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Ryota Takabe, Suguru Yachi, Weijie Du, Daichi Tsukahara, Hiroki Takeuchi, Kaoru Toko, and Takashi Suemasu

Citation: *AIP Advances* **6**, 085107 (2016); doi: 10.1063/1.4961063

View online: <http://dx.doi.org/10.1063/1.4961063>

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## Influence of air exposure duration and a-Si capping layer thickness on the performance of p-BaSi<sub>2</sub>/n-Si heterojunction solar cells

Ryota Takabe,<sup>1</sup> Suguru Yachi,<sup>1</sup> Weijie Du,<sup>2</sup> Daichi Tsukahara,<sup>1</sup>  
 Hiroki Takeuchi,<sup>1</sup> Kaoru Toko,<sup>1</sup> and Takashi Suemasu<sup>1,a</sup>

<sup>1</sup>*Institute of Applied Physics, University of Tsukuba, Tsukuba, Ibaraki 305-8573, Japan*

<sup>2</sup>*Key Laboratory of Optoelectronic Material and Device, College of Mathematics and Science, Shanghai Normal University, Shanghai 200234, China*

(Received 2 May 2016; accepted 2 August 2016; published online 10 August 2016)

Fabrication of p-BaSi<sub>2</sub>(20nm)/n-Si heterojunction solar cells was performed with different a-Si capping layer thicknesses ( $d_{\text{a-Si}}$ ) and varying air exposure durations ( $t_{\text{air}}$ ) prior to the formation of a 70-nm-thick indium-tin-oxide electrode. The conversion efficiencies ( $\eta$ ) reached approximately 4.7% regardless of  $t_{\text{air}}$  (varying from 12–150 h) for solar cells with  $d_{\text{a-Si}} = 5$  nm. In contrast,  $\eta$  increased from 5.3 to 6.6% with increasing  $t_{\text{air}}$  for those with  $d_{\text{a-Si}} = 2$  nm, in contrast to our prediction. For this sample, the reverse saturation current density ( $J_0$ ) and diode ideality factor decreased with  $t_{\text{air}}$ , resulting in the enhancement of  $\eta$ . The effects of the variation of  $d_{\text{a-Si}}$  (0.7, 2, 3, and 5 nm) upon the solar cell performance were examined while keeping  $t_{\text{air}} = 150$  h. The  $\eta$  reached a maximum of 9.0% when  $d_{\text{a-Si}}$  was 3 nm, wherein the open-circuit voltage and fill factor also reached a maximum. The series resistance, shunt resistance, and  $J_0$  exhibited a tendency to decrease as  $d_{\text{a-Si}}$  increased. These results demonstrate that a moderate oxidation of BaSi<sub>2</sub> is a very effective means to enhance the  $\eta$  of BaSi<sub>2</sub> solar cells. © 2016 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>). [<http://dx.doi.org/10.1063/1.4961063>]

### I. INTRODUCTION

For future deployment of terawatt-scale solar cells, extensive research has been conducted on thin-film solar cell materials such as chalcopyrite and cadmium telluride as well as Si-based materials because of their high energy conversion efficiency ( $\eta$ ) and low cost.<sup>1–6</sup> Perovskite-based solar cells have also gained increasing attention owing to their astonishing increase in efficiency.<sup>7,8</sup> However, these materials contain critical raw materials such as In, Cd, and Pb. There has also been growing interest in Si thin-film solar cells that employ an efficient light-trapping system,<sup>9–15</sup> but with this system it is not easy to achieve an  $\eta$  as high as 20%. Hence, it is necessary to explore alternative materials for thin-film solar cell applications. Among such materials, much attention has been given to the semiconductor BaSi<sub>2</sub>, which consists of safe and earth-abundant elements and has a band gap ( $E_g$ ) of 1.3 eV, matching the solar spectrum.<sup>16</sup> One of the most striking features of this material is that both its large absorption coefficient ( $\alpha$ ) and large minority-carrier diffusion length ( $L$ ) can be used. Because of its indirect band gap, undoped BaSi<sub>2</sub> can attain a minority-carrier lifetime ( $\tau$ ) and an  $L$  value as large as approximately 10  $\mu\text{s}$  and 10  $\mu\text{m}$ , respectively.<sup>17–19</sup> These values are sufficiently large for thin-film solar cell applications. Furthermore,  $\alpha$  exceeds  $3 \times 10^4 \text{ cm}^{-1}$  for photon energies greater than 1.5 eV because the direct transition occurs at energies slightly larger than  $E_g$ .<sup>20–22</sup> For these reasons, we expect  $\eta$  to be larger than 25% in a 2- $\mu\text{m}$ -thick BaSi<sub>2</sub> pn junction diode.<sup>23</sup>

