Measurement of the total cross section from elastic scattering in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

The ATLAS Collaboration ⋆

ARTICLE INFO

ABSTRACT

A measurement of the total pp cross section at the LHC at $\sqrt{s} = 8$ TeV is presented. An integrated luminosity of 500 μb^{-1} was accumulated in a special run with high-$\beta^\star$ beam optics to measure the differential elastic cross section as a function of the Mandelstam momentum transfer variable t. The measurement is performed with the ALFA sub-detector of ATLAS. Using a fit to the differential elastic cross section in the $-t$ range from 0.014 GeV^2 to 0.1 GeV^2 to extrapolate $t \to 0$, the total cross section, $\sigma_{\text{tot}}(pp \to X)$, is measured via the optical theorem to be

$$\sigma_{\text{tot}}(pp \to X) = 96.07 \pm 0.18 \, (\text{stat.}) \pm 0.85 \, (\text{exp.}) \pm 0.31 \, (\text{extr.}) \, \text{mb},$$

where the first error is statistical, the second accounts for all experimental systematic uncertainties and the last is related to uncertainties in the extrapolation $t \to 0$. In addition, the slope of the exponential function describing the elastic cross section at small t is determined to be $B = 19.74 \pm 0.05 \, (\text{stat.}) \pm 0.23 \, (\text{syst.})$ GeV^{-2}.

1. Introduction

The total cross section for proton–proton (pp) interactions characterizes a fundamental process of the strong interaction. Its energy evolution has been studied at each new range of centre-of-mass energies available. ATLAS has previously reported a measurement of the total cross section in pp collisions at $\sqrt{s} = 7$ TeV [1]. This paper details a measurement of the total cross section at $\sqrt{s} = 8$ TeV using data collected in 2012. The measurement methodology and analysis technique are very similar between the two measurements and the technical details are discussed thoroughly in Ref. [1].

Both measurements rely on the optical theorem:

$$\sigma_{\text{tot}} = 4\pi \, \text{Im}\,[f_{\text{el}}\,(t \to 0)] \tag{1}$$

which relates the total pp cross section σ_{tot} to the elastic-scattering amplitude extrapolated to the forward direction $f_{\text{el}}(t \to 0)$, with t being the four-momentum transfer squared. The total cross section can be extracted in different ways using the optical theorem. ATLAS uses the *luminosity-dependent* method which requires a measurement of the luminosity in order to normalize the elastic cross section. Here the measurement benefits from the high-precision luminosity measurement that ATLAS provides. With this method, σ_{tot} is given by the formula:

$$\sigma_{\text{tot}}^2 = \frac{16\pi\,(\hbar c)^2}{1 + \rho^2} \, \frac{\text{d}\sigma_{\text{el}}}{\text{d}t}\bigg|_{t \to 0}, \tag{2}$$

where ρ represents a small correction arising from the ratio of the real to the imaginary part of the elastic-scattering amplitude in the forward direction and is taken from global model extrapolations [2].

The first measurement of σ_{tot} at the LHC at 8 TeV was performed by the TOTEM Collaboration [3] using a *luminosity-independent* method and using data from the same LHC fill as ATLAS. At 7 TeV measurements of σ_{tot} were provided by TOTEM [4–6] and ATLAS [1]. In a recent publication a measurement in the Coulomb–nuclear interference region at very small t was also reported by TOTEM [7]. The inelastic cross section σ_{inel} can either be derived from the total and elastic cross section measurements as in Refs. [3–6,1] at 7 and 8 TeV, or be determined directly from the measurement of the inelastic rate without exploiting the optical theorem. These measurements of σ_{inel} were performed at 7 TeV by all LHC Collaborations [8–12] and recently also at 13 TeV by ATLAS [13].

2. Experimental setup

The ATLAS detector is described in detail elsewhere [14]. The elastic-scattering data were recorded with the ALFA sub-detector (Absolute Luminosity For ATLAS) [1]. It consists of Roman Pot (RP) tracking-detector stations placed at distances of 237 m (inner sta-

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tion) and 241 m (outer station) on either side of the ATLAS interaction point (IP). Each station houses two vertically moveable scintillating fibre detectors which are inserted in RPs and positioned close to the beam for data taking. Each detector consists of 10 modules of scintillating fibres with 64 fibres on both the front and back sides of a titanium support plate. The fibres are arranged orthogonally in a u–v-geometry at $\pm45°$ with respect to the y-axis.[1] The spatial resolution of the detectors is about 35 μm. Elastic scattering events are recorded in two independent arms of the spectrometer. Arm 1 consists of two upper detectors at the left side and two lower detectors at the right side, and arm 2 consists inversely of two lower detectors at the left and two upper detectors at the right side. Events with reconstructed tracks in all four detectors of an arm are referred to as "golden" events [1]. The detectors are supplemented with trigger counters consisting of plain scintillator tiles. The detector geometry is illustrated in Fig. 1 of Ref. [1]. All scintillation signals are detected by photomultipliers coupled to a compact assembly of front-end electronics including the MAROC chip [15,16] for signal amplification and discrimination. The entire experimental setup is depicted in Fig. 2 of Ref. [1].

