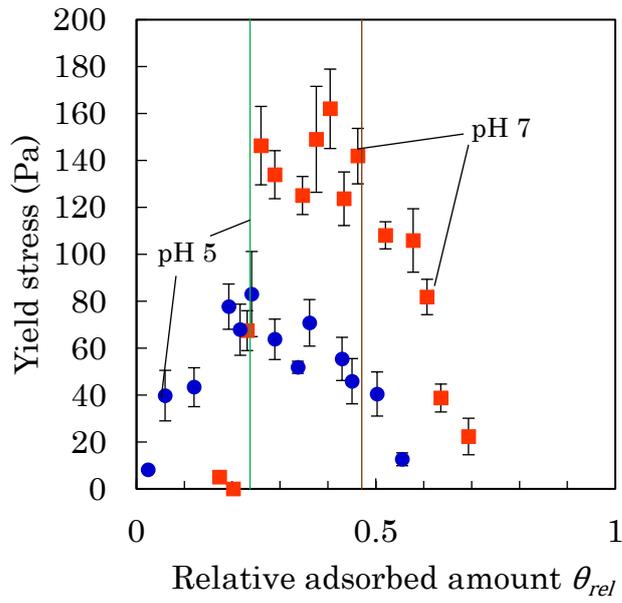


Figure 5 Yield stress of mixed suspension of lysozyme and silica at 10 mM KCl vs. relative adsorbed amount. The lines represent the relative adsorption mass corresponding to isoelectric point.