<sup>a</sup>[suemasu@bk.tsukuba.ac.jp](mailto:suemasu@bk.tsukuba.ac.jp)



An a-Si capping layer plays an important role in BaSi<sub>2</sub> solar cells, whereby an undoped n-BaSi<sub>2</sub> surface with a few nm-thick a-Si layer can exhibit a  $\tau \sim 10$   $\mu$ s with excellent repeatability.<sup>24</sup> Measurements of the valence band offset at the a-Si/BaSi<sub>2</sub> interface by hard x-ray photoelectron spectroscopy have shown that the barrier height of the a-Si layer for the minority carrier (i.e., holes) in the n-BaSi<sub>2</sub> is  $-0.2$  eV,<sup>25</sup> whereas that of the native oxide layer is  $3.9$  eV.<sup>26</sup> Therefore, the a-Si capping layer works as a good electrical contact for hole transport as well as a passivation layer. Very recently, we have attained  $\eta = 9.0\%$ , a short-circuit current density  $J_{SC} = 31.9$  mA/cm<sup>2</sup> and an open-circuit voltage  $V_{OC} = 0.46$  V for B-doped p-BaSi<sub>2</sub>/n-Si heterojunction solar cells using an a-Si capping layer.<sup>27</sup> The a-Si layer thickness ( $d_{a-Si}$ ) and the duration of air exposure ( $t_{air}$ ) after a-Si layer deposition are important parameters that may influence the performance of BaSi<sub>2</sub> solar cells. Previous studies have shown that  $t_{air}$  and  $d_{a-Si}$  have a significant effect upon the performance of a-Si:H(n)/a-Si:H(i)/c-Si(p) heterojunctions with intrinsic thin layer (HIT) solar cells.<sup>28</sup> This sensitivity arises from the significant influence the  $t_{air}$  and i-layer thickness have upon the minority-carrier transport and recombination at the interfaces in HIT solar cells.<sup>28</sup> However, there is limited information regarding their effects upon BaSi<sub>2</sub> solar cells. In this work, we fabricated a-Si/p-BaSi<sub>2</sub>(20 nm)/n-Si heterojunction solar cells via molecular beam epitaxy (MBE) prepared with various values of  $t_{air}$  and  $d_{a-Si}$ , and attempted to clarify the influence of these parameters upon the properties of p-BaSi<sub>2</sub>/n-Si solar cells.

## II. EXPERIMENTS

For the growth of the BaSi<sub>2</sub> layers, we used an ion-pumped MBE system (AVC Co., Ltd.) equipped with an electron-beam evaporation source for Si and with standard Knudsen cells for Ba and B. The deposition rates of Si ( $R_{Si}$ ) and Ba ( $R_{Ba}$ ) were controlled using an electron impact emission spectroscopy (EIES) feedback system (INFICON CO., Ltd). We first deposited Ba on a heated Czochralski n-Si(111) (resistivity  $\rho = 1-4$   $\Omega$ -cm) substrate at  $500^\circ\text{C}$  by reactive deposition epitaxy to form a 5-nm-thick BaSi<sub>2</sub> template layer,<sup>29</sup> where  $R_{Ba}$  was set at  $1.0$  nm/min. This template layer acted as seed crystals for the subsequent BaSi<sub>2</sub> layer. We next co-deposited Ba, Si and B on the templates at  $600^\circ\text{C}$  by MBE to form a 20-nm-thick B-doped p-BaSi<sub>2</sub> layer,<sup>30-32</sup> where  $R_{Si}$  and  $R_{Ba}$  were fixed at  $0.9$  and  $2.3$  nm/min, respectively. The B concentration was set to  $2 \times 10^{18}$  cm<sup>-3</sup> to be comparable to the hole concentration.<sup>27</sup> The epitaxial growth of the BaSi<sub>2</sub> layers in all of the samples was confirmed by reflection high-energy electron diffraction and x-ray diffraction (data not shown). We first prepared four samples (samples A–D) to examine the influence of  $t_{air}$  upon the solar cell performance. For this purpose, we deposited a 2 or 5-nm-thick a-Si layer on the BaSi<sub>2</sub> surface *in situ* at  $180^\circ\text{C}$  with  $R_{Si} = 0.9$  nm/min, followed by air exposure for  $t_{air} = 12$  or  $150$  h. In the previous study,<sup>25</sup> we found that the oxidation of BaSi<sub>2</sub> do not progress for sample capped with 5-nm-thick a-Si layers even for  $t_{air} = 24$  h. Therefore, we anticipated that the oxidation of BaSi<sub>2</sub> would be suppressed much further for  $t_{air} = 12$  h. On the other hand, we expected that the oxidation would progress for  $t_{air} = 150$  h. This is the reason why we chose these two air exposure durations. As described below, the  $\eta$  increased with  $t_{air}$  for samples with  $d_{a-Si} = 2$  nm. We next prepared samples in which  $d_{a-Si}$  was varied from  $0.7$  to  $5$  nm while keeping  $t_{air} = 150$  h to examine the influence of  $d_{a-Si}$ . Please note that the a-Si layers in this study were not hydrogenated. They were just evaporated from the solid source of Si by electron beam irradiation. After keeping samples in air for  $t_{air}$ , each sample was introduced into a radio-frequency (RF) sputtering chamber, and 1-mm-diameter and 70-nm-thick indium-tin-oxide (ITO) electrodes were sputtered on the front and Al electrodes on back surfaces at room temperature. The RF power was set to  $100$  W. The solar cell properties of samples A–D are summarized in Table I.