3. Experimental method

3.1. Measurement principle

The data were recorded in a single run of the LHC with special beam optics [17,18] of $\beta^\star = 90$ m.[2] The same optics were used at 7 TeV [1] and result in a small beam divergence with parallel-to-point focusing in the vertical plane. The four-momentum transfer t is calculated from the scattering angle $\theta^\star$ and the beam momentum p by:

$$-t = \left(\theta^\star \times p\right)^2 \ , \tag{3}$$

where for the nominal beam momentum $p = 3988 \pm 26$ GeV is assumed [19] and the scattering angle is calculated from the proton trajectories and beam optics parameters. The relevant beam optics parameters are incorporated in transport matrix elements which describe the particle trajectory from the interaction point through the magnetic lattice of the LHC to the RPs. Several methods were developed for the reconstruction of the scattering angle, as detailed in Ref. [1]. The *subtraction method* has the best resolution and is selected as the nominal method. It uses only the track positions ($w = \{x, y\}$) and the matrix element $M_{12} = \sqrt{\beta \times \beta^\star} \sin \psi$, where ψ refers to the phase advance of the betatron function at the RP:

$$\theta^\star_w = \frac{w_A - w_C}{M_{12,A} + M_{12,C}} \ . \tag{4}$$

Here A refers to the left side of the IP at positive z and C refers to the right side at negative z. Three alternative methods are defined in detail in Ref. [1]. The *local angle method* uses only the M_{22} matrix element and the track angle between the inner and outer detectors. The *local subtraction method* uses a combination of M_{11} and M_{12} matrix elements and both the local angle and track position. The *lattice method* also uses both track parameters and reconstructs the scattering angle by an inversion of the transport matrix. The alternative methods are used to impose constraints on the beam optics and to cross-check the subtraction method.

3.2. Data taking

The low-luminosity, high $\beta^\star$ run had 108 colliding bunches with about 7×10^{10} protons per bunch, but only 3 well-separated bunches of low emittance were selected for triggering. Precise positioning of the RPs is achieved with a beam-based alignment procedure which determines the position of the RPs with respect to the proton beams by monitoring the rate of the LHC beam-loss monitors during the RP insertion. The data were collected with the RPs at a distance of approximately 7.5 mm from the beam centre, corresponding to 9.5 times the vertical beam width. The beam centre and width monitored by LHC beam position monitors and the ATLAS beam-spot measurement [20] were found to be stable to within 10 μm during the run. The beam emittance was derived from the width of the luminous region in conjunction with the beam optics. It was supplemented by direct measurement from ALFA in the vertical plane. The luminosity-weighted average of the emittance in the vertical plane was determined to be 1.6 μm for both beams and between 1.8 μm and 2.5 μm for beam 1 and beam 2 respectively in the horizontal plane. The emittance uncertainty is about 10%.

To trigger on elastic-scattering events a coincidence was required between the A- and C-sides, where on each side at least one trigger signal in a detector of the corresponding arm was required. The trigger efficiency was determined from a data stream recorded with looser conditions to be 99.9% with negligible uncertainty. The dead-time fraction of the data acquisition system (DAQ) for the selected period was 0.4%.

3.3. Track reconstruction and alignment

A well-reconstructed elastic-scattering event consists of local tracks from the proton trajectory in all four ALFA stations. The reconstruction method assumes that the protons pass through the fibre detector perpendicularly. The average multiplicity per detector is about 23 hits, where typically 18–19 are attributed to the proton trajectory while the remaining 4–5 hits are due to beam-related background, cross-talk and electronic noise. Tracks are reconstructed in several steps from the overlap area of the hit fibres and several selections are applied [1] in order to reject events with hadronic shower developments.

The precise detector positions with respect to the circulating beams are crucial inputs for the reconstruction of the proton kinematics. First, the distance between the upper and the lower detectors is determined by the use of dedicated ALFA overlap detectors which allow simultaneous measurements of the same particle in the upper and lower half of a station. Then, the detector positions are directly determined from the elastic-scattering data, using the fact that the high-$\beta^\star$ optics and the azimuthal symmetry of the scattering angle result in elastic hit patterns that have an ellipsoidal shape elongated in the vertical direction. Three alignment parameters are determined for each detector: the horizontal and vertical offsets and the rotation angle around the beam axis. For the horizontal offset the centre of the x-distribution is taken and the rotation is obtained from a linear fit to a profile histogram of the x–y correlation. The vertical offset is obtained from a comparison of the yields in the upper and lower detectors using the sliding window technique [1]. The above procedures provide an independent alignment of each ALFA station. The vertical alignment parameters are in addition fine-tuned, exploiting the strong correlations between positions of tracks measured by different detectors in elastic events. First, the positions measured in one detector are extrapolated to the other detectors in the same arm using the ratio of the appropriate M_{12} matrix elements. Then, the extrapolated positions are compared to the corresponding measurements – the

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[1] ATLAS uses a right-handed coordinate system with its origin at the nominal IP in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring and the y-axis points upwards.

[2] The β-function determines the variation of the beam envelope around the ring and depends on the focusing properties of the magnetic lattice; its value at the IP is denoted by $\beta^\star$.
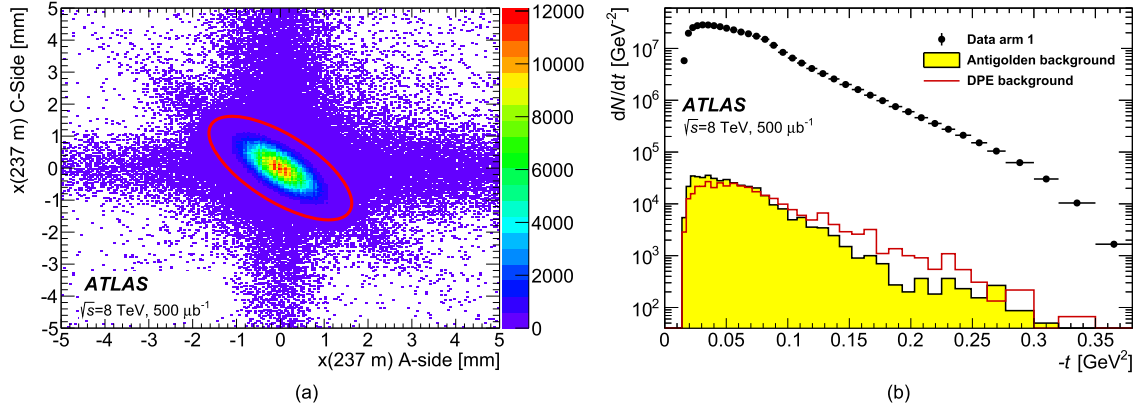
Fig. 1. (a) The correlation between the horizontal coordinates on the A- and C-sides. Elastic-scattering candidates after data quality, trigger and bunch selection but before acceptance and background rejection cuts are shown. Identified elastic events are required to lie inside the ellipse. (b) The distribution dN/dt, before corrections, as a function of t in arm 1 compared to the background spectrum determined using anti-golden events. The results of a simulation of the DPE background is also shown for comparison.

