



Search for Higgs boson decays into pairs of light (pseudo)scalar particles in the $\gamma\gamma jj$ final state in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

The ATLAS Collaboration*



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ABSTRACT

This Letter presents a search for exotic decays of the Higgs boson to a pair of new (pseudo)scalar particles, $H \rightarrow aa$, where the a particle has a mass in the range 20–60 GeV, and where one of the a bosons decays into a pair of photons and the other to a pair of gluons. The search is performed in event samples enhanced in vector-boson fusion Higgs boson production by requiring two jets with large invariant mass in addition to the Higgs boson candidate decay products. The analysis is based on the full dataset of pp collisions at $\sqrt{s} = 13$ TeV recorded in 2015 and 2016 with the ATLAS detector at the CERN Large Hadron Collider, corresponding to an integrated luminosity of 36.7 fb⁻¹. The data are in agreement with the Standard Model predictions and an upper limit at the 95% confidence level is placed on the production cross section times the branching ratio for the decay $H \rightarrow aa \rightarrow \gamma\gamma gg$. This limit ranges from 3.1 pb to 9.0 pb depending on the mass of the a boson.

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1. Introduction

The discovery or exclusion of the Standard Model (SM) Higgs boson was one of the main goals of the Large Hadron Collider (LHC) physics programme. A Higgs boson with mass of 125 GeV, and with properties compatible with those expected for the SM Higgs boson (H), was discovered by the ATLAS [1] and CMS [2] collaborations. Since its discovery, a comprehensive programme of measurements of the properties of this particle has been underway. These measurements could uncover deviations from branching ratios predicted by the SM or set a limit on the possible branching ratio for decays into new particles beyond the SM (BSM). Existing measurements constrain the branching ratio for such decays (B_{BSM}) to less than 34% at 95% confidence level (CL) [3], assuming that the absolute couplings to vector bosons are smaller than or equal to the SM ones.

Many BSM models predict exotic decays of the Higgs boson [4]. One possibility is that the Higgs boson decays into a pair of new (pseudo)scalar particles, a , which in turn decay to a pair of SM particles. Several searches have been performed for $H \rightarrow aa$ in various final states [5–9].

The results presented in this Letter cover the unexplored $\gamma\gamma jj$ final state in searches for $H \rightarrow aa$, where one of the a bosons decays into a pair of photons and the other decays into a pair

of gluons. This final state becomes relevant in models where the fermionic decays are suppressed and the a boson decays only into photons or gluons [4,10]. The ATLAS Run 1 search for $H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma$ [11] set a 95% CL limit $\sigma_H \times B(H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma) < 10^{-3} \sigma_{\text{SM}}$ for $10 < m_a < 62$ GeV, where σ_{SM} is the production cross-section for the SM Higgs boson. There is currently no direct limit set on $B(H \rightarrow aa \rightarrow \gamma\gamma gg)$; however, in combination with $B_{\text{BSM}} < 34\%$, the $H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma$ result sets an indirect limit on $B(H \rightarrow aa \rightarrow \gamma\gamma gg)$ to less than $\sim 4\%$. Assuming the same ratio of photon and gluon couplings to the a boson as to the SM Higgs boson, the $H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma$ decay occurs very rarely relative to the $H \rightarrow aa \rightarrow \gamma\gamma gg$ decay (a typical value for the ratio $B(H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma)/B(H \rightarrow aa \rightarrow \gamma\gamma gg)$ is 3.8×10^{-3} [10]) making $H \rightarrow aa \rightarrow \gamma\gamma jj$ an excellent unexplored final state for probing these fermion-suppressed coupling models. The branching ratio for $a \rightarrow \gamma\gamma$ can be enhanced in some scenarios. The two searches are therefore complementary, where the $H \rightarrow aa \rightarrow \gamma\gamma jj$ final state is more sensitive to photon couplings with the new physics sector similar to the photon coupling to the SM Higgs boson, while the $H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma$ final state is more sensitive to scenarios with enhanced photon couplings. In addition, the $H \rightarrow aa \rightarrow \gamma\gamma jj$ final state can probe models inaccessible by the $H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma$ final state, for example $H \rightarrow aa' \rightarrow \gamma\gamma jj$ where the a and a' are both (pseudo)scalar particles with similar masses with primary decays to photons and gluons, respectively.

Reference [10] shows that the search for $H \rightarrow \gamma\gamma gg$, where the Higgs boson is produced in association with a vector boson which

* E-mail address: atlas.publications@cern.ch.

decays leptonically, would require approximately 300 fb^{-1} of LHC data in order to be sensitive to branching ratios less than 4%. The gluon–gluon fusion (ggF) production mode has a larger cross-section, but is overwhelmed by the $\gamma\gamma$ +multi-jet background. The strategy described in this Letter consists in selecting events where vector-boson fusion (VBF) is the dominant Higgs boson production mode. Even though the production rate is lower than that for the ggF mode, the characteristic topology of the jets produced in association with the Higgs boson enables more effective suppression of the background.

2. Data and simulation

The search presented in this Letter is based on the 36.7 fb^{-1} dataset of proton–proton collisions recorded by the ATLAS experiment at the LHC at $\sqrt{s} = 13 \text{ TeV}$ during 2015 and 2016. The ATLAS detector [12] comprises an inner detector in a 2 T axial magnetic field, for tracking charged particles and a precise localisation of the interaction vertex, a finely segmented calorimeter, a muon spectrometer and a two-level trigger [13] that accepts events at a rate of about 1 kHz for data storage.

Monte Carlo (MC) event generators were used to simulate the $H \rightarrow aa \rightarrow \gamma\gamma gg$ signal. Signal samples for the ggF and VBF processes were generated at next-to-leading order using POWHEG-Box [14–16] interfaced with PYTHIA [17] for parton showering and hadronisation using the AZNLO set of tuned parameters set [18] and the CT10 parton distribution function (PDF) set [19]. Samples were generated in the m_a range¹ $20 \text{ GeV} < m_a < 60 \text{ GeV}$, assuming the a boson to be a (pseudo)scalar. All MC event samples were processed through a detailed simulation [20] of the ATLAS detector based on GEANT4 [21], and contributions from additional pp interactions (pile-up), simulated using PYTHIA and the MSTW2008LO PDF set [22], were overlaid onto the hard-scatter events.

3. Selection criteria

Events are selected by two diphoton triggers. One trigger path requires the presence in the electromagnetic (EM) calorimeter of two clusters of energy deposits with transverse energy² above 35 GeV and 25 GeV for the leading (highest transverse energy) and sub-leading (second-highest transverse energy) clusters, respectively. In the high-level trigger the shape of the energy deposit in both clusters is required to be loosely consistent with that expected from an EM shower initiated by a photon. The other trigger path requires the presence of two clusters with transverse energy above 22 GeV. In order to suppress the additional rate due to the lower transverse energy threshold, the shape requirements for the energy deposits are more stringent.

The photon candidates are reconstructed from the clusters of energy deposits in the EM calorimeter within the range $|\eta| < 2.37$. The energies of the clusters are calibrated to account for energy losses upstream of the calorimeter and for energy leakage outside the cluster, as well as other effects due to the detector geometry and response. The calibration is refined by applying η -dependent correction factors of the order of $\pm 1\%$, derived from $Z \rightarrow ee$ events [23]. As in the trigger selection, photon candidates are required to satisfy a set of identification criteria based on the shape

of the EM cluster [24]. Two working points are defined: a *Loose* working point, used for the preselection and the data-driven background estimation, and a *Tight* working point, with requirements that further reduce the misidentification of neutral hadrons decaying to two photons. In order to reject the hadronic jet background, photon candidates are required to be isolated from any other activity in the calorimeter. The calorimeter isolation is defined as the sum of the transverse energy in the calorimeter within a cone of $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.4$ centred around the photon candidate. The transverse energy of the photon candidate is subtracted from the calorimeter isolation. Contributions to the calorimeter isolation from the underlying event and pile-up are subtracted using the method proposed in Ref. [25]. Candidates with a calorimeter isolation larger than 2.2% of the photon’s transverse energy are rejected.

Jets are reconstructed from topological clusters [26] using the anti- k_t algorithm [27] implemented in the FastJet package [28] with a radius parameter of $R = 0.4$. Jets are calibrated using an energy- and η -dependent calibration scheme, and are required to have a transverse momentum (p_T) greater than 20 GeV and $|\eta| < 2.5$ or $p_T > 30 \text{ GeV}$ and $|\eta| < 4.4$. A track- and topology-based veto [29,30] is used to suppress jets originating from pile-up interactions. Jets must have an angular separation of $\Delta R > 0.4$ from any *Loose* photon candidate in the event.

Each event is required to have at least two photon candidates whose transverse energy requirements depend on the trigger path the event follows. In each path the offline transverse energy requirements are designed so that the trigger selections are fully efficient. For events passing the trigger with higher transverse energy thresholds, the leading photon is required to have $E_T > 40 \text{ GeV}$, and the sub-leading photon is required to have $E_T > 30 \text{ GeV}$. For events passing the trigger with lower thresholds, both the leading and sub-leading photons are required to have $E_T > 27 \text{ GeV}$. For events passing both triggers, the latter selection is applied. The invariant mass of the two leading photon candidates is denoted by $m_{\gamma\gamma}$.

In the VBF production mode, the Higgs boson is produced in association with two additional light-quark jets with a large opening angle and a large invariant mass. Selected events are therefore required to have at least four jets and the pair of jets with the highest invariant mass (m_{jj}^{VBF}) are referred to as *VBF jets*. In VBF signal events, these jets correspond to the light quarks emitting the vector bosons 55% of the time, as estimated in simulation. The VBF Higgs boson signal is further enhanced, relative to the dominant $\gamma\gamma$ +multi-jet background, by requiring m_{jj}^{VBF} to be greater than 500 GeV and the p_T of the leading VBF jet to be greater than 60 GeV. The discrimination power of these observables can be seen in the difference in shape between the VBF signal and the data, shown in Figs. 1(a) and 1(b). The two remaining highest- p_T jets are referred to as *signal jets*, with invariant mass m_{jj} . The two photon candidates and the two signal jets form the Higgs boson candidate with invariant mass $m_{\gamma\gamma jj}$, which is required to be in the range $100 < m_{\gamma\gamma jj} < 150 \text{ GeV}$. Fig. 1(d) shows that most of the selected signal events lie within this range, while the data have a broad distribution extending to higher values.

In order to take advantage of the $m_{\gamma\gamma}$ resolution of about 1.3 GeV to suppress the background with $m_{\gamma\gamma}$ far from the range of interest, five overlapping $m_{\gamma\gamma}$ regimes are defined as summarised in Table 1. The boundaries of the $m_{\gamma\gamma}$ regimes are chosen so that for any value of m_a considered in the scope of this search there is at least one regime where there is no significant signal acceptance loss due to the $m_{\gamma\gamma}$ requirement. For each $m_{\gamma\gamma}$ regime, the set of m_a values for which this requirement causes no significant signal acceptance loss is also indicated.

¹ The diphoton triggers considered for this search do not have acceptance for the lower mass range ($m_a < 20 \text{ GeV}$), where the two photons are collimated.

² ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$.

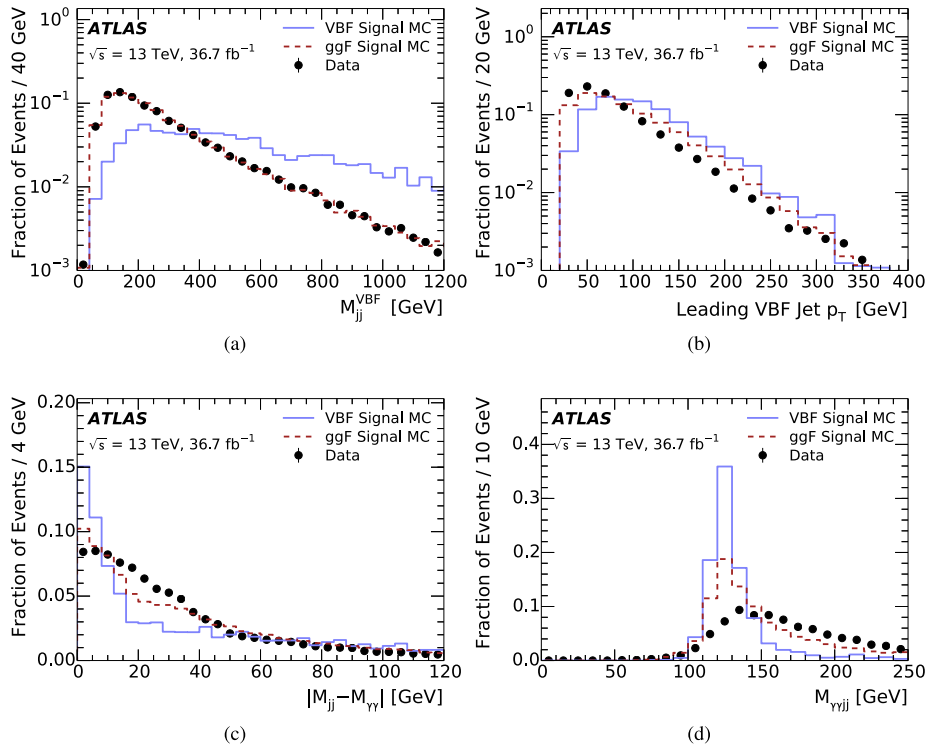


Fig. 1. Distributions of kinematic observables before the requirements on m_{jj}^{VBF} , leading VBF jet p_T , $m_{\gamma\gamma jj}$ and $|m_{jj} - m_{\gamma\gamma}|$ for: (a) m_{jj}^{VBF} ; (b) leading VBF jet p_T ; (c) $|m_{jj} - m_{\gamma\gamma}|$; and (d) $m_{\gamma\gamma jj}$ (with the additional requirement $|m_{jj} - m_{\gamma\gamma}| < 12$ GeV that defines the signal-enriched region). The quantities are shown separately for simulated signal events (with $m_a = 30$ GeV) produced in the VBF mode and compared with those produced in the ggF mode and the observed data.

4. Background estimation

The $\gamma\gamma$ +multi-jet background consists of multi-jet events with two reconstructed photon candidates, originating from isolated EM radiation or from jets. A data-driven estimation based on two-dimensional sidebands is used to predict the background yields. The method consists of using two uncorrelated observables to define four regions labelled A, B, C and D.

The first axis of the A/B/C/D plane separates events in regions C and D with both photons passing the *Tight* requirement from events in regions A and B with at most one photon passing the *Tight* requirement and at least one passing the *Loose* but not the *Tight* requirement. These regions are referred to respectively as *Tight-Tight* (C and D) and *Tight-Loose* (A and B).

The second axis separates events in regions B and D, satisfying $|m_{jj} - m_{\gamma\gamma}| < x_R$, from events in regions A and C, satisfying $|m_{jj} - m_{\gamma\gamma}| > x_R$. The value x_R depends on the $m_{\gamma\gamma}$ regime R to account for the degradation in resolution at higher mass. For $H \rightarrow aa \rightarrow \gamma\gamma gg$ signal events, where the a boson candidates have similar masses, the difference $|m_{jj} - m_{\gamma\gamma}|$ tends to be smaller than in the background, as shown in Fig. 1(c). The signal events that lie outside of the range $|m_{jj} - m_{\gamma\gamma}| < x_R$ are due to poor m_{jj} resolution or to incorrect assignment of the jets corresponding to the gluons originating from the a boson decay. Specific x_R values are given in Table 1. In each $m_{\gamma\gamma}$ regime, the boundary for $|m_{jj} - m_{\gamma\gamma}|$ is 0.4 times the central $m_{\gamma\gamma}$ value. An exception is made for the lowest $m_{\gamma\gamma}$ regime, where x_R is larger in order to increase the signal efficiency.