The plots of the current density versus voltage ( $J$ - $V$ ) were measured for as many electrodes as possible in an area of  $1 \times 1$  cm<sup>2</sup> on the sample wafer under AM1.5,  $100$  mW/cm<sup>2</sup> illumination at approximately  $25^\circ\text{C}$  using a mask with holes  $1$  mm in diameter. To accurately obtain the series resistance,  $R_S$ , diode ideality factor,  $\gamma$ , and reverse saturation current density,  $J_0$ , of a diode, we adopted a technique described in Ref. 33. Using the photodiode equation, the relationship between

TABLE I. Sample and solar cell parameters for samples A–D, giving the thickness of the a-Si layer ( $d_{\text{a-Si}}$ ), duration of exposure to air ( $t_{\text{air}}$ ), short-circuit current density ( $J_{\text{SC}}$ ), open-circuit voltage ( $V_{\text{OC}}$ ), fill factor ( $FF$ ), conversion efficiency ( $\eta$ ), series resistance ( $R_{\text{S}}$ ), shunt resistance ( $R_{\text{SH}}$ ), diode ideality factor ( $\gamma$ ) and reverse saturation current density ( $J_0$ ).

Sample	$d_{\text{a-Si}}$ (nm)	$t_{\text{air}}$ (h)	$J_{\text{SC}}$ (mA/cm <sup>2</sup> )	$V_{\text{OC}}$ (V)	$FF$	$\eta$ (%)	$R_{\text{S}}$ ( $\Omega$ )	$R_{\text{SH}}$ (k $\Omega$ )	$\gamma$	$J_0$ (mA/cm <sup>2</sup> )
A	2	12	30.5	0.34	0.51	5.3	178	36.1	1.58	$1.02 \times 10^{-2}$
B	2	150	29.5	0.44	0.51	6.6	336	20.4	1.29	$2.35 \times 10^{-4}$
C	5	12	22.5	0.42	0.50	4.8	125	12.3	1.26	$4.14 \times 10^{-5}$
D	5	150	28.1	0.40	0.42	4.7	322	11.8	1.44	$8.73 \times 10^{-4}$

$R_{\text{S}}$  and  $\gamma$  can be given as

$$\frac{dV}{dJ} = SR_{\text{S}} + \frac{\gamma k_{\text{B}}T}{q} \left[ \frac{1 - (SR_{\text{SH}})^{-1} dV/dJ}{J + J_{\text{SC}} - (SR_{\text{SH}})^{-1}V} \right]. \quad (1)$$

Here,  $T$  is the absolute temperature,  $q$  is the elemental charge,  $k_{\text{B}}$  is the Boltzmann constant,  $S$  is the area of the electrode,  $R_{\text{SH}}$  is the shunt resistance, and  $J_{\text{SC}}$  is the photocurrent density. Using the plot of  $dV/dJ$  versus the term in brackets in Eq. (1), we can directly deduce  $\gamma$  from the slope and  $R_{\text{S}}$  from the intercept. The external quantum efficiency (EQE) spectra were evaluated at 25°C using a lock-in technique with a xenon lamp and a 25-cm focal-length single monochromator (Bunko Keiki, SM-1700A).