average distance gives information about residual misalignments. The residuals obtained for all pairs of detectors are combined with the vertical offset and distance measurements in a global χ^2 fit, resulting in the final alignment parameters.

4. Model for elastic scattering simulation

Several parameterizations are available [21–31] for the differential elastic pp cross section. A conventional approach is adopted here by taking the following simplified formulae:

$$\frac{d\sigma}{dt} = \frac{1}{16\pi} \left| f_N(t) + f_C(t)e^{i\alpha\phi(t)} \right|^2 \;, \qquad (5)$$

$$f_C(t) = -8\pi\alpha\hbar c \frac{G^2(t)}{|t|} \;, \qquad (6)$$

$$f_N(t) = (\rho + i)\frac{\sigma_{tot}}{\hbar c}e^{-B|t|/2} \;, \qquad (7)$$

where G is the electric form factor of the proton, B the nuclear slope, f_C the Coulomb amplitude and f_N the nuclear amplitude with ϕ their relative phase shift. The value of $\rho = \mathrm{Re}(f_{el})/\mathrm{Im}(f_{el}) = 0.1362 \pm 0.0034$ is taken from a global fit to lower-energy data [2] and parameterizations for G and ϕ are given in Ref. [1]. This expression is used to fit the data and extract σ_{tot} and B.

Monte Carlo simulation of elastic-scattering events is performed with PYTHIA8 [32,33] version 8.186 with a t-spectrum generated according to Eq. (5). The simulation is used to calculate acceptance and unfolding corrections. In the simulation the angular divergence of beams at the IP and the spread of the production vertex are set to the measured values. Elastically scattered protons are transported from the interaction point to the RPs nominally by means of the transport matrix. For studies of systematic uncertainties this was also done by the tracking module of the MadX [34] beam optics calculation program. A fast parameterization of the detector response is used in the simulation and tuned to reproduce the measured difference in position between the outer detectors and their position as extrapolated from the inner detectors.

5. Data analysis

5.1. Event selection

Events are required to pass the trigger conditions for elastic-scattering events and have a reconstructed track in all four detectors of an arm in the golden topology. The fiducial volume is

defined by cuts on the vertical coordinate of the reconstructed track, which is required to be at least 90 µm from the detector edge near the beam and at least 1 mm away from the shadow of the beam screen, in each of the four detectors.[3] The values of cuts are chosen to obtain good agreement between data and simulation in the position distributions. The back-to-back topology of elastic events is further exploited to clean the sample by imposing cuts on the left-right acollinearity. The difference between the absolute value of the vertical coordinate at the A- and C-side is requested to be below 3 mm. For the horizontal coordinate the correlation of the A- and C-sides is used. Events are selected inside an ellipse with half-axis values of 3.5σ of the resolution determined by simulation, as illustrated in Fig. 1(a). Elastic events are concentrated inside a narrow ellipse with negative slope, whereas the beam-halo background appears in broad uncorrelated bands. The most efficient selection against background is obtained from the correlation between the position in the horizontal plane and the local angle between two stations, where events on either side are again required to be inside an ellipse of 3.5σ width. From an initial sample of 4.2 million elastic candidates, 3.8 million golden elastic events were selected after all cuts. The t-spectrum, before corrections, for selected elastic events in one arm is shown in Fig. 1(b).

5.2. Background estimate

A small fraction of the events inside the selected elliptical area shown in Fig. 1(a) are expected to be background, predominantly originating from double-Pomeron exchange (DPE) according to simulations based on the MBR model [35]. The background is estimated with a data-driven method [1] using events in the "anti-golden" topology with two tracks in both upper or both lower detectors at the A- and C-sides. This sample is free of signal and yields an estimate of background in the elastic sample with the golden topology. The shape of the t-spectrum for background events is obtained by flipping the sign of the vertical coordinate on either side. The resulting background distribution is shown in Fig. 1(b). In total 4400 background events are estimated to be in the selected sample, corresponding to a fraction of 0.12% of the selected events. The systematic uncertainty is about 50%, as derived in Ref. [1] from a comparison of different methods.

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[3] The beam screen is a protection element of the quadrupoles, which limits the acceptance of the detector at large $|y|$.