Region D is expected to contain the highest contribution of signal. In this region, 60% of the signal events are produced in the VBF mode and the remaining 40% in the ggF mode. Assuming no correlation in the background events between the two observables used to define the A/B/C/D regions, the number of background events in

Table 1

Definition of each $m_{\gamma\gamma}$ regime, the range of m_a values considered in the scope of this search with no significant signal loss acceptance due to the $m_{\gamma\gamma}$ requirement, and the corresponding boundary x_R for $|m_{jj} - m_{\gamma\gamma}|$.

| $m_{\gamma\gamma}$ regime | Definition | Range of m_a values | x_R [GeV] |
|---------------------------|--|---|-------------|
| 1 | $17.5 \text{ GeV} < m_{\gamma\gamma} < 27.5 \text{ GeV}$ | $20 \text{ GeV} \leq m_a \leq 25 \text{ GeV}$ | 12 |
| 2 | $22.5 \text{ GeV} < m_{\gamma\gamma} < 37.5 \text{ GeV}$ | $25 \text{ GeV} \leq m_a \leq 35 \text{ GeV}$ | 12 |
| 3 | $32.5 \text{ GeV} < m_{\gamma\gamma} < 47.5 \text{ GeV}$ | $35 \text{ GeV} \leq m_a \leq 45 \text{ GeV}$ | 16 |
| 4 | $42.5 \text{ GeV} < m_{\gamma\gamma} < 57.5 \text{ GeV}$ | $45 \text{ GeV} \leq m_a \leq 55 \text{ GeV}$ | 20 |
| 5 | $52.5 \text{ GeV} < m_{\gamma\gamma} < 65.0 \text{ GeV}$ | $55 \text{ GeV} \leq m_a \leq 60 \text{ GeV}$ | 24 |

the signal region D (N_D^{bkg}) is related to the number of background events in the control regions A, B and C, denoted by N_A^{bkg} , N_B^{bkg} and N_C^{bkg} , respectively, by the formula

$$N_D^{\text{bkg}} = \frac{N_B^{\text{bkg}} N_C^{\text{bkg}}}{N_A^{\text{bkg}}}. \quad (1)$$

In the following, the difference between the prediction N_D^{bkg} and the actual background yield in region D is referred to as *non-closure*. The non-closure results from residual correlations between the two observables used to define the A/B/C/D regions, and the uncertainty accounting for this effect is referred to as *closure uncertainty*. In order to quantify the non-closure, the data-driven estimation as described above is performed with the exception that the requirement on $m_{\gamma\gamma jj}$ is inverted. For each $m_{\gamma\gamma}$ regime, the closure uncertainty is defined to be the central value of the non-closure if it is found to be significant ($> 1\sigma$) in comparison with its statistical uncertainty; otherwise, the statistical uncertainty of its estimate is used.

Table 2

Efficiency of event selection on the inclusive $pp \rightarrow H \rightarrow aa \rightarrow \gamma\gamma gg$ signal, assuming the SM Higgs boson production cross-section and kinematics, in each of the A/B/C/D regions, for different m_a mass hypotheses. For each m_a value, all $m_{\gamma\gamma}$ regimes in which there is no significant signal acceptance loss due to the $m_{\gamma\gamma}$ requirement are shown.

| m_a [GeV] | $m_{\gamma\gamma}$ regime | Efficiency ($\times 10^{-5}$) | | | |
|----------------|------------------------------|---------------------------------|---------------------|---------------------|----------------|
| | | A | B | C | D |
| 20 | 1 | $0.50^{+0.16}_{-0.14}$ | 1.2 ± 0.4 | 3.9 ± 1.1 | 6.2 ± 1.8 |
| 25 | 1 | $0.67^{+0.27}_{-0.33}$ | $2.6^{+0.5}_{-0.6}$ | 5.8 ± 1.4 | 15 ± 4 |
| 25 | 2 | $0.67^{+0.27}_{-0.33}$ | $2.6^{+0.5}_{-0.6}$ | 5.8 ± 1.4 | 15 ± 4 |
| 30 | 2 | 1.22 ± 0.34 | 3.3 ± 0.9 | $7.6^{+1.4}_{-1.6}$ | 25^{+5}_{-6} |
| 35 | 2 | 1.8 ± 1.1 | 2.7 ± 1.2 | 9.3 ± 2.6 | 27 ± 6 |
| 35 | 3 | $0.53^{+1.20}_{-0.24}$ | 4.1 ± 1.2 | $6.1^{+1.2}_{-1.6}$ | 31 ± 7 |
| 40 | 3 | 1.2 ± 0.4 | 3.3 ± 1.0 | $7.9^{+1.7}_{-2.4}$ | 26 ± 6 |
| 45 | 3 | 2.5 ± 1.0 | 4.1 ± 1.3 | $7.7^{+1.7}_{-2.0}$ | 19 ± 5 |
| 45 | 4 | 2.2 ± 0.9 | 4.4 ± 1.4 | $5.9^{+1.5}_{-2.2}$ | 22 ± 5 |
| 50 | 4 | 0.93 ± 0.30 | 4.4 ± 1.2 | $5.0^{+1.3}_{-1.0}$ | 24 ± 5 |
| 55 | 4 | 0.37 ± 0.11 | 3.3 ± 0.9 | $5.4^{+1.3}_{-1.4}$ | 21 ± 5 |
| 55 | 5 | 0.23 ± 0.16 | 3.6 ± 1.0 | 3.4 ± 0.8 | 24 ± 6 |
| 60 | 5 | $0.77^{+0.32}_{-0.30}$ | 3.9 ± 1.0 | 4.9 ± 1.4 | 23 ± 6 |

5. Results

The efficiency of the event selection for the inclusive $pp \rightarrow H \rightarrow aa \rightarrow \gamma\gamma gg$ signal in each of the A/B/C/D regions is shown in Table 2, assuming the SM cross-section and kinematics for the ggF and VBF production modes, and the SM inclusive cross section as described in Ref. [31]; the contribution from all other production modes is expected to be negligible. The observed number of events in each of the A/B/C/D regions for each $m_{\gamma\gamma}$ regime is shown in Table 3 along with the predicted background in the signal region D, taking into account the closure uncertainty. Due to the low event counts in each of the A/B/C/D regions, the median expected background yield in region D estimated from pseudo-data experiments involving asymmetric Poisson uncertainties in the different regions slightly differs from the direct estimation from Eq. (1). No large excess is observed in region D when comparing the data yield to the background predicted from the A/B/C regions assuming that the signal is absent in these regions. However, given that a signal contamination is possible, a more refined procedure taking into account signal contributions in all regions is employed to set limits on the production rate of $H \rightarrow aa \rightarrow \gamma\gamma gg$.

A likelihood function, describing both the expected background and signal, is fit to all four A/B/C/D regions simultaneously. The free parameters of the likelihood are the numbers of signal and background events in region D, denoted μ_S and μ_{bkg} respectively, the ratio of background events expected in region B to that in region D, τ_B , and the ratio of background events expected in region C to that in region D, τ_C . The assumption of no correlation in the total background, Eq. (1), allows the background to be parameterised in terms of only three parameters. The closure uncertainty, which accounts for the uncertainty due to assuming non-correlation, is included in the likelihood function by applying a Gaussian prior to the expected number of background events in region A, $\tau_B \tau_C \mu_{\text{bkg}}$. The Gaussian width is determined by the size of the closure uncertainty summarized in Table 3. The parameter μ_S can be expressed as the product of the total integrated luminosity, the signal cross-section $\sigma_H \times B(H \rightarrow aa \rightarrow \gamma\gamma gg)$, and the signal selection efficiency estimated in MC simulation and quoted in Table 2. The signal contamination in the control regions A, B, and C is estimated

Table 3

Number of events observed in each of the A/B/C/D regions, the relative size of the closure uncertainty considered for each $m_{\gamma\gamma}$ regime, and the prediction for the number of background events in region D based on the control region yields. The median predicted background yield and its $\pm 1\sigma$ uncertainty in region D is also shown. The uncertainties in the prediction account for both the Poisson fluctuations of the number of events in the control regions and the closure uncertainty.

| $m_{\gamma\gamma}$ regime | A | B | C | D | Relative closure uncert. | Predicted background yield |
|------------------------------|----|----|----|----|-----------------------------|-------------------------------|
| 1 | 15 | 4 | 28 | 4 | 0.50 | 6^{+7}_{-4} |
| 2 | 22 | 6 | 34 | 15 | 0.32 | 8^{+7}_{-4} |
| 3 | 12 | 16 | 29 | 26 | 0.20 | 37^{+23}_{-14} |
| 4 | 8 | 12 | 19 | 38 | 0.21 | 27^{+22}_{-12} |
| 5 | 6 | 20 | 20 | 36 | 0.20 | 66^{+56}_{-28} |

from MC simulation and is varied coherently with μ_S in the likelihood fit.

The low number of observed events is the dominant source of uncertainty for this search. The second largest uncertainty is due to the closure uncertainty, also statistical in nature. Other sources of systematic uncertainty only affect the overall signal normalisation and the amount of signal contamination in control regions A, B and C. Dominant sources of experimental systematic uncertainty arise from the calibration and resolution of the energy of the jets [32,33]. Uncertainties associated with the photon energy calibration and resolution [23], as well as the photon identification and isolation efficiencies [24], are found to be negligible. Uncertainties associated with the estimation of the integrated luminosity and the simulation of pile-up interactions (*Lumi and Pile-up*) are found to be negligible. The systematic uncertainty associated with the modelling of the kinematics in signal events (*Modelling*) is evaluated by varying the choice of scales used in the generator program and assuming the SM Higgs boson production [34]. It is found to be similar in size to the experimental systematic uncertainty.

Nuisance parameters corresponding to each source of uncertainty are included in the profile likelihood with Gaussian constraints. Their effects on the estimated number of signal events μ_S are studied using Asimov [35] pseudo-datasets generated for an expected signal corresponding to the 95% CL upper limit obtained in this search and using the values of the background parameters maximising the likelihood in a fit to data which assumes no signal. Table 4 summarises the impact of each source of uncertainty varied by $\pm 1\sigma$ on the maximum-likelihood estimate for μ_S in each of the $m_{\gamma\gamma}$ regimes for an illustrative m_a hypothesis. The statistical uncertainty is the largest one for all regimes. The best-fit values of the parameters of the likelihood function are given in Table 5. The probability that the data are compatible with the background-only hypothesis is computed for each $m_{\gamma\gamma}$ regime and no significant excess is observed. The smallest local p -value, obtained for the $m_{\gamma\gamma}$ regime 2 ($m_a \approx 30$ GeV), is of the order of 4%. No significant excess is observed, and an upper limit is derived at 95% CL. The expected and observed exclusion limits on μ_S are given in Table 6. This is related to the limit on the $pp \rightarrow H \rightarrow aa \rightarrow \gamma\gamma gg$ cross-section by appropriately normalising to the measured total integrated luminosity and selection efficiencies relative to the inclusive signal production obtained from the ggF and VBF MC samples (Table 2). The limit is also expressed relative to the SM cross-section for the Higgs boson, shown in Fig. 2. Within a $m_{\gamma\gamma}$ analysis regime, limits are interpolated linearly in between simulated m_a values. Finally, for each mass point, the $m_{\gamma\gamma}$ regime that yields the best expected limit is used to provide the observed exclusion limit. The limit is calculated using a frequentist CL_s calculation [36].

Table 4

Maximum fractional impact on the fitted μ_S from sources of systematic uncertainty estimated using Asimov datasets. The signal injected in the Asimov datasets corresponds to the observed upper limit quoted in Table 6.

| Source of uncert. | $m_{\gamma\gamma}$ regime | | | | |
|-------------------|---------------------------|---------------------|---------------------|---------------------|---------------------|
| | 1 $m_a = 20$ GeV | 2 $m_a = 30$ GeV | 3 $m_a = 40$ GeV | 4 $m_a = 50$ GeV | 5 $m_a = 60$ GeV |
| Statistical | 0.73 | 0.51 | 0.89 | 1.13 | 0.92 |
| Closure | 0.44 | 0.27 | 0.39 | 0.64 | 0.89 |
| Modelling | 0.35 | 0.34 | 0.46 | 0.42 | 0.65 |
| Jet | 0.58 | 0.38 | 0.25 | 0.90 | 0.71 |
| Photon | 0.06 | 0.05 | 0.10 | 0.12 | 0.13 |
| Lumi and Pile-up | 0.06 | 0.04 | 0.27 | 0.14 | 0.32 |

Table 5

Maximum-likelihood fit values for each of the free parameters of the likelihood function in each $m_{\gamma\gamma}$ regime for a relevant signal m_a hypothesis. The estimated uncertainties in the fit parameters assume that the likelihood function is parabolic around the minimum of the fit.

| $m_{\gamma\gamma}$ regime | m_a [GeV] | μ_S | μ_{bkg} | τ_B | τ_C |
|---------------------------|-------------|----------------|--------------------|-----------------|-----------------|
| 1 | 20 | -7 ± 18 | 11 ± 17 | 0.5 ± 0.4 | 2.9 ± 3.1 |
| 2 | 30 | 8 ± 8 | 7 ± 6 | 0.68 ± 0.32 | 4.3 ± 3.1 |
| 3 | 40 | -30 ± 80 | 60 ± 70 | 0.35 ± 0.19 | 0.67 ± 0.33 |
| 4 | 50 | 22 ± 28 | 16 ± 23 | 0.5 ± 0.4 | 0.9 ± 1.0 |
| 5 | 60 | -290 ± 260 | 340 ± 340 | 0.21 ± 0.05 | 0.24 ± 0.05 |

6. Conclusions

In summary, a search for exotic decays of the Higgs boson into a pair of new (pseudo)scalar particles, $H \rightarrow aa$, in final states with two photons and two jets is conducted using 36.7 fb^{-1} of pp collisions at $\sqrt{s} = 13$ TeV recorded with the ATLAS detector at the LHC. The search for $H \rightarrow aa \rightarrow \gamma\gamma gg$ is performed in the mass range $20 < m_a < 60$ GeV and with additional jet requirements to enhance VBF-produced signal while suppressing the $\gamma\gamma$ +jets background. No significant excess of data is observed relative to the SM predictions. An upper limit is set for the product of the production cross-section for $pp \rightarrow H$ and the branching ratio for the decay $H \rightarrow aa \rightarrow \gamma\gamma gg$. The upper limit ranges from 3.1 pb to

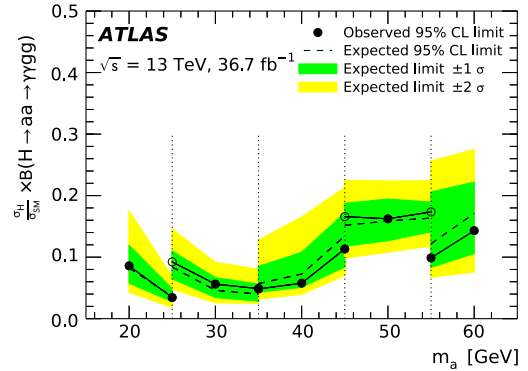


Fig. 2. The observed (solid line) and expected (dashed line) 95% CL exclusion upper limit on the $pp \rightarrow H \rightarrow aa \rightarrow \gamma\gamma gg$ cross-section times branching ratio as a function of m_a , normalised to the SM inclusive $pp \rightarrow H$ cross-section [31]. The vertical lines indicate the boundaries between the different $m_{\gamma\gamma}$ analysis regimes. At the boundaries, the $m_{\gamma\gamma}$ regime that yields the best expected limit is used to provide the observed exclusion limit (filled circles); the observed limit provided by the regime that yields the worse limit is also indicated (empty circles).

9.0 pb depending on m_a , and is mostly driven by the statistical uncertainties. These results complement the previous upper limit on $H \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma$ and further constrains the BSM parameter space for exotic decays of the Higgs boson.