5.3. Reconstruction efficiency

The rate of elastic-scattering events is corrected for reconstruction inefficiencies. These events may not be reconstructed when protons or halo particles interact with the stations or detectors, causing a shower to develop and resulting in high fibre hit multiplicities. This correction is called the event reconstruction efficiency and is given by

$$\varepsilon_{\text{rec}} = \frac{N_{\text{reco}}}{N_{\text{reco}} + N_{\text{fail}}} \, , \qquad (8)$$

for each arm where N_{reco} is the number of reconstructed events and N_{fail} the number of events for which the reconstruction failed because a shower developed. The sample of failed events is split into different categories depending on the number of detectors with reconstruction failures, because the event background is different for each category. The fraction of elastic events in the subsample where one out of four detectors failed to reconstruct a track is above 99%, whereas this fraction is 95% for the subsample where two detectors failed to reconstruct a track on one side. The event yields in the different categories are calculated with a data-driven method, for which the details are given in Ref. [1]. The background fraction in the case with only two detectors with reconstructed tracks is estimated with background templates of the x distribution, obtained from data by selecting single diffractive events. In the case of a successful track reconstruction in three detectors, where a good t-measurement is still possible, the partial reconstruction efficiency was verified to be independent of t, which is then also assumed for the other categories. Events falling outside the acceptance, but faking a signal through shower development, were eliminated from the reconstruction efficiency calculation by applying another template analysis using the y distribution obtained from golden elastic events.

The event reconstruction efficiencies in arm 1 and arm 2 are determined to be $\varepsilon_{\text{rec},1} = 0.9050 \pm 0.0003$ (stat.) ± 0.0034 (syst.) and $\varepsilon_{\text{rec},2} = 0.8883 \pm 0.0003$ (stat.) ± 0.0045 (syst.), respectively. The lower reconstruction efficiency in arm 2 originates from a different amount of material which induces a higher probability of shower development. The systematic uncertainty is estimated by a variation of the selection criteria and templates, as described in Ref. [1].

5.4. Beam optics

The precision of the t-reconstruction depends on knowledge of the transport matrix elements. A data-driven method was developed [1] to tune the relevant matrix elements using constraints on the beam optics derived from measured correlations in the ALFA data. These constraints are incorporated in a fit of the strength of the inner triplet quadrupole magnets Q1 and Q3, which yields an effective beam optics used in the simulation. The values of the constraints are compatible with those published in Ref. [1] within 15% and the resulting magnet strength offsets are in good agreement with the values found at 7 TeV.

5.5. Acceptance and unfolding

The acceptance is defined as the ratio of events passing all geometrical and fiducial acceptance cuts to all generated events, and is calculated as a function of t. The form of the acceptance curve as shown in Fig. 2 results from the different contributions of the vertical and horizontal scattering angles to the value of t and the impact of the fiducial volume cuts on these contributions. In particular, the position of the peak depends on the cut at large $|y|$ at the beam screen, which is slightly different for the two arms. The
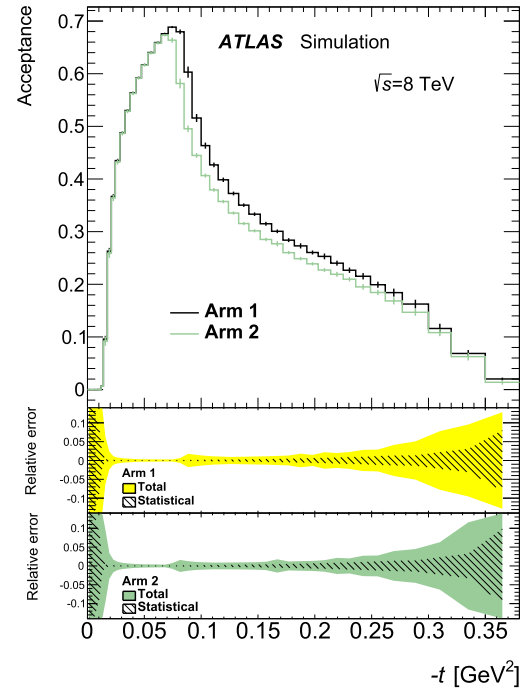

Fig. 2. The acceptance as a function of the true value of t for each arm with total uncertainties shown as error bars. The lower panels show relative total and statistical uncertainties.

rise of the acceptance at small t is different in the two arms because of different detector distances, between 8 and 8.4 mm, to the beam.

The measured t-spectrum is affected by detector resolution and beam divergence effects, which are corrected with an unfolding procedure. The t-resolution of the subtraction method is about 10% at small t and 3% at large t. The alternative methods have a t-resolution which is a factor of 2–3 worse [1]. The background-subtracted distributions in each arm are corrected for migration effects using an iterative, dynamically stabilized, unfolding method [36], which is based on a simulated transition matrix describing the resolution-induced migration between bins of the t-spectrum. The corrections induced by the unfolding are small ($<2\%$) for the subtraction method except at small t where they rise to 30%. For the other methods the corrections are generally t-dependent and increase to 50% at large t.

5.6. Luminosity

The ATLAS luminosity measurement at high luminosity ($L > 10^{33}$ cm^{-2} s^{-1}) is described in detail in Ref. [37]. Unlike that measurement, the run in this analysis had an instantaneous luminosity $L \sim 0.05 \cdot 10^{30}$ cm^{-2} s^{-1}, about five orders of magnitude lower. Only three bunches were present in this run, whereas more than a thousand bunches are common at high luminosity. The average number of interactions per bunch-crossing (pile-up) in this sample is $\mu \sim 0.1$, which is also low compared to the values of $\mu = 10$–40 reached routinely in normal conditions. At such low values of the luminosity, some of the standard algorithms are unusable due to lack of sensitivity. On the other hand, an additional method based on vertex counting in the inner detector (ID) can be exploited, which is most effective at low pile-up. Another consequence of the low luminosity is the relative importance of the background sources: the beam–gas contribution, normally negligible, can become comparable with the collision rate, while the "afterglow"

background (see Ref. [37]) becomes conversely less important, due to the small number of colliding bunches.