Table 6

Observed (expected) upper limits at the 95% CL, for each of the m_a values considered in the search. In each case, the $m_{\gamma\gamma}$ regime used to calculate the limits is also indicated. The limits reflect both the statistical and systematic sources of uncertainty in the fit, and the $\pm 1\sigma$ widths of the expected limit distributions are also indicated.

| $m_{\gamma\gamma}$ regime | m_a [GeV] | μ_S | $\sigma_H \times B(H \rightarrow aa \rightarrow \gamma\gamma gg)$ [pb] | $\frac{\sigma_H}{\sigma_{SM}} \times B(H \rightarrow aa \rightarrow \gamma\gamma gg)$ |
|---------------------------|-------------|----------------------|--|---|
| 1 | 20 | $10.8^{+4.6}_{-3.1}$ | $4.8^{+2.1}_{-1.4}$ | $0.086^{+0.037}_{-0.025}$ |
| 1 | 25 | $10.4^{+3.8}_{-2.5}$ | $1.9^{+0.7}_{-0.5}$ | $0.034^{+0.013}_{-0.008}$ |
| 2 | 25 | 28^{+8}_{-6} | $5.1^{+1.4}_{-1.1}$ | $0.092^{+0.026}_{-0.019}$ |
| 2 | 30 | 29^{+11}_{-6} | $3.1^{+1.1}_{-0.7}$ | $0.056^{+0.021}_{-0.012}$ |
| 2 | 35 | 27^{+9}_{-6} | $2.7^{+0.9}_{-0.6}$ | $0.049^{+0.016}_{-0.011}$ |
| 3 | 35 | 30^{+18}_{-9} | $2.7^{+1.6}_{-0.8}$ | $0.048^{+0.028}_{-0.014}$ |
| 3 | 40 | 31^{+19}_{-12} | $3.2^{+2.0}_{-1.2}$ | $0.058^{+0.035}_{-0.022}$ |
| 3 | 45 | 45^{+15}_{-20} | $6.3^{+2.1}_{-2.8}$ | $0.113^{+0.038}_{-0.050}$ |
| 4 | 45 | 74^{+16}_{-15} | $9.2^{+2.0}_{-1.9}$ | $0.166^{+0.036}_{-0.034}$ |
| 4 | 50 | 79^{+17}_{-16} | $9.0^{+2.0}_{-1.8}$ | $0.162^{+0.036}_{-0.032}$ |
| 4 | 55 | 73^{+11}_{-10} | $9.7^{+1.5}_{-1.2}$ | $0.173^{+0.026}_{-0.022}$ |
| 5 | 55 | 48^{+41}_{-19} | $5.5^{+4.7}_{-2.1}$ | $0.10^{+0.08}_{-0.04}$ |
| 5 | 60 | 67^{+24}_{-31} | $8.0^{+2.8}_{-3.6}$ | $0.14^{+0.05}_{-0.07}$ |

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The ATLAS Collaboration

M. Aaboud^{34d}, G. Aad⁹⁹, B. Abbott¹²⁴, O. Abdinov^{13,*}, B. Abeloos¹²⁸, S.H. Abidi¹⁶⁴, O.S. AbouZeid¹⁴³, N.L. Abraham¹⁵³, H. Abramowicz¹⁵⁸, H. Abreu¹⁵⁷, Y. Abulaiti⁶, B.S. Acharya^{67a,67b,m}, S. Adachi¹⁶⁰, L. Adamczyk^{41a}, J. Adelman¹¹⁹, M. Adersberger¹¹², T. Adye¹⁴⁰, A.A. Affolder¹⁴³, Y. Afik¹⁵⁷, C. Agheorghiesei^{27c}, J.A. Aguilar-Saavedra^{135f,135a}, F. Ahmadov^{80,ah}, G. Aielli^{74a,74b}, S. Akatsuka⁸³, T.P.A. Åkesson⁹⁵, E. Akilli⁵⁵, A.V. Akimov¹⁰⁸, G.L. Alberghi^{23b,23a}, J. Albert¹⁷⁴, P. Albicocco⁵², M.J. Alconada Verzini⁸⁶, S. Alderweireldt¹¹⁷, M. Aleksa³⁵, I.N. Aleksandrov⁸⁰, C. Alexa^{27b}, G. Alexander¹⁵⁸, T. Alexopoulos¹⁰, M. Alhroob¹²⁴, B. Ali¹³⁷, M. Aliev^{68a,68b}, G. Alimonti^{69a}, J. Alison³⁶, S.P. Alkire³⁸, C. Allaire¹²⁸, B.M.M. Allbrooke¹⁵³, B.W. Allen¹²⁷, P.P. Allport²¹, A. Aloisio^{70a,70b}, A. Alonso³⁹, F. Alonso⁸⁶, C. Alpigiani¹⁴⁵, A.A. Alshehri⁵⁸, M.I. Alstady⁹⁹, B. Alvarez Gonzalez³⁵, D. Álvarez Piqueras¹⁷², M.G. Alviggi^{70a,70b}, B.T. Amadio¹⁸, Y. 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I. Brock²⁴, R. Brock¹⁰⁴, G. Brooijmans³⁸, T. Brooks⁹¹, W.K. Brooks^{144b}, E. Brost¹¹⁹, J.H. Broughton²¹, P.A. Bruckman de Renstrom⁴², D. Bruncko^{28b}, A. Bruni^{23b}, G. Bruni^{23b}, L.S. Bruni¹¹⁸, S. Bruno^{74a,74b}, B.H. Brunt³¹, M. Bruschi^{23b}, N. Brusino¹³⁴, P. Bryant³⁶, L. Bryngemark⁴⁶, T. Buanes¹⁷, Q. Buat³⁵, P. Buchholz¹⁴⁸, A.G. Buckley⁵⁸, I.A. Budagov⁸⁰, F. Buehrer⁵³, M.K. Bugge¹³⁰, O. Bulekov¹¹⁰, D. Bullock⁸, T.J. Burch¹¹⁹, S. Burdin⁸⁸, C.D. Burgard¹¹⁸, A.M. Burger⁵, B. Burghgrave¹¹⁹, K. Burka⁴², S. Burke¹⁴⁰, I. Burmeister⁴⁷, J.T.P. Burr¹³¹, D. Büscher⁵³, V. Büscher⁹⁷, E. Buschmann⁵⁴, P. Bussey⁵⁸, J.M. Butler²⁵, C.M. Buttar⁵⁸, J.M. Butterworth⁹², P. Butti³⁵, W. Buttinger³⁵, A. Buzatu¹⁵⁵, A.R. Buzykaev^{120b,120a}, G. Cabras^{23b,23a}, S. Cabrera Urbán¹⁷², D. Caforio¹³⁷, H. Cai¹⁷¹, V.M.M. Cairo², O. Cakir^{4a}, N. Calace⁵⁵, P. Calafiura¹⁸, A. Calandri⁹⁹, G. Calderini⁹⁴, P. Calfayan⁶⁶, G. Callea^{40b,40a}, L.P. Caloba^{141a}, S. Calvente Lopez⁹⁶, D. Calvet³⁷, S. Calvet³⁷, T.P. Calvet⁹⁹, M. Calvetti^{72a,72b}, R. Camacho Toro³⁶, S. Camarda³⁵, P. Camarri^{74a,74b}, D. Cameron¹³⁰, R. Caminal Armadans¹⁰⁰, C. Camincher⁵⁹, S. Campana³⁵, M. Campanelli⁹², A. Camplani^{69a,69b}, A. Campoverde¹⁴⁸, V. Canale^{70a,70b}, M. Cano Bret^{61c}, J. Cantero¹²⁵, T. Cao¹⁵⁸, Y. Cao¹⁷¹, M.D.M. Capeans Garrido³⁵, I. Caprini^{27b}, M. Caprini^{27b}, M. Capua^{40b,40a}, R.M. Carbone³⁸, R. Cardarelli^{74a}, F. Cardillo⁵³, I. Carli¹³⁸, T. Carli³⁵, G. Carlino^{70a}, B.T. Carlson¹³⁴, L. Carminati^{69a,69b}, R.M.D. Carney^{45a,45b}, S. Caron¹¹⁷, E. Carquin^{144b}, S. Carrá^{69a,69b}, G.D. Carrillo-Montoya³⁵, D. Casadei²¹, M.P. Casado^{14,e}, A.F. Casha¹⁶⁴, M. Casolino¹⁴, D.W. Casper¹⁶⁹, R. Castelijin¹¹⁸, V. Castillo Gimenez¹⁷², N.F. Castro^{135a}, A. Catinaccio³⁵, J.R. Catmore¹³⁰, A. Cattai³⁵, J. Caudron²⁴, V. Cavaliere²⁹, E. Cavallaro¹⁴, D. Cavalli^{69a}, M. Cavalli-Sforza¹⁴, V. Cavasinni^{72a,72b}, E. Celebi^{12b}, F. Ceradini^{75a,75b}, L. Cerda Alberich¹⁷², A.S. Cerqueira^{141b}, A. Cerri¹⁵³, L. Cerrito^{74a,74b}, F. Cerutti¹⁸, A. Cervelli^{23b,23a}, S.A. Cetin^{12b}, A. Chafaq^{34a}, DC Chakraborty¹¹⁹, S.K. Chan⁶⁰, W.S. Chan¹¹⁸, Y.L. Chan^{64a}, P. Chang¹⁷¹, J.D. Chapman³¹, D.G. Charlton²¹, C.C. Chau³³, C.A. Chavez Barajas¹⁵³, S. Che¹²², A. Chegwiddden¹⁰⁴, S. Chekanov⁶, S.V. Chekulaev^{165a}, G.A. Chelkov^{80,af}, M.A. Chelstowska³⁵, C. Chen^{61a}, C. Chen⁷⁹, H. Chen²⁹, J. Chen^{61a}, J. Chen³⁸, S. Chen¹³², S.J. Chen^{15b}, X. Chen^{15c,as}, Y. Chen⁸², H.C. Cheng¹⁰³, H.J. Cheng^{15d}, A. Cheplakov⁸⁰, E. Cheremushkina¹³⁹, R. Cherkaoui El Moursli^{34e}, E. Cheu⁷, K. Cheung⁶⁵, L. Chevalier¹⁴², V. Chiarella⁵², G. Chiarelli^{72a}, G. Chiodini^{68a}, A.S. Chisholm³⁵, A. Chitan^{27b}, I. Chiu¹⁶⁰, Y.H. Chiu¹⁷⁴, M.V. Chizhov⁸⁰, K. Choi⁶⁶, A.R. Chomont³⁷, S. Chouridou¹⁵⁹, Y.S. Chow¹¹⁸, V. Christodoulou⁹², M.C. Chu^{64a}, J. Chudoba¹³⁶, A.J. Chuinard¹⁰¹, J.J. Chwastowski⁴², L. Chytka¹²⁶, D. Cinca⁴⁷, V. Cindro⁸⁹, I.A. Cioară²⁴, A. Ciocio¹⁸, F. Ciotto^{70a,70b}, Z.H. Citron¹⁷⁸, M. Citterio^{69a}, A. Clark⁵⁵, M.R. Clark³⁸, P.J. Clark⁵⁰, R.N. Clarke¹⁸, C. Clement^{45a,45b}, Y. Coadou⁹⁹, M. Cobal^{67a,67c}, A. Coccaro^{56b,56a}, J. Cochran⁷⁹, L. Colasurdo¹¹⁷, B. Cole³⁸, A.P. Colijn¹¹⁸, J. Collot⁵⁹, P. Conde Muiño^{135a,135b}, E. Coniavitis⁵³, S.H. Connell^{32b}, I.A. Connelly⁹⁸, S. Constantinescu^{27b}, G. Conti³⁵, F. Conventi^{70a,av}, A.M. Cooper-Sarkar¹³¹, F. Cormier¹⁷³, K.J.R. Cormier¹⁶⁴, M. Corradi^{73a,73b}, E.E. Corrigan⁹⁵, F. Corriveau^{101,af}, A. Cortes-Gonzalez³⁵, M.J. Costa¹⁷², D. Costanzo¹⁴⁶, G. Cottin³¹, G. Cowan⁹¹, B.E. Cox⁹⁸, K. Cranmer¹²¹, S.J. Crawley⁵⁸, R.A. Creager¹³², G. Cree³³, S. Crépe-Renaudin⁵⁹, F. Crescioli⁹⁴, M. Cristinziani²⁴, V. Croft¹²¹, G. Crosetti^{40b,40a}, A. Cueto⁹⁶, T. Cuhadar Donszelmann¹⁴⁶, A.R. Cukierman¹⁵⁰, J. Cummings¹⁸¹, M. Curatolo⁵², J. Cúth⁹⁷, S. Czekierda⁴², P. Czodrowski³⁵, M.J. Da Cunha Sargedas De Sousa^{135a,135b}, C. Da Via⁹⁸, W. Dabrowski^{41a}, T. Dado^{28a,z}, S. Dahbi^{34e}, T. Dai¹⁰³, O. Dale¹⁷, F. Dallaire¹⁰⁷, C. Dallapiccola¹⁰⁰, M. Dam³⁹, G. D'amen^{23b,23a}, J.R. Dandoy¹³², M.F. Daneri³⁰, N.P. Dang^{179,i}, N.D. Dann⁹⁸, M. Danninger¹⁷³, M. Dano Hoffmann¹⁴², V. Dao³⁵, G. Darbo^{56b}, S. Darmora⁸, O. Dartsis⁵, A. Dattagupta¹²⁷, T. Daubney⁴⁶, S. D'Auria⁵⁸, W. Davey²⁴, C. David⁴⁶, T. Davidek¹³⁸, D.R. Davis⁴⁹, P. Davison⁹², E. Dawe¹⁰², I. Dawson¹⁴⁶, K. De⁸, R. de Asmundis^{70a}, A. De Benedetti¹²⁴, S. De Castro^{23b,23a}, S. De Cecco⁹⁴, N. De Groot¹¹⁷, P. de Jong¹¹⁸, H. De la Torre¹⁰⁴, F. De Lorenzi⁷⁹, A. De Maria^{54,r}, D. De Pedis^{73a}, A. De Salvo^{73a}, U. De Sanctis^{74a,74b}, A. De Santo¹⁵³, K. De Vasconcelos Corga⁹⁹, J.B. De Vivie De Regie¹²⁸, C. Debenedetti¹⁴³, D.V. Dedovich⁸⁰, N. Dehghanian³, I. Deigaard¹¹⁸, M. Del Gaudio^{40b,40a}, J. Del Peso⁹⁶, D. Delgove¹²⁸, F. Deliot¹⁴², C.M. Delitzsch⁷, M. Della Pietra^{70a,70b}, D. della Volpe⁵⁵, A. Dell'Acqua³⁵, L. Dell'Asta²⁵, M. Delmastro⁵, C. Delporte¹²⁸, P.A. Delsart⁵⁹, D.A. DeMarco¹⁶⁴, S. Demers¹⁸¹, M. Demichev⁸⁰, S.P. Denisov¹³⁹, D. Denysiuk¹¹⁸, L. D'Eramo⁹⁴, D. Derendarz⁴², J.E. Derkaoui^{34d}, F. Derue⁹⁴, P. Dervan⁸⁸, K. Desch²⁴, C. Deterre⁴⁶, K. Dette¹⁶⁴, M.R. Devesa³⁰, P.O. Deviveiros³⁵, A. Dewhurst¹⁴⁰, S. Dhaliwal²⁶, F.A. Di Bello⁵⁵, A. Di Ciaccio^{74a,74b}, L. Di Ciaccio⁵, W.K. Di Clemente¹³², C. Di Donato^{70a,70b}, A. Di Girolamo³⁵, B. Di Micco^{75a,75b}, R. Di Nardo³⁵, K.F. Di Petrillo⁶⁰, A. Di Simone⁵³, R. Di Sipio¹⁶⁴,

D. Di Valentino³³, C. Diaconu⁹⁹, M. Diamond¹⁶⁴, F.A. Dias³⁹, M.A. Diaz^{144a}, J. Dickinson¹⁸,
 E.B. Diehl¹⁰³, J. Dietrich¹⁹, S. Díez Cornell⁴⁶, A. Dimitrievska¹⁸, J. Dingfelder²⁴, P. Dita^{27b}, S. Dita^{27b},
 F. Dittus³⁵, F. Djama⁹⁹, T. Djobava^{156b}, J.I. Djuvsland^{62a}, M.A.B. do Vale^{141c}, M. Dobre^{27b},
 D. Dodsworth²⁶, C. Doglioni⁹⁵, J. Dolejsi¹³⁸, Z. Dolezal¹³⁸, M. Donadelli^{141d}, J. Donini³⁷,
 M. D’Onofrio⁸⁸, J. Dopke¹⁴⁰, A. Doria^{70a}, M.T. Dova⁸⁶, A.T. Doyle⁵⁸, E. Drechsler⁵⁴, E. Dreyer¹⁴⁹,
 M. Dris¹⁰, Y. Du^{61b}, J. Duarte-Campderros¹⁵⁸, F. Dubinin¹⁰⁸, A. Dubreuil⁵⁵, E. Duchovni¹⁷⁸,
 G. Duckeck¹¹², A. Ducourthial⁹⁴, O.A. Ducu^{107,y}, D. Duda¹¹⁸, A. Dudarev³⁵, A.Ch. Dudder⁹⁷,
 E.M. Duffield¹⁸, L. Duflot¹²⁸, M. Dührssen³⁵, C. Dülsen¹⁸⁰, M. Dumancic¹⁷⁸, A.E. Dumitriu^{27b,d},
 A.K. Duncan⁵⁸, M. Dunford^{62a}, A. Duperrin⁹⁹, H. Duran Yildiz^{4a}, M. Düren⁵⁷, A. Durglishvili^{156b},
 D. Duschinger⁴⁸, B. Dutta⁴⁶, D. Duvnjak¹, M. Dyndal⁴⁶, B.S. Dziedzic⁴², C. Eckardt⁴⁶, K.M. Ecker¹¹³,
 R.C. Edgar¹⁰³, T. Eifert³⁵, G. Eigen¹⁷, K. Einsweiler¹⁸, T. Ekelof¹⁷⁰, M. El Kacimi^{34c}, R. El Kosseifi⁹⁹,
 V. Ellajosyula⁹⁹, M. Ellert¹⁷⁰, F. Ellinghaus¹⁸⁰, A.A. Elliot¹⁷⁴, N. Ellis³⁵, J. Elmsheuser²⁹, M. Elsing³⁵,
 D. Emeliyanov¹⁴⁰, Y. Enari¹⁶⁰, J.S. Ennis¹⁷⁶, M.B. Epland⁴⁹, J. Erdmann⁴⁷, A. Ereditato²⁰, S. Errede¹⁷¹,
 M. Escalier¹²⁸, C. Escobar¹⁷², B. Esposito⁵², O. Estrada Pastor¹⁷², A.I. Etienvre¹⁴², E. Etzion¹⁵⁸,
 H. Evans⁶⁶, A. Ezhilov¹³³, M. Ezzi^{34e}, F. Fabbri^{23b,23a}, L. Fabbri^{23b,23a}, V. Fabiani¹¹⁷, G. Facini⁹²,
 R.M. Fakhruddinov¹³⁹, S. Falciano^{73a}, J. Faltova¹³⁸, Y. Fang^{15a}, M. Fanti^{69a,69b}, A. Farbin⁸, A. Farilla^{75a},
 E.M. Farina^{71a,71b}, T. Farooque¹⁰⁴, S. Farrell¹⁸, S.M. Farrington¹⁷⁶, P. Farthouat³⁵, F. Fassi^{34e},
 P. Fassnacht³⁵, D. Fassouliotis⁹, M. Fauci Giannelli⁵⁰, A. Favareto^{56b,56a}, W.J. Fawcett⁵⁵, L. Fayard¹²⁸,
 O.L. Fedin^{133,n}, W. Fedorko¹⁷³, M. Feickert⁴³, S. Feigl¹³⁰, L. Feligioni⁹⁹, C. Feng^{61b}, E.J. Feng³⁵,
 M. Feng⁴⁹, M.J. Fenton⁵⁸, A.B. Fenyuk¹³⁹, L. Feremenga⁸, P. Fernandez Martinez¹⁷², J. Ferrando⁴⁶,
 A. Ferrari¹⁷⁰, P. Ferrari¹¹⁸, R. Ferrari^{71a}, D.E. Ferreira de Lima^{62b}, A. Ferrer¹⁷², D. Ferrere⁵⁵,
 C. Ferretti¹⁰³, F. Fiedler⁹⁷, A. Filipčič⁸⁹, F. Filthaut¹¹⁷, M. Fincke-Keeler¹⁷⁴, K.D. Finelli²⁵,
 M.C.N. Fiolhais^{135a,135c,a}, L. Fiorini¹⁷², C. Fischer¹⁴, J. Fischer¹⁸⁰, W.C. Fisher¹⁰⁴, N. Flaschel⁴⁶,
 I. Fleck¹⁴⁸, P. Fleischmann¹⁰³, R.R.M. Fletcher¹³², T. Flick¹⁸⁰, B.M. Flierl¹¹², L.M. Flores¹³²,
 L.R. Flores Castillo^{64a}, N. Fomin¹⁷, G.T. Forcolin⁹⁸, A. Formica¹⁴², F.A. Förster¹⁴, A.C. Forti⁹⁸,
 A.G. Foster²¹, D. Fournier¹²⁸, H. Fox⁸⁷, S. Fracchia¹⁴⁶, P. Francavilla^{72a,72b}, M. Franchini^{23b,23a},
 S. Franchino^{62a}, D. Francis³⁵, L. Franconi¹³⁰, M. Franklin⁶⁰, M. Frate¹⁶⁹, M. Fraternali^{71a,71b},
 D. Freeborn⁹², S.M. Fressard-Batraneanu³⁵, B. Freund¹⁰⁷, W.S. Freund^{141a}, D. Froidevaux³⁵, J.A. Frost¹³¹,
 C. Fukunaga¹⁶¹, T. Fusayasu¹¹⁴, J. Fuster¹⁷², O. Gabizon¹⁵⁷, A. Gabrielli^{23b,23a}, A. Gabrielli¹⁸,
 G.P. Gach^{41a}, S. Gadatsch⁵⁵, S. Gadomski⁵⁵, P. Gadow¹¹³, G. Gagliardi^{56b,56a}, L.G. Gagnon¹⁰⁷,
 C. Galea¹¹⁷, B. Galhardo^{135a,135c}, E.J. Gallas¹³¹, B.J. Gallop¹⁴⁰, P. Gallus¹³⁷, G. Galster³⁹,
 R. Gamboa Goni⁹⁰, K.K. Gan¹²², S. Ganguly¹⁷⁸, Y. Gao⁸⁸, Y.S. Gao^{150,k}, C. García¹⁷²,
 J.E. García Navarro¹⁷², J.A. García Pascual^{15a}, M. Garcia-Sciveres¹⁸, R.W. Gardner³⁶, N. Garelli¹⁵⁰,
 V. Garonne¹³⁰, K. Gasnikova⁴⁶, A. Gaudiello^{56b,56a}, G. Gaudio^{71a}, I.L. Gavrilenko¹⁰⁸, C. Gay¹⁷³,
 G. Gaycken²⁴, E.N. Gazis¹⁰, C.N.P. Gee¹⁴⁰, J. Geisen⁵⁴, M. Geisen⁹⁷, M.P. Geisler^{62a}, K. Gellerstedt^{45a,45b},
 C. Gemme^{56b}, M.H. Genest⁵⁹, C. Geng¹⁰³, S. Gentile^{73a,73b}, C. Gentsos¹⁵⁹, S. George⁹¹, D. Gerbaudo¹⁴,
 G. Gessner⁴⁷, S. Ghasemi¹⁴⁸, M. Ghneimat²⁴, B. Giacobbe^{23b}, S. Giagu^{73a,73b}, N. Giangiacomi^{23b,23a},
 P. Giannetti^{72a}, S.M. Gibson⁹¹, M. Gignac¹⁴³, M. Gilchriese¹⁸, D. Gillberg³³, G. Gilles¹⁸⁰,
 D.M. Gingrich^{3,au}, M.P. Giordani^{67a,67c}, F.M. Giorgi^{23b}, P.F. Giraud¹⁴², P. Giromini⁶⁰,
 G. Giugliarelli^{67a,67c}, D. Giugni^{69a}, F. Giuli¹³¹, M. Giulini^{62b}, S. Gkaitatzis¹⁵⁹, I. Gkialas^{9,h},
 E.L. Gkougkousis¹⁴, P. Gkoutoumis¹⁰, L.K. Gladilin¹¹¹, C. Glasman⁹⁶, J. Glatzer¹⁴, P.C.F. Glaysher⁴⁶,
 A. Glazov⁴⁶, M. Goblirsch-Kolb²⁶, J. Godlewski⁴², S. Goldfarb¹⁰², T. Golling⁵⁵, D. Golubkov¹³⁹,
 A. Gomes^{135a,135b,135d}, R. Goncalves Gama^{141b}, R. Gonçalves^{135a}, G. Gonella⁵³, L. Gonella²¹,
 A. Gongadze⁸⁰, F. Gonnella²¹, J.L. Gonski⁶⁰, S. González de la Hoz¹⁷², S. Gonzalez-Sevilla⁵⁵,
 L. Goossens³⁵, P.A. Gorbounov¹⁰⁹, H.A. Gordon²⁹, B. Gorini³⁵, E. Gorini^{68a,68b}, A. Gorišek⁸⁹,
 A.T. Goshaw⁴⁹, C. Gössling⁴⁷, M.I. Gostkin⁸⁰, C.A. Gottardo²⁴, C.R. Goudet¹²⁸, D. Goujdami^{34c},
 A.G. Goussiou¹⁴⁵, N. Govender^{32b,b}, C. Goy⁵, E. Gozani¹⁵⁷, I. Grabowska-Bold^{41a}, P.O.J. Gradin¹⁷⁰,
 E.C. Graham⁸⁸, J. Gramling¹⁶⁹, E. Gramstad¹³⁰, S. Grancagnolo¹⁹, V. Gratchev¹³³, P.M. Gravila^{27f},
 C. Gray⁵⁸, H.M. Gray¹⁸, Z.D. Greenwood^{93,ak}, C. Greife²⁴, K. Gregersen⁹², I.M. Gregor⁴⁶, P. Grenier¹⁵⁰,
 K. Grevtsov⁴⁶, J. Griffiths⁸, A.A. Grillo¹⁴³, K. Grimm¹⁵⁰, S. Grinstein^{14,aa}, Ph. Gris³⁷, J.-F. Grivaz¹²⁸,
 S. Groh⁹⁷, E. Gross¹⁷⁸, J. Grosse-Knetter⁵⁴, G.C. Grossi⁹³, Z.J. Grout⁹², A. Grummer¹¹⁶, L. Guan¹⁰³,
 W. Guan¹⁷⁹, J. Guenther³⁵, A. Guerguichon¹²⁸, F. Guescini^{165a}, D. Guest¹⁶⁹, O. Gueta¹⁵⁸, R. Gugel⁵³,