In 2012, the beam conditions monitor (BCM) was used as the baseline detector for luminosity measurements. It consists of diamond-sensor detectors placed on both sides of the IP. It measures the luminosity using an event-counting method based on the requirement of having activity in either side (BCM_EventOR). LUCID (LUminosity measurement with a Cherenkov Integrating Detector) is also located on both sides of the IP and uses the same algorithm to measure the luminosity (LUCID_EventOR). A third method for measuring the per-bunch luminosity is provided by the ID. Tracks are reconstructed requiring at least nine hits and no missing hits along the track trajectory, and a transverse momentum $p_T > 900$ MeV. Then, at least five selected tracks are required to form a primary vertex (VTX5). The number of primary vertices per event is proportional to the luminosity and provides an independent method with respect to LUCID and BCM.

The absolute luminosity scale of each algorithm was calibrated by the van der Meer (vdM) method [38] at an intermediate luminosity regime ($L \sim 10^{30}$ cm^{-2} s^{-1}). The treatment of both afterglow and beam–gas background is described in detail in Ref. [37]. The first is evaluated by measuring the detector activity in unfilled bunches preceding the colliding bunches, while the second is estimated from the so-called unpaired bunches, in which bunches in only one of the two beams are filled and no beam–beam collisions occur. In the high-$\beta^\star$ run and for BCM_EventOR, the afterglow background is evaluated to be 0.05% and the beam–gas contribution is 0.4%.

BCM_EventOR was chosen as the baseline algorithm for the luminosity determination, whereas the LUCID_EventOR and VTX5 methods are only used for the evaluation of systematic uncertainties. It proved to be the most stable, both by comparing the various vdM calibration sessions performed during the year and by studying its long-term behaviour at high luminosity. This choice also ensures maximum compatibility with the high-luminosity case. By comparing the LUCID_EventOR and VTX5 results with BCM_EventOR, a maximum difference of 0.3% is found. No change of this difference with time, or equivalently μ, is observed.

The following contributions to the systematic uncertainty of the luminosity determination are considered:

- The absolute luminosity scale, common to all algorithms, is determined by the vdM method. Its uncertainty of 1.2% is dominated by the beam conditions. This uncertainty is fully correlated between low- and high-luminosity data sets [37].
- The BCM calibration stability between the high-$\beta^\star$ run and the vdM session is estimated to be 0.8% by comparing with the VTX5 method among the various vdM scans.
- The afterglow background uncertainty is conservatively taken as 100% of the afterglow level itself, which leads to an uncertainty of 0.05% in the luminosity.
- The beam–gas background uncertainty is obtained using LUCID by comparing the difference in the off-time activity (i.e. produced by beam–gas interactions and not by collisions at the IP) between the colliding and the unpaired bunches. It is estimated to be 0.3%.

The total systematic uncertainty is therefore 1.5%. The final integrated luminosity is measured to be $L_{int} = 496.3 \pm 0.3$ (stat.) $\pm$ 7.3 (syst.) μb^{-1}.

6. Results

6.1. Elastic cross section

The differential elastic cross section in a given bin t_i is calculated from the following formula:

$$\frac{d\sigma_{el}}{dt_i} = \frac{1}{\Delta t_i} \times \frac{\mathcal{M}^{-1}[N_i - B_i]}{A_i \times \epsilon^{reco} \times \epsilon^{trig} \times \epsilon^{DAQ} \times L_{int}} \, , \qquad (9)$$

where Δt_i is the width of the bins in t, $\mathcal{M}^{-1}$ symbolizes the unfolding procedure applied to the background-subtracted number of events $N_i - B_i$, A_i is the acceptance, ϵ^{reco} is the event reconstruction efficiency, ϵ^{trig} is the trigger efficiency, ϵ^{DAQ} is the dead-time correction and L_{int} is the integrated luminosity. The binning in t is chosen to yield a purity above 50%, which corresponds to 1.5 times the resolution at small t. It is enlarged at large t in order to account for the lower number of events. The numerical values for the resulting differential elastic cross section are given in Table 1.

The experimental systematic uncertainties are derived according to the methods detailed in Ref. [1] as follows:

- The value of the beam momentum used in the t-reconstruction (Eq. (3)) and in the simulation is varied by 0.65%, as recommended in Ref. [19].
- The uncertainty in the luminosity of 1.5% is applied to the cross-section normalization.
- The event reconstruction efficiency is varied by its uncertainty of about 0.5% and the uncertainty in the tracking efficiency is estimated by varying the reconstruction criteria.
- The uncertainties originating from the effective beam optics are calculated from variations of the optics constraints, of the strength of the quadrupoles not adjusted in the fit, and of the quadrupole alignment constants. Additional uncertainties are related to the error of the optics fit, to the beam transport scheme used in the simulation, and to the impact from a residual beam crossing angle assumed to vary within its uncertainty of ± 10 μrad.
- The uncertainties from the alignment of the ALFA detectors are evaluated by varying the correction constants for horizontal and vertical offsets as well as the rotation within their uncertainties as determined from variations of the alignment procedures, and by taking the difference between different optimization configurations for the vertical alignment parameters.
- The background normalization uncertainty of 50% is applied in the background subtraction and the background shape is varied by inverting the sign of different detector combinations.
- The detector resolution values in the fast simulation are replaced by estimates from GEANT4 [39,40] and test-beam measurements, and a y-dependent resolution is used instead of a constant value.
- The value of the nuclear slope in the simulation is varied around the nominal value of 19.7 GeV^{-2} by ± 1 GeV^{-2}, corresponding to about five times the uncertainty of the measured B value.
- The beam emittance value in the simulation is varied by its uncertainty of about 7%. Additionally, the ratio of the emittance in beam 1 to the emittance in beam 2, which are measured by wire scans after injection only, is set to unity.
- The intrinsic unfolding uncertainty is estimated from a data-driven closure test.