T. Koffas³³, E. Koffeman¹¹⁸, N.M. Köhler¹¹³, T. Koi¹⁵⁰, M. Kolb^{62b}, I. Koletsou⁵, T. Kondo⁸¹, N. Kondrashova^{61c}, K. Köneke⁵³, A.C. König¹¹⁷, T. Kono^{81,ap}, R. Konoplich^{121,al}, N. Konstantinidis⁹², B. Konya⁹⁵, R. Kopeliansky⁶⁶, S. Koperny^{41a}, K. Korcyl⁴², K. Kordas¹⁵⁹, A. Korn⁹², I. Korolkov¹⁴, E.V. Korolkova¹⁴⁶, O. Kortner¹¹³, S. Kortner¹¹³, T. Kosek¹³⁸, V.V. Kostyukhin²⁴, A. Kotwal⁴⁹, A. Koulouris¹⁰, A. Kourkoumeli-Charalampidi^{71a,71b}, C. Kourkoumelis⁹, E. Kourlitis¹⁴⁶, V. Kouskoura²⁹, A.B. Kowalewska⁴², R. Kowalewski¹⁷⁴, T.Z. Kowalski^{41a}, C. Kozakai¹⁶⁰, W. Kozanecki¹⁴², A.S. Kozhin¹³⁹, V.A. Kramarenko¹¹¹, G. Kramberger⁸⁹, D. Krasnopevtsev¹¹⁰, M.W. Krasny⁹⁴, A. Krasznahorkay³⁵, D. Krauss¹¹³, J.A. Kremer^{41a}, J. Kretzschmar⁸⁸, K. Kreutzfeldt⁵⁷, P. Krieger¹⁶⁴, K. Krizka¹⁸, K. Kroeninger⁴⁷, H. Kroha¹¹³, J. Kroll¹³⁶, J. Kroll¹³², J. Kroseberg²⁴, J. Krstic¹⁶, U. Kruchonak⁸⁰, H. Krüger²⁴, N. Krumnack⁷⁹, M.C. Kruse⁴⁹, T. Kubota¹⁰², S. Kудay^{4b}, J.T. Kuechler¹⁸⁰, S. Kuehn³⁵, A. Kugel^{62a}, F. Kuger¹⁷⁵, T. Kuhl⁴⁶, V. Kukhtin⁸⁰, R. Kukla⁹⁹, Y. Kulchitsky¹⁰⁵, S. Kuleshov^{144b}, Y.P. Kulinich¹⁷¹, M. Kuna⁵⁹, T. Kunigo⁸³, A. Kupco¹³⁶, T. Kupfer⁴⁷, O. Kuprash¹⁵⁸, H. Kurashige⁸², L.L. Kurchaninov^{165a}, Y.A. Kurochkin¹⁰⁵, M.G. Kurth^{15d}, E.S. Kuwertz¹⁷⁴, M. Kuze¹⁶², J. Kvita¹²⁶, T. Kwan¹⁷⁴, A. La Rosa¹¹³, J.L. La Rosa Navarro^{141d}, L. La Rotonda^{40b,40a}, F. La Ruffa^{40b,40a}, C. Lacasta¹⁷², F. Lacava^{73a,73b}, J. Lacey⁴⁶, D.P.J. Lack⁹⁸, H. Lacker¹⁹, D. Lacour⁹⁴, E. Ladygin⁸⁰, R. Lafaye⁵, B. Laforge⁹⁴, S. Lai⁵⁴, S. Lammers⁶⁶, W. Lampl⁷, E. Lançon²⁹, U. Landgraf⁵³, M.P.J. Landon⁹⁰, M.C. Lanfermann⁵⁵, V.S. Lang⁴⁶, J.C. Lange¹⁴, R.J. Langenberg³⁵, A.J. Lankford¹⁶⁹, F. Lanni²⁹, K. Lantzsch²⁴, A. Lanza^{71a}, A. Lapertosa^{56b,56a}, S. Laplace⁹⁴, J.F. Laporte¹⁴², T. Lari^{69a}, F. Lasagni Manghi^{23b,23a}, M. Lassnig³⁵, T.S. Lau^{64a}, A. Laudrain¹²⁸, A.T. Law¹⁴³, P. Laycock⁸⁸, M. Lazzaroni^{69a,69b}, B. Le¹⁰², O. Le Dortz⁹⁴, E. Le Guirriec⁹⁹, E.P. Le Quilleuc¹⁴², M. LeBlanc⁷, T. LeCompte⁶, F. Ledroit-Guillon⁵⁹, C.A. Lee²⁹, G.R. Lee^{144a}, L. Lee⁶⁰, S.C. Lee¹⁵⁵, B. Lefebvre¹⁰¹, M. Lefebvre¹⁷⁴, F. Legger¹¹², C. Leggett¹⁸, G. Lehmann Miotto³⁵, W.A. Leight⁴⁶, A. Leisos^{159,x}, M.A.L. Leite^{141d}, R. Leitner¹³⁸, D. Lellouch¹⁷⁸, B. Lemmer⁵⁴, K.J.C. Leney⁹², T. Lenz²⁴, B. Lenzi³⁵, R. Leone⁷, S. Leone^{72a}, C. Leonidopoulos⁵⁰, G. Lerner¹⁵³, C. Leroy¹⁰⁷, R. Les¹⁶⁴, A.A.J. Lesage¹⁴², C.G. Lester³¹, M. Levchenko¹³³, J. Levêque⁵, D. Levin¹⁰³, L.J. Levinson¹⁷⁸, M. Levy²¹, D. Lewis⁹⁰, B. Li^{61a,q}, C-Q. Li^{61a}, H. Li^{61b}, L. Li^{61c}, Q. Li^{15d}, Q. Li^{61a}, S. Li^{61d,61c}, X. Li^{61c}, Y. Li¹⁴⁸, Z. Liang^{15a}, B. Liberti^{74a}, A. Liblong¹⁶⁴, K. Lie^{64c}, A. Limosani¹⁵⁴, C.Y. Lin³¹, K. Lin¹⁰⁴, S.C. Lin¹⁶⁸, T.H. Lin⁹⁷, R.A. Linck⁶⁶, B.E. Lindquist¹⁵², A.L. Lioni⁵⁵, E. Lipeles¹³², A. Lipniacka¹⁷, M. Lisovsky^{62b}, T.M. Liss^{171,ar}, A. Lister¹⁷³, A.M. Litke¹⁴³, J.D. Little⁸, B. Liu⁷⁹, H. Liu²⁹, H. Liu¹⁰³, J.B. Liu^{61a}, J.K.K. Liu¹³¹, K. Liu⁹⁴, M. Liu^{61a}, P. Liu¹⁸, Y. Liu^{61a}, Y.L. Liu^{61a}, M. Livan^{71a,71b}, A. Lleres⁵⁹, J. Llorente Merino^{15a}, S.L. Lloyd⁹⁰, C.Y. Lo^{64b}, F. Lo Sterzo⁴³, E.M. Lobodzinska⁴⁶, P. Loch⁷, F.K. Loebinger⁹⁸, A. Loesle⁵³, K.M. Loew²⁶, T. Lohse¹⁹, K. Lohwasser¹⁴⁶, M. Lokajicek¹³⁶, B.A. Long²⁵, J.D. Long¹⁷¹, R.E. Long⁸⁷, L. Longo^{68a,68b}, K.A. Looper¹²², J.A. Lopez^{144b}, I. Lopez Paz¹⁴, A. Lopez Solis⁹⁴, J. Lorenz¹¹², N. Lorenzo Martinez⁵, M. Losada²², P.J. Lösel¹¹², X. Lou^{15a}, A. Lounis¹²⁸, J. Love⁶, P.A. Love⁸⁷, H. Lu^{64a}, N. Lu¹⁰³, Y.J. Lu⁶⁵, H.J. Lubatti¹⁴⁵, C. Luci^{73a,73b}, A. Lucotte⁵⁹, C. Luedtke⁵³, F. Luehring⁶⁶, W. Lukas⁷⁷, L. Luminari^{73a}, B. Lund-Jensen¹⁵¹, M.S. Lutz¹⁰⁰, P.M. Luzzi⁹⁴, D. Lynn²⁹, R. Lysak¹³⁶, E. Lytken⁹⁵, F. Lyu^{15a}, V. Lyubushkin⁸⁰, H. Ma²⁹, L.L. Ma^{61b}, Y. Ma^{61b}, G. Maccarrone⁵², A. Macchiolo¹¹³, C.M. Macdonald¹⁴⁶, J. Machado Miguens^{132,135b}, D. Madaffari¹⁷², R. Madar³⁷, W.F. Mader⁴⁸, A. Madsen⁴⁶, N. Madysa⁴⁸, J. Maeda⁸², S. Maeland¹⁷, T. Maeno²⁹, A.S. Maevskiy¹¹¹, V. Magerl⁵³, C. Maidantchik^{141a}, T. Maier¹¹², A. Maio^{135a,135b,135d}, O. Majersky^{28a}, S. Majewski¹²⁷, Y. Makida⁸¹, N. Makovec¹²⁸, B. Malaescu⁹⁴, Pa. Malecki⁴², V.P. Maleev¹³³, F. Malek⁵⁹, U. Mallik⁷⁸, D. Malon⁶, C. Malone³¹, S. Maltezos¹⁰, S. Malyukov³⁵, J. Mamuzic¹⁷², G. Mancini⁵², I. Mandić⁸⁹, J. Maneira^{135a,135b}, L. Manhaes de Andrade Filho^{141b}, J. Manjarres Ramos⁴⁸, K.H. Mankinen⁹⁵, A. Mann¹¹², A. Manousos³⁵, B. Mansoulie¹⁴², J.D. Mansour^{15a}, R. Mantifel¹⁰¹, M. Mantoani⁵⁴, S. Manzoni^{69a,69b}, G. Marceca³⁰, L. March⁵⁵, L. Marchese¹³¹, G. Marchiori⁹⁴, M. Marcisovsky¹³⁶, C.A. Marin Tobon³⁵, M. Marjanovic³⁷, D.E. Marley¹⁰³, F. Marroquim^{141a}, Z. Marshall¹⁸, M.U.F. Martensson¹⁷⁰, S. Marti-Garcia¹⁷², C.B. Martin¹²², T.A. Martin¹⁷⁶, V.J. Martin⁵⁰, B. Martin dit Latour¹⁷, M. Martinez^{14,aa}, V.I. Martinez Outschoorn¹⁰⁰, S. Martin-Haugh¹⁴⁰, V.S. Martoiu^{27b}, A.C. Martyniuk⁹², A. Marzin³⁵, L. Masetti⁹⁷, T. Mashimo¹⁶⁰, R. Mashinistov¹⁰⁸, J. Masik⁹⁸, A.L. Maslennikov^{120b,120a}, L.H. Mason¹⁰², L. Massa^{74a,74b}, P. Mastrandrea⁵, A. Mastroberardino^{40b,40a}, T. Masubuchi¹⁶⁰, P. Mättig¹⁸⁰, J. Maurer^{27b}, B. Maček⁸⁹, S.J. Maxfield⁸⁸, D.A. Maximov^{120b,120a}, R. Mazini¹⁵⁵, I. Maznas¹⁵⁹, S.M. Mazza¹⁴³, N.C. Mc Fadden¹¹⁶, G. Mc Goldrick¹⁶⁴, S.P. Mc Kee¹⁰³, A. McCann¹⁰³, T.G. McCarthy¹¹³, L.I. McClymont⁹², E.F. McDonald¹⁰²,

J.A. McFayden³⁵, G. Mchedlidze⁵⁴, M.A. McKay⁴³, S.J. McMahon¹⁴⁰, P.C. McNamara¹⁰², C.J. McNicol¹⁷⁶,
 R.A. McPherson^{174,af}, Z.A. Meadows¹⁰⁰, S. Meehan¹⁴⁵, T. Megy⁵³, S. Mehlhase¹¹², A. Mehta⁸⁸,
 T. Meideck⁵⁹, B. Meirose⁴⁴, D. Melini^{172,f}, B.R. Mellado Garcia^{32c}, J.D. Mellenthin⁵⁴, M. Melo^{28a},
 F. Meloni²⁰, A. Melzer²⁴, S.B. Menary⁹⁸, L. Meng⁸⁸, X.T. Meng¹⁰³, A. Mengarelli^{23b,23a}, S. Menke¹¹³,
 E. Meoni^{40b,40a}, S. Mergelmeyer¹⁹, C. Merlassino²⁰, P. Mermod⁵⁵, L. Merola^{70a,70b}, C. Meroni^{69a},
 F.S. Merritt³⁶, A. Messina^{73a,73b}, J. Metcalfe⁶, A.S. Mete¹⁶⁹, C. Meyer¹³², J. Meyer¹¹⁸, J.-P. Meyer¹⁴²,
 H. Meyer Zu Theenhausen^{62a}, F. Miano¹⁵³, R.P. Middleton¹⁴⁰, S. Miglioranza^{56b,56a}, L. Mijović⁵⁰,
 G. Mikenberg¹⁷⁸, M. Mikestikova¹³⁶, M. Mikuž⁸⁹, M. Milesi¹⁰², A. Milic¹⁶⁴, D.A. Millar⁹⁰, D.W. Miller³⁶,
 A. Milov¹⁷⁸, D.A. Milstead^{45a,45b}, A.A. Minaenko¹³⁹, I.A. Minashvili^{156b}, A.I. Mincer¹²¹, B. Mindur^{41a},
 M. Mineev⁸⁰, Y. Minegishi¹⁶⁰, Y. Ming¹⁷⁹, L.M. Mir¹⁴, A. Mirto^{68a,68b}, K.P. Mistry¹³², T. Mitani¹⁷⁷,
 J. Mitrevski¹¹², V.A. Mitsou¹⁷², A. Miucci²⁰, P.S. Miyagawa¹⁴⁶, A. Mizukami⁸¹, J.U. Mjörnmark⁹⁵,
 T. Mkrtchyan¹⁸², M. Mlynarikova¹³⁸, T. Moa^{45a,45b}, K. Mochizuki¹⁰⁷, P. Mogg⁵³, S. Mohapatra³⁸,
 S. Molander^{45a,45b}, R. Moles-Valls²⁴, M.C. Mondragon¹⁰⁴, K. Mönig⁴⁶, J. Monk³⁹, E. Monnier⁹⁹,
 A. Montalbano¹⁴⁹, J. Montejo Berlingen³⁵, F. Monticelli⁸⁶, S. Monzani^{69a}, R.W. Moore³, N. Morange¹²⁸,
 D. Moreno²², M. Moreno Llácer³⁵, P. Morettini^{56b}, M. Morgenstern¹¹⁸, S. Morgenstern³⁵, D. Mori¹⁴⁹,
 T. Mori¹⁶⁰, M. Morii⁶⁰, M. Morinaga¹⁷⁷, V. Morisbak¹³⁰, A.K. Morley³⁵, G. Mornacchi³⁵, J.D. Morris⁹⁰,
 L. Morvaj¹⁵², P. Moschovakos¹⁰, M. Mosidze^{156b}, H.J. Moss¹⁴⁶, J. Moss^{150,l}, K. Motohashi¹⁶²,
 R. Mount¹⁵⁰, E. Mountricha²⁹, E.J.W. Moyse¹⁰⁰, S. Muanza⁹⁹, F. Mueller¹¹³, J. Mueller¹³⁴,
 R.S.P. Mueller¹¹², D. Muenstermann⁸⁷, P. Mullen⁵⁸, G.A. Mullier²⁰, F.J. Munoz Sanchez⁹⁸, P. Murin^{28b},
 W.J. Murray^{176,140}, A. Murrone^{69a,69b}, M. Muškinja⁸⁹, C. Mwewa^{32a}, A.G. Myagkov^{139,am}, J. Myers¹²⁷,
 M. Myska¹³⁷, B.P. Nachman¹⁸, O. Nackenhorst⁴⁷, K. Nagai¹³¹, R. Nagai^{81,ap}, K. Nagano⁸¹, Y. Nagasaka⁶³,
 K. Nagata¹⁶⁶, M. Nagel⁵³, E. Nagy⁹⁹, A.M. Nairz³⁵, Y. Nakahama¹¹⁵, K. Nakamura⁸¹, T. Nakamura¹⁶⁰,
 I. Nakano¹²³, R.F. Naranjo Garcia⁴⁶, R. Narayan¹¹, D.I. Narrias Villar^{62a}, I. Naryshkin¹³³, T. Naumann⁴⁶,
 G. Navarro²², R. Nayyar⁷, H.A. Neal¹⁰³, P.Yu. Nechaeva¹⁰⁸, T.J. Neep¹⁴², A. Negri^{71a,71b}, M. Negrini^{23b},
 S. Nektarijevic¹¹⁷, C. Nellist⁵⁴, M.E. Nelson¹³¹, S. Nemecek¹³⁶, P. Nemethy¹²¹, M. Nessi^{35,g},
 M.S. Neubauer¹⁷¹, M. Neumann¹⁸⁰, P.R. Newman²¹, T.Y. Ng^{64c}, Y.S. Ng¹⁹, H.D.N. Nguyen⁹⁹,
 T. Nguyen Manh¹⁰⁷, R.B. Nickerson¹³¹, R. Nicolaidou¹⁴², J. Nielsen¹⁴³, N. Nikiporou¹¹,
 V. Nikolaenko^{139,am}, I. Nikolic-Audit⁹⁴, K. Nikolopoulos²¹, P. Nilsson²⁹, Y. Ninomiya⁸¹, A. Nisati^{73a},
 N. Nishu^{61c}, R. Nisius¹¹³, I. Nitsche⁴⁷, T. Nitta¹⁷⁷, T. Nobe¹⁶⁰, Y. Noguchi⁸³, M. Nomachi¹²⁹,
 I. Nomidis³³, M.A. Nomura²⁹, T. Nooney⁹⁰, M. Nordberg³⁵, N. Norjoharuddeen¹³¹, T. Novak⁸⁹,
 O. Novgorodova⁴⁸, R. Novotny¹³⁷, M. Nozaki⁸¹, L. Nozka¹²⁶, K. Ntekas¹⁶⁹, E. Nurse⁹², F. Nuti¹⁰²,
 F.G. Oakham^{33,au}, H. Oberlack¹¹³, T. Obermann²⁴, J. Ocariz⁹⁴, A. Ochi⁸², I. Ochoa³⁸,
 J.P. Ochoa-Ricoux^{144a}, K. O'Connor²⁶, S. Oda⁸⁵, S. Odaka⁸¹, A. Oh⁹⁸, S.H. Oh⁴⁹, C.C. Ohm¹⁵¹,
 H. Ohman¹⁷⁰, H. Oide^{56b,56a}, H. Okawa¹⁶⁶, Y. Okumura¹⁶⁰, T. Okuyama⁸¹, A. Olariu^{27b},
 L.F. Oleiro Seabra^{135a}, S.A. Olivares Pino^{144a}, D. Oliveira Damazio²⁹, J.L. Oliver¹, M.J.R. Olsson³⁶,
 A. Olszewski⁴², J. Olszowska⁴², D.C. O'Neil¹⁴⁹, A. Onofre^{135a,135e}, K. Onogi¹¹⁵, P.U.E. Onyisi¹¹,
 H. Oppen¹³⁰, M.J. Oreglia³⁶, Y. Oren¹⁵⁸, D. Orestano^{75a,75b}, E.C. Orgill⁹⁸, N. Orlando^{64b},
 A.A. O'Rourke⁴⁶, R.S. Orr¹⁶⁴, B. Osculati^{56b,56a,*}, V. O'Shea⁵⁸, R. Ospanov^{61a}, G. Otero y Garzon³⁰,
 H. Otono⁸⁵, M. Ouchrif^{34d}, F. Ould-Saada¹³⁰, A. Ouraou¹⁴², K.P. Oussoren¹¹⁸, Q. Ouyang^{15a}, M. Owen⁵⁸,
 R.E. Owen²¹, V.E. Ozcan^{12c}, N. Ozturk⁸, K. Pachal¹⁴⁹, A. Pacheco Pages¹⁴, L. Pacheco Rodriguez¹⁴²,
 C. Padilla Aranda¹⁴, S. Pagan Griso¹⁸, M. Paganini¹⁸¹, F. Paige²⁹, G. Palacino⁶⁶, S. Palazzo^{40b,40a},
 S. Palestini³⁵, M. Palka^{41b}, D. Pallin³⁷, E.St. Panagiotopoulou¹⁰, I. Panagoulas¹⁰, C.E. Pandini⁵⁵,
 J.G. Panduro Vazquez⁹¹, P. Pani³⁵, D. Pantea^{27b}, L. Paolozzi⁵⁵, Th.D. Papadopoulou¹⁰,
 K. Papageorgiou^{9,h}, A. Paramonov⁶, D. Paredes Hernandez^{64b}, B. Parida^{61c}, A.J. Parker⁸⁷, K.A. Parker⁴⁶,
 M.A. Parker³¹, F. Parodi^{56b,56a}, J.A. Parsons³⁸, U. Parzefall⁵³, V.R. Pascuzzi¹⁶⁴, J.M.P. Pasner¹⁴³,
 E. Pasqualucci^{73a}, S. Passaggio^{56b}, Fr. Pastore⁹¹, S. Patariaia⁹⁷, J.R. Pater⁹⁸, T. Pauly³⁵, B. Pearson¹¹³,
 S. Pedraza Lopez¹⁷², R. Pedro^{135a,135b}, S.V. Peleganchuk^{120b,120a}, O. Penc¹³⁶, C. Peng^{15d}, H. Peng^{61a},
 J. Penwell⁶⁶, B.S. Peralva^{141b}, M.M. Perego¹⁴², D.V. Perepelitsa²⁹, F. Peri¹⁹, L. Perini^{69a,69b},
 H. Pernegger³⁵, S. Perrella^{70a,70b}, V.D. Peshekhonov^{80,*}, K. Peters⁴⁶, R.F.Y. Peters⁹⁸, B.A. Petersen³⁵,
 T.C. Petersen³⁹, E. Petit⁵⁹, A. Petridis¹, C. Petridou¹⁵⁹, P. Petroff¹²⁸, E. Petrolo^{73a}, M. Petrov¹³¹,
 F. Petrucci^{75a,75b}, N.E. Pettersson¹⁰⁰, A. Peyaud¹⁴², R. Pezoa^{144b}, T. Pham¹⁰², F.H. Phillips¹⁰⁴,
 P.W. Phillips¹⁴⁰, G. Piacquadio¹⁵², E. Pianori¹⁷⁶, A. Picazio¹⁰⁰, M.A. Pickering¹³¹, R. Piegaia³⁰,

J.E. Pilcher³⁶, A.D. Pilkington⁹⁸, M. Pinamonti^{74a,74b}, J.L. Pinfeld³, M. Pitt¹⁷⁸, M-A. Pleier²⁹, V. Pleskot¹³⁸, E. Plotnikova⁸⁰, D. Pluth⁷⁹, P. Podberezko^{120b,120a}, R. Poettgen⁹⁵, R. Poggi^{71a,71b}, L. Poggioli¹²⁸, I. Pogrebnyak¹⁰⁴, D. Pohl²⁴, I. Pokharel⁵⁴, G. Polesello^{71a}, A. Poley⁴⁶, A. Policicchio^{40b,40a}, R. Polifka³⁵, A. Polini^{23b}, C.S. Pollard⁴⁶, V. Polychronakos²⁹, D. Ponomarenko¹¹⁰, L. Pontecorvo^{73a}, G.A. Popeneciu^{27d}, D.M. Portillo Quintero⁹⁴, S. Pospisil¹³⁷, K. Potamianos⁴⁶, I.N. Potrap⁸⁰, C.J. Potter³¹, H. Potti¹¹, T. Poulsen⁹⁵, J. Poveda³⁵, M.E. Pozo Astigarraga³⁵, P. Pralavorio⁹⁹, S. Prell⁷⁹, D. Price⁹⁸, M. Primavera^{68a}, S. Prince¹⁰¹, N. Proklova¹¹⁰, K. Prokofiev^{64c}, F. Prokoshin^{144b}, S. Protopopescu²⁹, J. Proudfoot⁶, M. Przybycien^{41a}, A. Puri¹⁷¹, P. Puzo¹²⁸, J. Qian¹⁰³, Y. Qin⁹⁸, A. Quadt⁵⁴, M. Queitsch-Maitland⁴⁶, A. Qureshi¹, V. Radeka²⁹, S.K. Radhakrishnan¹⁵², P. Rados¹⁰², F. Ragusa^{69a,69b}, G. Rahal⁵¹, J.A. Raine⁹⁸, S. Rajagopalan²⁹, T. Rashid¹²⁸, S. Raspopov⁵, M.G. Ratti^{69a,69b}, D.M. Rauch⁴⁶, F. Rauscher¹¹², S. Rave⁹⁷, I. Ravinovich¹⁷⁸, J.H. Rawling⁹⁸, M. Raymond³⁵, A.L. Read¹³⁰, N.P. Readoff⁵⁹, M. Reale^{68a,68b}, D.M. Rebuzzi^{71a,71b}, A. Redelbach¹⁷⁵, G. Redlinger²⁹, R. Reece¹⁴³, R.G. Reed^{32c}, K. Reeves⁴⁴, L. Rehnisch¹⁹, J. Reichert¹³², A. Reiss⁹⁷, C. Rembser³⁵, H. Ren^{15d}, M. Rescigno^{73a}, S. Resconi^{69a}, E.D. Resseguie¹³², S. Rettie¹⁷³, E. Reynolds²¹, O.L. Rezanova^{120b,120a}, P. Reznicek¹³⁸, R. Richter¹¹³, S. Richter⁹², E. Richter-Was^{41b}, O. Ricken²⁴, M. Ridel⁹⁴, P. Rieck¹¹³, C.J. Riegel¹⁸⁰, O. Rifki⁴⁶, M. Rijssenbeek¹⁵², A. Rimoldi^{71a,71b}, M. Rimoldi²⁰, L. Rinaldi^{23b}, G. Ripellino¹⁵¹, B. Ristić³⁵, E. Ritsch³⁵, I. Riu¹⁴, J.C. Rivera Vergara^{144a}, F. Rizatdinova¹²⁵, E. Rizvi⁹⁰, C. Rizzi¹⁴, R.T. Roberts⁹⁸, S.H. Robertson^{101,af}, A. Robichaud-Veronneau¹⁰¹, D. Robinson³¹, J.E.M. Robinson⁴⁶, A. Robson⁵⁸, E. Rocco⁹⁷, C. Roda^{72a,72b}, Y. Rodina^{99,ab}, S. Rodriguez Bosca¹⁷², A. Rodriguez Perez¹⁴, D. Rodriguez Rodriguez¹⁷², A.M. Rodríguez Vera^{165b}, S. Roe³⁵, C.S. Rogan⁶⁰, O. Røhne¹³⁰, R. Röhrig¹¹³, J. Roloff⁶⁰, A. Romaniouk¹¹⁰, M. Romano^{23b,23a}, S.M. Romano Saez³⁷, E. Romero Adam¹⁷², N. Rompotis⁸⁸, M. Ronzani⁵³, L. Roos⁹⁴, S. Rosati^{73a}, K. Rosbach⁵³, P. Rose¹⁴³, N-A. Rosien⁵⁴, E. Rossi^{70a,70b}, L.P. Rossi^{56b}, L. Rossini^{69a,69b}, J.H.N. Rosten³¹, R. Rosten¹⁴⁵, M. Rotaru^{27b}, J. Rothberg¹⁴⁵, D. Rousseau¹²⁸, D. Roy^{32c}, A. Rozanov⁹⁹, Y. Rozen¹⁵⁷, X. Ruan^{32c}, F. Rubbo¹⁵⁰, F. Rühr⁵³, A. Ruiz-Martinez³³, Z. Rurikova⁵³, N.A. Rusakovich⁸⁰, H.L. Russell¹⁰¹, J.P. Rutherford⁷, N. Ruthmann³⁵, E.M. Rüttinger^{46,j}, Y.F. Ryabov¹³³, M. Rybar¹⁷¹, G. Rybkin¹²⁸, S. Ryu⁶, A. Ryzhov¹³⁹, G.F. Rzehorz⁵⁴, G. Sabato¹¹⁸, S. Sacerdoti¹²⁸, H.F-W. Sadrozinski¹⁴³, R. Sadykov⁸⁰, F. Safai Tehrani^{73a}, P. Saha¹¹⁹, M. Sahinsoy^{62a}, M. Saimpert⁴⁶, M. Saito¹⁶⁰, T. Saito¹⁶⁰, H. Sakamoto¹⁶⁰, A. Sakharov^{121,al}, G. Salamanna^{75a,75b}, J.E. Salazar Loyola^{144b}, D. Salek¹¹⁸, P.H. Sales De Bruin¹⁷⁰, D. Salihagic¹¹³, A. Salnikov¹⁵⁰, J. Salt¹⁷², D. Salvatore^{40b,40a}, F. Salvatore¹⁵³, A. Salvucci^{64a,64b,64c}, A. Salzburger³⁵, D. Sammel⁵³, D. Sampsonidis¹⁵⁹, D. Sampsonidou¹⁵⁹, J. Sánchez¹⁷², A. Sanchez Pineda^{67a,67c}, H. Sandaker¹³⁰, C.O. Sander⁴⁶, M. Sandhoff¹⁸⁰, C. Sandoval²², D.P.C. Sankey¹⁴⁰, M. Sannino^{56b,56a}, Y. Sano¹¹⁵, A. Sansoni⁵², C. Santoni³⁷, H. Santos^{135a}, I. Santoyo Castillo¹⁵³, A. Saponov⁸⁰, J.G. Saraiva^{135a,135d}, O. Sasaki⁸¹, K. Sato¹⁶⁶, E. Sauvan⁵, P. Savard^{164,au}, N. Savic¹¹³, R. Sawada¹⁶⁰, C. Sawyer¹⁴⁰, L. Sawyer^{93,ak}, C. Sbarra^{23b}, A. Sbrizzi^{23b,23a}, T. Scanlon⁹², D.A. Scannicchio¹⁶⁹, J. Schaarschmidt¹⁴⁵, P. Schacht¹¹³, B.M. Schachtner¹¹², D. Schaefer³⁶, L. Schaefer¹³², J. Schaeffer⁹⁷, S. Schaepe³⁵, U. Schäfer⁹⁷, A.C. Schaffer¹²⁸, D. Schaile¹¹², R.D. Schamberger¹⁵², V.A. Schegelsky¹³³, D. Scheirich¹³⁸, F. Schenck¹⁹, M. Schernau¹⁶⁹, C. Schiavi^{56b,56a}, S. Schier¹⁴³, L.K. Schildgen²⁴, Z.M. Schillaci²⁶, C. Schillo⁵³, E.J. Schioppa³⁵, M. Schioppa^{40b,40a}, K.E. Schleicher⁵³, S. Schlenker³⁵, K.R. Schmidt-Sommerfeld¹¹³, K. Schmieden³⁵, C. Schmitt⁹⁷, S. Schmitt⁴⁶, S. Schmitz⁹⁷, U. Schnoor⁵³, L. Schoeffel¹⁴², A. Schoening^{62b}, E. Schopf²⁴, M. Schott⁹⁷, J.F.P. Schouwenberg¹¹⁷, J. Schovancova³⁵, S. Schramm⁵⁵, N. Schuh⁹⁷, A. Schulte⁹⁷, H-C. Schultz-Coulon^{62a}, M. Schumacher⁵³, B.A. Schumm¹⁴³, Ph. Schune¹⁴², A. Schwartzman¹⁵⁰, T.A. Schwarz¹⁰³, H. Schweiger⁹⁸, Ph. Schwemling¹⁴², R. Schwienhorst¹⁰⁴, A. Sciandra²⁴, G. Sciolla²⁶, M. Scornajenghi^{40b,40a}, F. Scuri^{72a}, F. Scutti¹⁰², L.M. Scyboz¹¹³, J. Searcy¹⁰³, P. Seema²⁴, S.C. Seidel¹¹⁶, A. Seiden¹⁴³, J.M. Seixas^{141a}, G. Sekhniaidze^{70a}, K. Sekhon¹⁰³, S.J. Sekula⁴³, N. Semprini-Cesari^{23b,23a}, S. Senkin³⁷, C. Serfon¹³⁰, L. Serin¹²⁸, L. Serkin^{67a,67b}, M. Sessa^{75a,75b}, H. Severini¹²⁴, F. Sforza¹⁶⁷, A. Sfyrlla⁵⁵, E. Shabalina⁵⁴, J.D. Shahinian¹⁴³, N.W. Shaikh^{45a,45b}, L.Y. Shan^{15a}, R. Shang¹⁷¹, J.T. Shank²⁵, M. Shapiro¹⁸, A.S. Sharma¹, P.B. Shatalov¹⁰⁹, K. Shaw^{67a,67b}, S.M. Shaw⁹⁸, A. Shcherbakova^{45a,45b}, C.Y. Shehu¹⁵³, Y. Shen¹²⁴, N. Sherafati³³, A.D. Sherman²⁵, P. Sherwood⁹², L. Shi^{155,aq}, S. Shimizu⁸², C.O. Shimmin¹⁸¹, M. Shimojima¹¹⁴, I.P.J. Shipsey¹³¹, S. Shirabe⁸⁵, M. Shiyakova^{80,ad}, J. Shlomi¹⁷⁸, A. Shmeleva¹⁰⁸, D. Shoaleh Saadi¹⁰⁷, M.J. Shochet³⁶, S. Shojaii¹⁰², D.R. Shope¹²⁴, S. Shrestha¹²², E. Shulga¹¹⁰,