The main sources of systematic uncertainty are the beam momentum uncertainty and the luminosity uncertainty. For each system-

Table 1
The measured values of the differential elastic cross section with statistical and systematic uncertainties. The central t-values in each bin are calculated from simulation, in which a slope parameter of $B = 19.7$ GeV^{-2} is used.

Low $|t|$ edge [GeV^2]	High $|t|$ edge [GeV^2]	Central $|t|$ [GeV^2]	$d\sigma_{el}/dt$ [mb/GeV^2]	Stat. uncert. [mb/GeV^2]	Syst. uncert. [mb/GeV^2]	Total uncert. [mb/GeV^2]
0.0090	0.0120	0.0105	387	29	14	32
0.0120	0.0140	0.0130	370	5.6	12	13
0.0140	0.0175	0.0157	352.3	1.4	8.7	8.9
0.0175	0.0210	0.0192	329.8	0.8	6.5	6.5
0.0210	0.0245	0.0227	306.9	0.6	5.7	5.8
0.0245	0.0285	0.0265	284.6	0.5	5.0	5.1
0.0285	0.0330	0.0307	261.7	0.4	4.6	4.6
0.0330	0.0375	0.0352	239.3	0.4	4.1	4.1
0.0375	0.0425	0.0400	218.0	0.4	3.6	3.6
0.0425	0.0475	0.0450	197.3	0.3	3.3	3.3
0.0475	0.0530	0.0502	178.0	0.3	3.0	3.0
0.0530	0.0590	0.0559	158.8	0.2	2.7	2.7
0.0590	0.0650	0.0619	141.1	0.2	2.4	2.4
0.0650	0.0710	0.0679	126.0	0.2	2.2	2.2
0.0710	0.0780	0.0744	111.1	0.2	2.0	2.0
0.0780	0.0850	0.0814	96.8	0.2	2.0	2.0
0.0850	0.0920	0.0884	84.7	0.2	1.7	1.7
0.0920	0.1000	0.0959	72.9	0.2	1.6	1.6
0.1000	0.1075	0.1037	62.7	0.2	1.5	1.5
0.1075	0.1150	0.1112	54.1	0.2	1.3	1.4
0.1150	0.1240	0.1194	46.11	0.14	1.13	1.13
0.1240	0.1330	0.1284	38.76	0.14	1.0	1.01
0.1330	0.1420	0.1374	32.60	0.12	0.92	0.93
0.1420	0.1520	0.1468	27.10	0.11	0.82	0.83
0.1520	0.1620	0.1568	22.48	0.11	0.74	0.74
0.1620	0.1720	0.1668	18.48	0.10	0.68	0.68
0.1720	0.1820	0.1768	15.25	0.09	0.67	0.68
0.1820	0.1930	0.1873	12.36	0.08	0.57	0.58
0.1930	0.2030	0.1978	10.08	0.08	0.48	0.48
0.2030	0.2140	0.2083	8.20	0.07	0.43	0.43
0.2140	0.2250	0.2193	6.58	0.06	0.33	0.33
0.2250	0.2360	0.2303	5.34	0.06	0.27	0.28
0.2360	0.2490	0.2422	4.28	0.05	0.24	0.24
0.2490	0.2620	0.2552	3.30	0.05	0.22	0.23
0.2620	0.2770	0.2691	2.47	0.04	0.18	0.18
0.2770	0.3000	0.2877	1.69	0.03	0.14	0.14
0.3000	0.3200	0.3094	1.06	0.03	0.10	0.1
0.3200	0.3500	0.3335	0.62	0.02	0.08	0.08
0.3500	0.3800	0.3635	0.36	0.04	0.04	0.05

atic uncertainty source the shift of the cross-section value in each t-bin is recorded. The most important shifts are shown in Fig. 3(a).

6.2. Total cross section

A profile fit [41] is used to determine σ_{tot}. It includes statistical and systematic uncertainties and their correlations across the t-spectrum. For each shift due to a systematic uncertainty a nuisance parameter is fitted in a procedure described in Ref. [1].

The theoretical prediction of Eq. (5) including the Coulomb and interference terms is fitted to the data to extract σ_{tot} and B alongside the nuisance parameters, as shown in Fig. 3(b). The fit range is chosen to be from $-t = 0.014$ GeV^2 to $-t = 0.1$ GeV^2, where the lower bound is set by requiring the acceptance to exceed 10% and the upper bound is chosen to exclude the large-t region where theoretical models predict deviations from a single exponential function [42]. The fit yields $\sigma_{tot} = 96.07 \pm 0.86$ mb and $B = 19.74 \pm 0.17$ GeV^{-2} with $\chi^2/N_{dof} = 17.8/14$ and the uncertainties include all statistical and experimental systematic contributions. The most important uncertainty component is the luminosity error for σ_{tot} and the beam energy error for B. Additional uncertainties arising from the extrapolation $t \to 0$ are estimated from a variation of the upper end of the fit range respectively up to $-t = 0.152$ GeV^2 and up to $-t = 0.065$ GeV^2, and from a variation of the lower end, i.e. from $-t = 0.009$ GeV^2 to $-t = 0.0245$ GeV^2. Further theoretical uncertainties considered

include: a variation of the ρ-parameter in Eq. (1) by ± 0.0034; the replacement of the dipole parameterization by a double-dipole parameterization [43] for the proton electric form factor; the replacement of the Coulomb phase from West and Yennie [22] by parameterizations from Refs. [24,27]; the inclusion of a term related to the magnetic moment of the proton in the Coulomb amplitude [23]. The dominant extrapolation uncertainty is induced by the fit range variation. The final results for σ_{tot} and B are:

$$\sigma_{tot} = 96.07 \pm 0.18 \text{ (stat.)} \pm 0.85 \text{ (exp.)} \pm 0.31 \text{ (extr.) mb}, \quad (10)$$

$$B = 19.74 \pm 0.05 \text{ (stat.)} \pm 0.16 \text{ (exp.)} \pm 0.15 \text{ (extr.) GeV}^{-2}. \quad (11)$$

A summary of the results for σ_{tot} from four different t-reconstruction methods is given in Table 2. The results from the nominal subtraction method are in good agreement with the other methods, considering the uncorrelated uncertainty of 0.3–0.4 mb. The alternative methods are correlated through the common use of the local angle variable.

Further stability checks are carried out in order to cross-check the fitting method. A fit using only the covariance matrix of statistical uncertainties yields $\sigma_{tot} = 96.34 \pm 0.07$ (stat.) in good agreement with the results from the profile fit Eq. (10). The same fit with only statistical uncertainties was also performed for the two arms of ALFA independently and gave consistent results within one standard deviation of the statistical uncertainty. The data sample was split into ten sub-periods with roughly equal numbers of selected events and no dependence of the measured value of σ_{tot} on
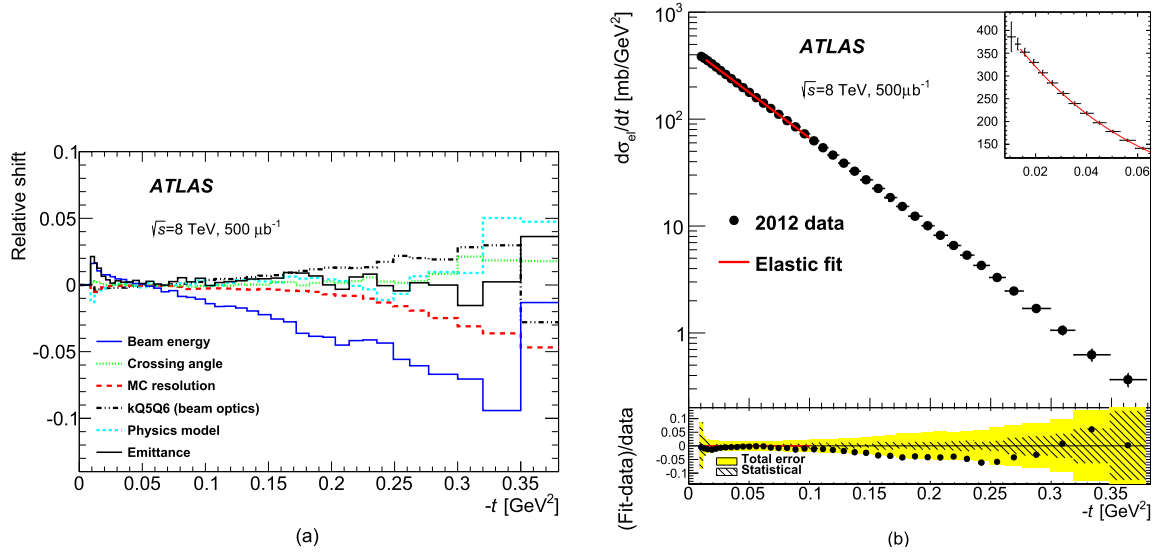
Fig. 3. (a) Relative shifts in the differential elastic cross section as a function of t for selected systematic uncertainty sources. Shown are the uncertainties related to the beam energy, to the crossing angle, to the modelling of the detector resolution in the simulation (MC resolution), to the beam optics (kQ5Q6, magnet strength), to the value of B in the simulation (Physics model) and to the emittance. (b) The fit of the theoretical prediction to the differential elastic cross section with σ_{tot} and B as free parameters. In the lower plot the points represent the relative difference between fit and data, the yellow area represents the total experimental uncertainty and the hatched area the statistical component. The red line indicates the fit range; the fit result is extrapolated in the lower plot outside the fit range. The upper right insert shows a zoom of the data and fit at small t. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

Table 2
The total cross section and uncertainties for four different t-reconstruction methods. The nominal results are based on the subtraction method, quoted in the second column.

	σ_{tot} [mb]			
	Subtraction	Local angle	Lattice	Local subtraction
Total cross section	96.07	96.52	96.56	96.58
Statistical error	0.18	0.15	0.16	0.15
Experimental error	0.85	0.94	0.88	0.89
Extrapolation error	0.31	0.42	0.23	0.23
Total error	0.92	0.98	0.93	0.93

time was observed. Also, the data from the three different bunches were investigated independently and found to give consistent results. Finally the stability of the analysis was tested by a wide variation of the event selection cuts. The largest change of σ_{tot} with these cut variations was observed for the cut on the correlation between x and θ_x. That produced a change of ± 0.3 mb, well within the t-dependent experimental systematic uncertainty of about 0.5 mb. Several alternative parameterizations [22,26,28, 27,29–31] of the differential elastic cross section, including nonexponential forms at large t, were used to fit the spectrum up to $-t = 0.3$ GeV^2 in order to assess the impact on the value of the total cross section. The RMS of the values obtained is 0.28 mb, in good agreement with the quoted extrapolation uncertainty of 0.31 mb assigned to the simple exponential form.