P. Sicho ¹³⁶, A.M. Sickles ¹⁷¹, P.E. Sidebo ¹⁵¹, E. Sideras Haddad ^{32c}, O. Sidiropoulou ¹⁷⁵, A. Sidoti ^{23b,23a}, F. Siegert ⁴⁸, Dj. Sijacki ¹⁶, J. Silva ^{135a,135d}, M. Silva Jr. ¹⁷⁹, S.B. Silverstein ^{45a}, L. Simic ⁸⁰, S. Simion ¹²⁸, E. Simioni ⁹⁷, B. Simmons ⁹², M. Simon ⁹⁷, P. Sinervo ¹⁶⁴, N.B. Sinev ¹²⁷, M. Sioli ^{23b,23a}, G. Siragusa ¹⁷⁵, I. Siral ¹⁰³, S.Yu. Sivoklokov ¹¹¹, J. Sjölin ^{45a,45b}, M.B. Skinner ⁸⁷, P. Skubic ¹²⁴, M. Slater ²¹, T. Slavicek ¹³⁷, M. Slawinska ⁴², K. Sliwa ¹⁶⁷, R. Slovak ¹³⁸, V. Smakhtin ¹⁷⁸, B.H. Smart ⁵, J. Smiesko ^{28a}, N. Smirnov ¹¹⁰, S.Yu. Smirnov ¹¹⁰, Y. Smirnov ¹¹⁰, L.N. Smirnova ^{111,s}, O. Smirnova ⁹⁵, J.W. Smith ⁵⁴, M.N.K. Smith ³⁸, R.W. Smith ³⁸, M. Smizanska ⁸⁷, K. Smolek ¹³⁷, A.A. Snesarev ¹⁰⁸, I.M. Snyder ¹²⁷, S. Snyder ²⁹, R. Sobie ^{174,af}, F. Socher ⁴⁸, A.M. Soffa ¹⁶⁹, A. Soffer ¹⁵⁸, A. Søgaard ⁵⁰, D.A. Soh ¹⁵⁵, G. Sokhrannyi ⁸⁹, C.A. Solans Sanchez ³⁵, M. Solar ¹³⁷, E.Yu. Soldatov ¹¹⁰, U. Soldevila ¹⁷², A.A. Solodkov ¹³⁹, A. Soloshenko ⁸⁰, O.V. Solovyanov ¹³⁹, V. Solovyev ¹³³, P. Sommer ¹⁴⁶, H. Son ¹⁶⁷, W. Song ¹⁴⁰, A. Sopczak ¹³⁷, F. Sopkova ^{28b}, D. Sosa ^{62b}, C.L. Sotiropoulou ^{72a,72b}, S. Sottocornola ^{71a,71b}, R. Soualah ^{67a,67c}, A.M. Soukharev ^{120b,120a}, D. South ⁴⁶, B.C. Sowden ⁹¹, S. Spagnolo ^{68a,68b}, M. Spalla ¹¹³, M. Spangenberg ¹⁷⁶, F. Spanò ⁹¹, D. Sperlich ¹⁹, F. Spettel ¹¹³, T.M. Spieker ^{62a}, R. Spighi ^{23b}, G. Spigo ³⁵, L.A. Spiller ¹⁰², M. Spousta ¹³⁸, R.D. St. Denis ^{58,*}, A. Stabile ^{69a,69b}, R. Stamen ^{62a}, S. Stamm ¹⁹, E. Stanecka ⁴², R.W. Stanek ⁶, C. Stanescu ^{75a}, M.M. Stanitzki ⁴⁶, B.S. Stapf ¹¹⁸, S. Stapnes ¹³⁰, E.A. Starchenko ¹³⁹, G.H. Stark ³⁶, J. Stark ⁵⁹, S.H. Stark ³⁹, P. Staroba ¹³⁶, P. Starovoitov ^{62a}, S. Stärz ³⁵, R. Staszewski ⁴², M. Stegler ⁴⁶, P. Steinberg ²⁹, B. Stelzer ¹⁴⁹, H.J. Stelzer ³⁵, O. Stelzer-Chilton ^{165a}, H. Stenzel ⁵⁷, T.J. Stevenson ⁹⁰, G.A. Stewart ⁵⁸, M.C. Stockton ¹²⁷, G. Stoicea ^{27b}, P. Stolte ⁵⁴, S. Stonjek ¹¹³, A. Straessner ⁴⁸, M.E. Stramaglia ²⁰, J. Strandberg ¹⁵¹, S. Strandberg ^{45a,45b}, M. Strauss ¹²⁴, P. Strizenec ^{28b}, R. Ströhmer ¹⁷⁵, D.M. Strom ¹²⁷, R. Stroynowski ⁴³, A. Strubig ⁵⁰, S.A. Stucci ²⁹, B. Stugu ¹⁷, N.A. Styles ⁴⁶, D. Su ¹⁵⁰, J. Su ¹³⁴, S. Suchek ^{62a}, Y. Sugaya ¹²⁹, M. Suk ¹³⁷, V.V. Sulin ¹⁰⁸, D.M.S. Sultan ⁵⁵, S. Sultansoy ^{4c}, T. Sumida ⁸³, S. Sun ¹⁰³, X. Sun ³, K. Suruliz ¹⁵³, C.J.E. Suster ¹⁵⁴, M.R. Sutton ¹⁵³, S. Suzuki ⁸¹, M. Svatos ¹³⁶, M. Swiatlowski ³⁶, S.P. Swift ², A. Sydorenko ⁹⁷, I. Sykora ^{28a}, T. Sykora ¹³⁸, D. Ta ⁹⁷, K. Tackmann ⁴⁶, J. Taenzer ¹⁵⁸, A. Taffard ¹⁶⁹, R. Tafirout ^{165a}, E. Tahirovic ⁹⁰, N. Taiblum ¹⁵⁸, H. Takai ²⁹, R. Takashima ⁸⁴, E.H. Takasugi ¹¹³, K. Takeda ⁸², T. Takeshita ¹⁴⁷, Y. Takubo ⁸¹, M. Talby ⁹⁹, A.A. Talyshev ^{120b,120a}, J. Tanaka ¹⁶⁰, M. Tanaka ¹⁶², R. Tanaka ¹²⁸, R. Tanioka ⁸², B.B. Tannenwald ¹²², S. Tapia Araya ^{144b}, S. Tapprogge ⁹⁷, A. Tarek Abouelfadl Mohamed ⁹⁴, S. Tarem ¹⁵⁷, G. Tarna ^{27b,d}, G.F. Tartarelli ^{69a}, P. Tas ¹³⁸, M. Tasevsky ¹³⁶, T. Tashiro ⁸³, E. Tassi ^{40b,40a}, A. Tavares Delgado ^{135a,135b}, Y. Tayalati ^{34e}, A.C. Taylor ¹¹⁶, A.J. Taylor ⁵⁰, G.N. Taylor ¹⁰², P.T.E. Taylor ¹⁰², W. Taylor ^{165b}, P. Teixeira-Dias ⁹¹, D. Temple ¹⁴⁹, H. Ten Kate ³⁵, P.K. Teng ¹⁵⁵, J.J. Teoh ¹²⁹, F. Tepel ¹⁸⁰, S. Terada ⁸¹, K. Terashi ¹⁶⁰, J. Terron ⁹⁶, S. Terzo ¹⁴, M. Testa ⁵², R.J. Teuscher ^{164,af}, S.J. Thais ¹⁸¹, T. Theveneaux-Pelzer ⁴⁶, F. Thiele ³⁹, J.P. Thomas ²¹, A.S. Thompson ⁵⁸, P.D. Thompson ²¹, L.A. Thomsen ¹⁸¹, E. Thomson ¹³², Y. Tian ³⁸, R.E. Ticse Torres ⁵⁴, V.O. Tikhomirov ^{108,an}, Yu.A. Tikhonov ^{120b,120a}, S. Timoshenko ¹¹⁰, P. Tipton ¹⁸¹, S. Tisserant ⁹⁹, K. Todome ¹⁶², S. Todorova-Nova ⁵, S. Todt ⁴⁸, J. Tojo ⁸⁵, S. Tokár ^{28a}, K. Tokushuku ⁸¹, E. Tolley ¹²², M. Tomoto ¹¹⁵, L. Tompkins ^{150,o}, K. Toms ¹¹⁶, B. Tong ⁶⁰, P. Tornambe ⁵³, E. Torrence ¹²⁷, H. Torres ⁴⁸, E. Torró Pastor ¹⁴⁵, J. Toth ^{99,ae}, F. Touchard ⁹⁹, D.R. Tovey ¹⁴⁶, C.J. Treado ¹²¹, T. Trefzger ¹⁷⁵, F. Tresoldi ¹⁵³, A. Tricoli ²⁹, I.M. Trigger ^{165a}, S. Trincaz-Duvoid ⁹⁴, M.F. Tripiana ¹⁴, W. Trischuk ¹⁶⁴, B. Trocme ⁵⁹, A. Trofymov ⁴⁶, C. Troncon ^{69a}, M. Trovatelli ¹⁷⁴, L. Truong ^{32b}, M. Trzebinski ⁴², A. Trzupek ⁴², K.W. Tsang ^{64a}, J.C.-L. Tseng ¹³¹, P.V. Tsiarehka ¹⁰⁵, N. Tsirintanis ⁹, S. Tsiskaridze ¹⁴, V. Tsiskaridze ¹⁵², E.G. Tskhadadze ^{156a}, I.I. Tsukerman ¹⁰⁹, V. Tsulaia ¹⁸, S. Tsuno ⁸¹, D. Tsybychev ¹⁵², Y. Tu ^{64b}, A. Tudorache ^{27b}, V. Tudorache ^{27b}, T.T. Tulbure ^{27a}, A.N. Tuna ⁶⁰, S. Turchikhin ⁸⁰, D. Turgeman ¹⁷⁸, I. Turk Cakir ^{4b,v}, R. Turra ^{69a}, P.M. Tuts ³⁸, G. Uccielli ^{23b,23a}, I. Ueda ⁸¹, M. Ughetto ^{45a,45b}, F. Ukegawa ¹⁶⁶, G. Unal ³⁵, A. Undrus ²⁹, G. Unel ¹⁶⁹, F.C. Ungaro ¹⁰², Y. Unno ⁸¹, K. Uno ¹⁶⁰, J. Urban ^{28b}, P. Urquijo ¹⁰², P. Urrejola ⁹⁷, G. Usai ⁸, J. Usui ⁸¹, L. Vacavant ⁹⁹, V. Vacek ¹³⁷, B. Vachon ¹⁰¹, K.O.H. Vadla ¹³⁰, A. Vaidya ⁹², C. Valderanis ¹¹², E. Valdes Santurio ^{45a,45b}, M. Valente ⁵⁵, S. Valentini ^{23b,23a}, A. Valero ¹⁷², L. Valéry ¹⁴, A. Vallier ⁵, J.A. Valls Ferrer ¹⁷², W. Van Den Wollenberg ¹¹⁸, H. van der Graaf ¹¹⁸, P. van Gemmeren ⁶, J. Van Nieuwkoop ¹⁴⁹, I. van Vulpen ¹¹⁸, M.C. van Woerden ¹¹⁸, M. Vanadia ^{74a,74b}, W. Vandelli ³⁵, A. Vaniachine ¹⁶³, P. Vankov ¹¹⁸, R. Vari ^{73a}, E.W. Varnes ⁷, C. Varni ^{56b,56a}, T. Varol ⁴³, D. Varouchas ¹²⁸, A. Vartapetian ⁸, K.E. Varvell ¹⁵⁴, G.A. Vasquez ^{144b}, J.G. Vasquez ¹⁸¹, F. Vazeille ³⁷, D. Vazquez Furelos ¹⁴, T. Vazquez Schroeder ¹⁰¹, J. Veatch ⁵⁴, L.M. Veloce ¹⁶⁴, F. Veloso ^{135a,135c}, S. Veneziano ^{73a}, A. Ventura ^{68a,68b}, M. Venturi ¹⁷⁴, N. Venturi ³⁵, V. Vercesi ^{71a}, M. Verducci ^{75a,75b}, W. Verkerke ¹¹⁸,

A.T. Vermeulen¹¹⁸, J.C. Vermeulen¹¹⁸, M.C. Vetterli^{149,au}, N. Viaux Maira^{144b}, O. Viazlo⁹⁵,
 I. Vichou^{171,*}, T. Vickey¹⁴⁶, O.E. Vickey Boeriu¹⁴⁶, G.H.A. Viehhauser¹³¹, S. Viel¹⁸, L. Vignani¹³¹,
 M. Villa^{23b,23a}, M. Villaplana Perez^{69a,69b}, E. Vilucchi⁵², M.G. Vincter³³, V.B. Vinogradov⁸⁰,
 A. Vishwakarma⁴⁶, C. Vittori^{23b,23a}, I. Vivarelli¹⁵³, S. Vlachos¹⁰, M. Vogel¹⁸⁰, P. Vokac¹³⁷, G. Volpi¹⁴,
 S.E. von Buddenbrock^{32c}, E. von Toerne²⁴, V. Vorobel¹³⁸, K. Vorobev¹¹⁰, M. Vos¹⁷², J.H. Vosseveld⁸⁸,
 N. Vranjes¹⁶, M. Vranjes Milosavljevic¹⁶, V. Vrba¹³⁷, M. Vreeswijk¹¹⁸, T. Šfiligoj⁸⁹, R. Vuillermet³⁵,
 I. Vukotic³⁶, T. Ženiš^{28a}, L. Živković¹⁶, P. Wagner²⁴, W. Wagner¹⁸⁰, J. Wagner-Kuhr¹¹², H. Wahlberg⁸⁶,
 S. Wahrenmund⁴⁸, K. Wakamiya⁸², J. Walder⁸⁷, R. Walker¹¹², W. Walkowiak¹⁴⁸, V. Wallangen^{45a,45b},
 A.M. Wang⁶⁰, C. Wang^{61b,d}, F. Wang¹⁷⁹, H. Wang¹⁸, H. Wang³, J. Wang¹⁵⁴, J. Wang^{62b}, Q. Wang¹²⁴,
 R.-J. Wang⁹⁴, R. Wang⁶, S.M. Wang¹⁵⁵, T. Wang³⁸, W. Wang^{15b}, W. Wang^{61a,ag}, Z. Wang^{61c},
 C. Wanotayaroj⁴⁶, A. Warburton¹⁰¹, C.P. Ward³¹, D.R. Wardrope⁹², A. Washbrook⁵⁰, P.M. Watkins²¹,
 A.T. Watson²¹, M.F. Watson²¹, G. Watts¹⁴⁵, S. Watts⁹⁸, B.M. Waugh⁹², A.F. Webb¹¹, S. Webb⁹⁷,
 M.S. Weber²⁰, S.A. Weber³³, S.M. Weber^{62a}, J.S. Webster⁶, A.R. Weidberg¹³¹, B. Weinert⁶⁶,
 J. Weingarten⁵⁴, M. Weirich⁹⁷, C. Weiser⁵³, P.S. Wells³⁵, T. Wenaus²⁹, T. Wengler³⁵, S. Wenig³⁵,
 N. Wermes²⁴, M.D. Werner⁷⁹, P. Werner³⁵, M. Wessels^{62a}, T.D. Weston²⁰, K. Whalen¹²⁷,
 N.L. Whallon¹⁴⁵, A.M. Wharton⁸⁷, A.S. White¹⁰³, A. White⁸, M.J. White¹, R. White^{144b}, D. Whiteson¹⁶⁹,
 B.W. Whitmore⁸⁷, F.J. Wickens¹⁴⁰, W. Wiedenmann¹⁷⁹, M. Wielers¹⁴⁰, C. Wiglesworth³⁹,
 L.A.M. Wiik-Fuchs⁵³, A. Wildauer¹¹³, F. Wilk⁹⁸, H.G. Wilkens³⁵, H.H. Williams¹³², S. Williams³¹,
 C. Willis¹⁰⁴, S. Willocq¹⁰⁰, J.A. Wilson²¹, I. Wingerter-Seez⁵, E. Winkels¹⁵³, F. Winklmeier¹²⁷,
 O.J. Winston¹⁵³, B.T. Winter²⁴, M. Wittgen¹⁵⁰, M. Wobisch⁹³, A. Wolf⁹⁷, T.M.H. Wolf¹¹⁸, R. Wolff⁹⁹,
 M.W. Wolter⁴², H. Wolters^{135a,135c}, V.W.S. Wong¹⁷³, N.L. Woods¹⁴³, S.D. Worm²¹, B.K. Wosiek⁴²,
 K.W. Woźniak⁴², M. Wu³⁶, S.L. Wu¹⁷⁹, X. Wu⁵⁵, Y. Wu^{61a}, T.R. Wyatt⁹⁸, B.M. Wynne⁵⁰, S. Xella³⁹,
 Z. Xi¹⁰³, L. Xia^{15c}, D. Xu^{15a}, H. Xu^{61a}, L. Xu²⁹, T. Xu¹⁴², W. Xu¹⁰³, B. Yabsley¹⁵⁴, S. Yacoub^{32a},
 K. Yajima¹²⁹, D.P. Yallup⁹², D. Yamaguchi¹⁶², Y. Yamaguchi¹⁶², A. Yamamoto⁸¹, T. Yamanaka¹⁶⁰,
 F. Yamane⁸², M. Yamatani¹⁶⁰, T. Yamazaki¹⁶⁰, Y. Yamazaki⁸², Z. Yan²⁵, H. Yang^{61c,61d}, H. Yang¹⁸,
 S. Yang⁷⁸, Y. Yang¹⁶⁰, Y. Yang¹⁵⁵, Z. Yang¹⁷, W.-M. Yao¹⁸, Y.C. Yap⁴⁶, Y. Yasu⁸¹, E. Yatsenko⁵,
 K.H. Yau Wong²⁴, J. Ye⁴³, S. Ye²⁹, I. Yeletsikh⁸⁰, E. Yigitbasi²⁵, E. Yildirim⁹⁷, K. Yorita¹⁷⁷,
 K. Yoshihara¹³², C.J.S. Young³⁵, C. Young¹⁵⁰, J. Yu⁸, J. Yu⁷⁹, S.P.Y. Yuen²⁴, I. Yusuff^{31,aw}, B. Zabinski⁴²,
 G. Zacharis¹⁰, R. Zaidan¹⁴, A.M. Zaitsev^{139,am}, N. Zakharchuk⁴⁶, J. Zalieckas¹⁷, S. Zambito⁶⁰, D. Zanzi³⁵,
 C. Zeitnitz¹⁸⁰, G. Zemaityte¹³¹, J.C. Zeng¹⁷¹, Q. Zeng¹⁵⁰, O. Zenin¹³⁹, D. Zerwas¹²⁸, D. Zhang¹⁰³,
 D. Zhang^{61b}, F. Zhang¹⁷⁹, G. Zhang^{61a,ag}, H. Zhang¹²⁸, J. Zhang⁶, L. Zhang⁵³, L. Zhang^{61a}, M. Zhang¹⁷¹,
 P. Zhang^{15b}, R. Zhang^{61a,d}, R. Zhang²⁴, X. Zhang^{61b}, Y. Zhang^{15d}, Z. Zhang¹²⁸, X. Zhao⁴³, Y. Zhao^{61b,aj},
 Z. Zhao^{61a}, A. Zhemchugov⁸⁰, B. Zhou¹⁰³, C. Zhou¹⁷⁹, L. Zhou⁴³, M. Zhou^{15d}, M. Zhou¹⁵², N. Zhou^{61c},
 Y. Zhou⁷, C.G. Zhu^{61b}, H. Zhu^{15a}, J. Zhu¹⁰³, Y. Zhu^{61a}, X. Zhuang^{15a}, K. Zhukov¹⁰⁸, V. Zhulanov^{120b,120a},
 A. Zibell¹⁷⁵, D. Zieminska⁶⁶, N.I. Zimine⁸⁰, S. Zimmermann⁵³, Z. Zinonos¹¹³, M. Zinser⁹⁷,
 M. Ziolkowski¹⁴⁸, G. Zoernig¹⁷⁹, A. Zoccoli^{23b,23a}, T.G. Zorbas¹⁴⁶, R. Zou³⁶, M. zur Nedden¹⁹,
 L. Zwalinski³⁵