The TOTEM Collaboration exploited data from the same LHC fill for a measurement of σ_{tot} using the luminosity-independent method. Their result is $\sigma_{tot} = 101.7 \pm 2.9$ mb [3], higher than the measurement presented here. The difference corresponds to 1.9σ assuming uncorrelated uncertainties. Better agreement is observed in the nuclear slope measurement, where TOTEM reports $B = 19.9 \pm 0.3$ GeV^{-2}, a value very close to the present result $B = 19.74 \pm 0.19$ GeV^{-2}, which indicates that the difference is confined to the normalization. The measurements of ATLAS and TOTEM are compared to measurements at lower energy and to a global fit [2] in Fig. 4(a) for σ_{tot} and in Fig. 4(b) for B. TOTEM also reported evidence of non-exponential behaviour of the differential

elastic cross section [49] in the $-t$-range below 0.2 GeV^2, where deviations from the single exponential form of the order of one percent are observed. Such effects cannot be substantiated with this data set because their size is below the systematic uncertainties of the present measurement.

As well as the total cross section, the total integrated elastic cross section can be calculated, provided that the Coulomb amplitude is neglected. In this case, σ_{el} can be obtained from the formula

$$\sigma_{el} = \frac{\sigma_{tot}^2}{B} \frac{1 + \rho^2}{16\pi (\hbar c)^2} , \qquad (12)$$

and the result is $\sigma_{el} = 24.33 \pm 0.04$ (stat.) ± 0.39 (syst.) mb. The measured integrated elastic cross section in the fiducial range from $-t = 0.009$ GeV^2 to $-t = 0.38$ GeV^2 corresponds to 80% of this total elastic cross section $\sigma_{el}^{observed} = 19.67 \pm 0.02$ (stat.) ± 0.33 (syst.) mb. The total inelastic cross section is determined by subtraction of the total elastic cross section from the total cross section. The resulting value is $\sigma_{inel} = 71.73 \pm 0.15$ (stat.) ± 0.69 (syst.) mb.

7. Conclusion

ATLAS has performed a measurement of the total cross section from elastic pp scattering at $\sqrt{s} = 8$ TeV. The measurement is based on 500 μb^{-1} of collision data collected in a high-$\beta^\star$ run at the LHC in 2012 with the ALFA Roman Pot sub-detector. The optical theorem is used to extract the total cross section from the differential elastic cross section by extrapolating $t \to 0$. The differential cross section is also used to determine the nuclear slope. The analysis uses data-driven methods to determine relevant beam optics parameters and event reconstruction efficiency, and to tune the simulation. The detailed evaluation of the associated systematic uncertainties is supplemented by a comparison of t-reconstruction methods with different sensitivities to beam optics. The absolute luminosity for this run is determined in a dedicated analysis, taking into account the special conditions with a very low number of interactions per bunch crossing. The total cross section at $\sqrt{s} = 8$ TeV is determined to be
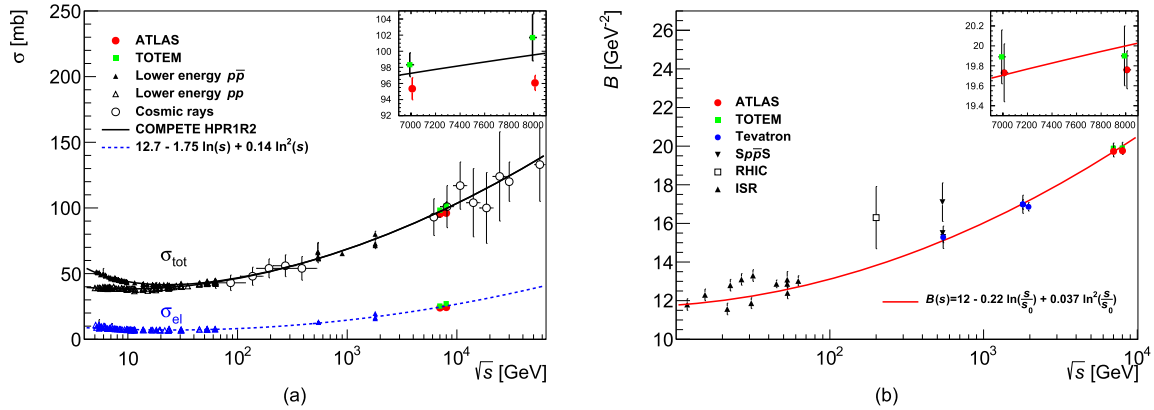
Fig. 4. (a) Comparison of total and elastic cross-section measurements presented here with other published measurements [2,5,44–47] and model predictions as a function of the centre-of-mass energy. (b) Comparison of the measurement of the nuclear slope B presented here with other published measurements at the ISR, at the S$p\bar{p}$S, at RHIC, at the Tevatron and with the measurement from TOTEM at the LHC. The red line shows a model calculation [48], which contains a linear term and quadratic term in $\ln s$. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

$$\sigma_{\text{tot}}(pp \to X) = 96.07 \pm 0.18 \text{ (stat.)} \pm 0.85 \text{ (exp.)}$$
$$\pm 0.31 \text{ (extr.) mb ,}$$

where the first error is statistical, the second accounts for all experimental systematic uncertainties and the last is related to uncertainties in the extrapolation $t \to 0$. In addition, the slope of the elastic differential cross section at small t is determined to be $B = 19.74 \pm 0.05 \text{ (stat.)} \pm 0.23 \text{ (syst.)} \text{ GeV}^{-2}$.

The total elastic cross section is extracted from the fitted parameterization as $\sigma_{\text{el}}(pp \to pp) = 24.33 \pm 0.04$ (stat.) ± 0.39 (syst.) mb and the inelastic cross section is obtained by subtraction from the total cross section as $\sigma_{\text{inel}} = 71.73 \pm 0.15$ (stat.) ± 0.69 (syst.) mb. The measurements at 8 TeV are significantly more precise than the previous measurements at 7 TeV because of the smaller luminosity uncertainty and a larger data sample.

Acknowledgements

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