¹ Department of Physics, University of Adelaide, Adelaide, Australia

² Physics Department, SUNY Albany, Albany, NY, United States of America

³ Department of Physics, University of Alberta, Edmonton, AB, Canada

⁴ (a) Department of Physics, Ankara University, Ankara; (b) Istanbul Aydin University, Istanbul; (c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey

⁵ LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy-le-Vieux, France

⁶ High Energy Physics Division, Argonne National Laboratory, Argonne, IL, United States of America

⁷ Department of Physics, University of Arizona, Tucson, AZ, United States of America

⁸ Department of Physics, University of Texas at Arlington, Arlington, TX, United States of America

⁹ Physics Department, National and Kapodistrian University of Athens, Athens, Greece

¹⁰ Physics Department, National Technical University of Athens, Zografou, Greece

¹¹ Department of Physics, University of Texas at Austin, Austin, TX, United States of America

¹² (a) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; (b) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; (c) Department of Physics, Bogazici University, Istanbul; (d) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey

¹³ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

¹⁴ Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain

¹⁵ (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Department of Physics, Nanjing University, Nanjing; (c) Physics Department, Tsinghua University, Beijing;

(d) University of Chinese Academy of Science (UCAS), Beijing, China

¹⁶ Institute of Physics, University of Belgrade, Belgrade, Serbia

¹⁷ Department for Physics and Technology, University of Bergen, Bergen, Norway

¹⁸ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, United States of America

¹⁹ Institut für Physik, Humboldt Universität, Berlin, Germany

²⁰ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland

- ²¹ School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- ²² Centro de Investigaciones, Universidad Antonio Nariño, Bogota, Colombia
- ²³ ^(a) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna; ^(b) INFN Sezione di Bologna, Italy
- ²⁴ Physikalisches Institut, Universität Bonn, Bonn, Germany
- ²⁵ Department of Physics, Boston University, Boston, MA, United States of America
- ²⁶ Department of Physics, Brandeis University, Waltham, MA, United States of America
- ²⁷ ^(a) Transilvania University of Brasov, Brasov; ^(b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; ^(c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; ^(d) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; ^(e) University Politehnica Bucharest, Bucharest; ^(f) West University in Timisoara, Timisoara, Romania
- ²⁸ ^(a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; ^(b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
- ²⁹ Physics Department, Brookhaven National Laboratory, Upton, NY, United States of America
- ³⁰ Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
- ³¹ Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
- ³² ^(a) Department of Physics, University of Cape Town, Cape Town; ^(b) Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg; ^(c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa
- ³³ Department of Physics, Carleton University, Ottawa, ON, Canada
- ³⁴ ^(a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies – Université Hassan II, Casablanca; ^(b) Centre National de l’Energie des Sciences Techniques Nucleaires (CNESTEN), Rabat; ^(c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA, Marrakech; ^(d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; ^(e) Faculté des Sciences, Université Mohammed V, Rabat, Morocco
- ³⁵ CERN, Geneva, Switzerland
- ³⁶ Enrico Fermi Institute, University of Chicago, Chicago, IL, United States of America
- ³⁷ LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
- ³⁸ Nevis Laboratory, Columbia University, Irvington, NY, United States of America
- ³⁹ Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- ⁴⁰ ^(a) Dipartimento di Fisica, Università della Calabria, Rende; ^(b) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
- ⁴¹ ^(a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; ^(b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
- ⁴² Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
- ⁴³ Physics Department, Southern Methodist University, Dallas, TX, United States of America
- ⁴⁴ Physics Department, University of Texas at Dallas, Richardson, TX, United States of America
- ⁴⁵ ^(a) Department of Physics, Stockholm University; ^(b) Oskar Klein Centre, Stockholm, Sweden
- ⁴⁶ Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany
- ⁴⁷ Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- ⁴⁸ Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
- ⁴⁹ Department of Physics, Duke University, Durham, NC, United States of America
- ⁵⁰ SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- ⁵¹ Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
- ⁵² INFN e Laboratori Nazionali di Frascati, Frascati, Italy
- ⁵³ Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
- ⁵⁴ II. Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
- ⁵⁵ Département de Physique Nucléaire et Corpusculaire, Université de Genève, Geneva, Switzerland
- ⁵⁶ ^(a) Dipartimento di Fisica, Università di Genova, Genova; ^(b) INFN Sezione di Genova, Italy
- ⁵⁷ II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- ⁵⁸ SUPA – School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- ⁵⁹ LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France
- ⁶⁰ Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States of America
- ⁶¹ ^(a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; ^(b) School of Physics, Shandong University, Shandong; ^(c) School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai; ^(d) Tsung-Dao Lee Institute, Shanghai, China
- ⁶² ^(a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ⁶³ Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- ⁶⁴ ^(a) Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; ^(b) Department of Physics, University of Hong Kong, Hong Kong; ^(c) Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
- ⁶⁵ Department of Physics, National Tsing Hua University, Hsinchu, Taiwan
- ⁶⁶ Department of Physics, Indiana University, Bloomington, IN, United States of America
- ⁶⁷ ^(a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; ^(b) ICTP, Trieste; ^(c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
- ⁶⁸ ^(a) INFN Sezione di Lecce; ^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- ⁶⁹ ^(a) INFN Sezione di Milano; ^(b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- ⁷⁰ ^(a) INFN Sezione di Napoli; ^(b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy
- ⁷¹ ^(a) INFN Sezione di Pavia; ^(b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- ⁷² ^(a) INFN Sezione di Pisa; ^(b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- ⁷³ ^(a) INFN Sezione di Roma; ^(b) Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
- ⁷⁴ ^(a) INFN Sezione di Roma Tor Vergata; ^(b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
- ⁷⁵ ^(a) INFN Sezione di Roma Tre; ^(b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
- ⁷⁶ ^(a) INFN-TIFPA; ^(b) Università degli Studi di Trento, Trento, Italy
- ⁷⁷ Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- ⁷⁸ University of Iowa, Iowa City, IA, United States of America
- ⁷⁹ Department of Physics and Astronomy, Iowa State University, Ames, IA, United States of America
- ⁸⁰ Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
- ⁸¹ KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- ⁸² Graduate School of Science, Kobe University, Kobe, Japan
- ⁸³ Faculty of Science, Kyoto University, Kyoto, Japan
- ⁸⁴ Kyoto University of Education, Kyoto, Japan
- ⁸⁵ Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan
- ⁸⁶ Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- ⁸⁷ Physics Department, Lancaster University, Lancaster, United Kingdom
- ⁸⁸ Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- ⁸⁹ Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia
- ⁹⁰ School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom

- ⁹¹ Department of Physics, Royal Holloway University of London, Egham, United Kingdom
- ⁹² Department of Physics and Astronomy, University College London, London, United Kingdom
- ⁹³ Louisiana Tech University, Ruston, LA, United States of America
- ⁹⁴ LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France
- ⁹⁵ Fysiska institutionen, Lunds universitet, Lund, Sweden
- ⁹⁶ Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain
- ⁹⁷ Institut für Physik, Universität Mainz, Mainz, Germany
- ⁹⁸ School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- ⁹⁹ CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- ¹⁰⁰ Department of Physics, University of Massachusetts, Amherst, MA, United States of America
- ¹⁰¹ Department of Physics, McGill University, Montreal, QC, Canada
- ¹⁰² School of Physics, University of Melbourne, Victoria, Australia
- ¹⁰³ Department of Physics, University of Michigan, Ann Arbor, MI, United States of America
- ¹⁰⁴ Department of Physics and Astronomy, Michigan State University, East Lansing, MI, United States of America
- ¹⁰⁵ B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus
- ¹⁰⁶ Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Belarus
- ¹⁰⁷ Group of Particle Physics, University of Montreal, Montreal, QC, Canada
- ¹⁰⁸ P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia
- ¹⁰⁹ Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
- ¹¹⁰ National Research Nuclear University MEPhI, Moscow, Russia
- ¹¹¹ D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- ¹¹² Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- ¹¹³ Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- ¹¹⁴ Nagasaki Institute of Applied Science, Nagasaki, Japan
- ¹¹⁵ Graduate School of Science and Kobayashi–Maskawa Institute, Nagoya University, Nagoya, Japan
- ¹¹⁶ Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States of America
- ¹¹⁷ Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- ¹¹⁸ Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- ¹¹⁹ Department of Physics, Northern Illinois University, DeKalb, IL, United States of America
- ¹²⁰ ^(a) Budker Institute of Nuclear Physics, SB RAS, Novosibirsk; ^(b) Novosibirsk State University Novosibirsk, Russia
- ¹²¹ Department of Physics, New York University, New York, NY, United States of America
- ¹²² Ohio State University, Columbus, OH, United States of America
- ¹²³ Faculty of Science, Okayama University, Okayama, Japan
- ¹²⁴ Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States of America
- ¹²⁵ Department of Physics, Oklahoma State University, Stillwater, OK, United States of America
- ¹²⁶ Palacký University, RCPTM, Olomouc, Czech Republic
- ¹²⁷ Center for High Energy Physics, University of Oregon, Eugene, OR, United States of America
- ¹²⁸ LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
- ¹²⁹ Graduate School of Science, Osaka University, Osaka, Japan
- ¹³⁰ Department of Physics, University of Oslo, Oslo, Norway
- ¹³¹ Department of Physics, Oxford University, Oxford, United Kingdom
- ¹³² Department of Physics, University of Pennsylvania, Philadelphia, PA, United States of America
- ¹³³ Konstantinov Nuclear Physics Institute of National Research Centre “Kurchatov Institute”, PNPI, St. Petersburg, Russia
- ¹³⁴ Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States of America
- ¹³⁵ ^(a) Laboratório de Instrumentação e Física Experimental de Partículas – LIP; ^(b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; ^(c) Departamento de Física, Universidade de Coimbra, Coimbra; ^(d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; ^(e) Departamento de Física, Universidade do Minho, Braga; ^(f) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); ^(g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
- ¹³⁶ Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
- ¹³⁷ Czech Technical University in Prague, Prague, Czech Republic
- ¹³⁸ Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic
- ¹³⁹ State Research Center Institute for High Energy Physics, NRC KI, Protvino, Russia
- ¹⁴⁰ Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- ¹⁴¹ ^(a) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; ^(b) Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; ^(c) Universidade Federal de São João del Rei (UFSJ), São João del Rei; ^(d) Instituto de Física, Universidade de São Paulo, São Paulo, Brazil
- ¹⁴² DRF/IRFU, CEA Saclay, Gif-sur-Yvette, France
- ¹⁴³ Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, United States of America
- ¹⁴⁴ ^(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; ^(b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- ¹⁴⁵ Department of Physics, University of Washington, Seattle, WA, United States of America
- ¹⁴⁶ Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
- ¹⁴⁷ Department of Physics, Shinshu University, Nagano, Japan
- ¹⁴⁸ Department Physik, Universität Siegen, Siegen, Germany
- ¹⁴⁹ Department of Physics, Simon Fraser University, Burnaby, BC, Canada
- ¹⁵⁰ SLAC National Accelerator Laboratory, Stanford, CA, United States of America
- ¹⁵¹ Physics Department, Royal Institute of Technology, Stockholm, Sweden
- ¹⁵² Departments of Physics and Astronomy, Stony Brook University, Stony Brook, NY, United States of America
- ¹⁵³ Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
- ¹⁵⁴ School of Physics, University of Sydney, Sydney, Australia
- ¹⁵⁵ Institute of Physics, Academia Sinica, Taipei, Taiwan
- ¹⁵⁶ ^(a) E. Andronikashvili Institute of Physics, Iv. Javakishvili Tbilisi State University, Tbilisi; ^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
- ¹⁵⁷ Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel
- ¹⁵⁸ Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
- ¹⁵⁹ Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
- ¹⁶⁰ International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan
- ¹⁶¹ Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
- ¹⁶² Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
- ¹⁶³ Tomsk State University, Tomsk, Russia
- ¹⁶⁴ Department of Physics, University of Toronto, Toronto, ON, Canada
- ¹⁶⁵ ^(a) TRIUMF, Vancouver, BC; ^(b) Department of Physics and Astronomy, York University, Toronto, ON, Canada
- ¹⁶⁶ Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan

- ¹⁶⁷ Department of Physics and Astronomy, Tufts University, Medford, MA, United States of America
¹⁶⁸ Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
¹⁶⁹ Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States of America
¹⁷⁰ Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
¹⁷¹ Department of Physics, University of Illinois, Urbana, IL, United States of America
¹⁷² Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia – CSIC, Valencia, Spain
¹⁷³ Department of Physics, University of British Columbia, Vancouver, BC, Canada
¹⁷⁴ Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
¹⁷⁵ Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
¹⁷⁶ Department of Physics, University of Warwick, Coventry, United Kingdom
¹⁷⁷ Waseda University, Tokyo, Japan
¹⁷⁸ Department of Particle Physics, Weizmann Institute of Science, Rehovot, Israel
¹⁷⁹ Department of Physics, University of Wisconsin, Madison, WI, United States of America
¹⁸⁰ Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany
¹⁸¹ Department of Physics, Yale University, New Haven, CT, United States of America
¹⁸² Yerevan Physics Institute, Yerevan, Armenia

- ^a Also at Borough of Manhattan Community College, City University of New York, New York City, United States of America.
^b Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town, South Africa.
^c Also at CERN, Geneva, Switzerland.
^d Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France.
^e Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain.
^f Also at Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain), Spain.
^g Also at Departement de Physique Nucléaire et Corpusculaire, Université de Genève, Geneva, Switzerland.
^h Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.
ⁱ Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY, United States of America.
^j Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom.
^k Also at Department of Physics, California State University, Fresno CA, United States of America.
^l Also at Department of Physics, California State University, Sacramento CA, United States of America.
^m Also at Department of Physics, King's College London, London, United Kingdom.
ⁿ Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.
^o Also at Department of Physics, Stanford University, Stanford CA, United States of America.
^p Also at Department of Physics, University of Fribourg, Fribourg, Switzerland.
^q Also at Department of Physics, University of Michigan, Ann Arbor MI, United States of America.
^r Also at Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy.
^s Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia.
^t Also at Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany.
^u Also at Georgian Technical University (GTU), Tbilisi, Georgia.
^v Also at Giresun University, Faculty of Engineering, Turkey.
^w Also at Graduate School of Science, Osaka University, Osaka, Japan.
^x Also at Hellenic Open University, Patras, Greece.
^y Also at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania.
^z Also at II. Physikalisches Institut, Georg-August-Universität, Göttingen, Germany.
^{aa} Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.
^{ab} Also at Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain.
^{ac} Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands.
^{ad} Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria.
^{ae} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
^{af} Also at Institute of Particle Physics (IPP), Canada.
^{ag} Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
^{ah} Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.
^{ai} Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.
^{aj} Also at LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France.
^{ak} Also at Louisiana Tech University, Ruston LA, United States of America.
^{al} Also at Manhattan College, New York NY, United States of America.
^{am} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.
^{an} Also at National Research Nuclear University MEPhI, Moscow, Russia.
^{ao} Also at Near East University, Nicosia, North Cyprus, Mersin 10, Turkey.
^{ap} Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan.
^{aq} Also at School of Physics, Sun Yat-sen University, Guangzhou, China.
^{ar} Also at The City College of New York, New York NY, United States of America.
^{as} Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing, China.
^{at} Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.
^{au} Also at TRIUMF, Vancouver, BC, Canada.
^{av} Also at Università di Napoli Parthenope, Napoli, Italy.
^{aw} Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.
* Deceased.