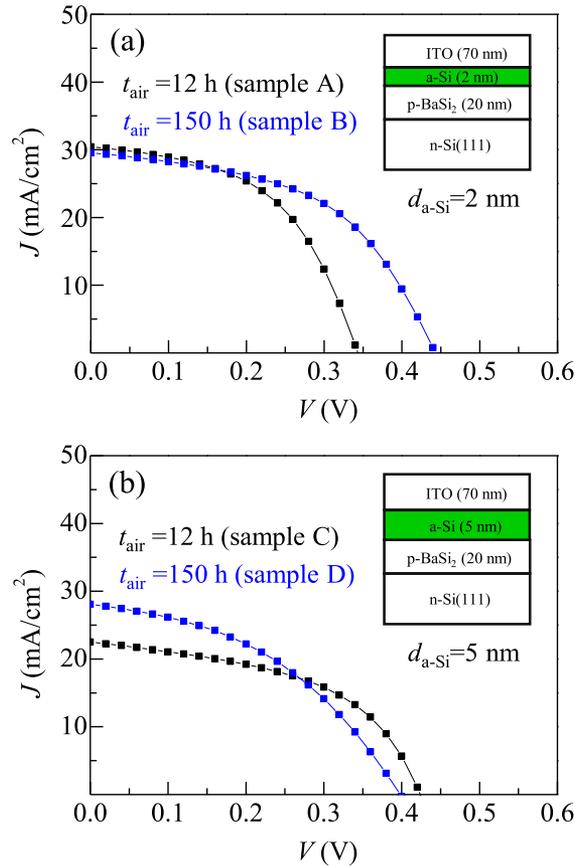


FIG. 1.  $J$ - $V$  characteristics under AM1.5 illumination measured for samples with  $d_{\text{a-Si}}$  = (a) 2 nm (samples A and B) and (b) 5 nm (samples C and D). The  $t_{\text{air}}$  was varied as 12 or 150 h.

### III. RESULTS AND DISCUSSION

First we discuss the influence of  $t_{\text{air}}$  on the solar cell performance. Figure 1(a) and 1(b) show typical examples of the  $J$ - $V$  characteristics under AM1.5 illumination on samples with  $d_{\text{a-Si}} = 2$  nm (samples A and B, Fig. 1(a)) and 5 nm (samples C and D, Fig. 1(b)), for  $t_{\text{air}}$  of 12 or 150 h. As shown in Fig. 1(a),  $V_{\text{OC}}$  drastically increases with increasing  $t_{\text{air}}$  for samples with  $d_{\text{a-Si}} = 2$  nm, and  $\eta$  was improved from 5.3 to 6.6%, as shown in Table I. Meanwhile,  $V_{\text{OC}}$  decreases with increasing  $t_{\text{air}}$  for samples with  $d_{\text{a-Si}} = 5$  nm, while  $J_{\text{SC}}$  increases and  $\eta$  remains approximately constant at 4.7%. As shown in Table I,  $R_{\text{S}}$  increases with  $t_{\text{air}}$  regardless of  $d_{\text{a-Si}}$ , meaning that part of the a-Si and/or BaSi<sub>2</sub> layer became oxidized during the exposure of the samples to air. In our previous study, we found that  $\tau$  was improved when the oxygen composition in the region close to the BaSi<sub>2</sub> surface became large.<sup>22</sup> We thus speculate that the oxygen concentration became higher in the BaSi<sub>2</sub> region close to the a-Si/BaSi<sub>2</sub> interface for samples with  $d_{\text{a-Si}} = 2$  nm than for those with  $d_{\text{a-Si}} = 5$  nm when  $t_{\text{air}}$  was increased. This increase in oxygen concentration may lead to a reduction of the surface recombination and thereby a decrease in  $J_0$ . It was indeed found that  $J_0$  decreased by approximately 1/50 in sample B ( $d_{\text{a-Si}} = 2$  nm) compared with sample A ( $d_{\text{a-Si}} = 2$  nm) after 150 h, as shown in Table I. In an ideal case,  $V_{\text{OC}}$  is given by

$$V_{\text{OC}} \approx \frac{k_{\text{B}}T}{q} \exp\left(\frac{J_{\text{SC}}}{J_0}\right). \quad (2)$$

It is therefore reasonable that  $V_{\text{OC}}$  becomes larger by increasing  $t_{\text{air}}$ . Meanwhile,  $J_{\text{SC}}$  increased but the fill factor ( $FF$ ) and  $V_{\text{OC}}$  decreased with increasing  $t_{\text{air}}$  for samples at  $d_{\text{a-Si}} = 5$  nm because  $J_0$  increased by more than 20 times and  $R_{\text{S}}$  increased. The increase of  $J_{\text{SC}}$  likely arises from the decrease of absorption in the a-Si layer owing to its partial oxidation, as will be discussed later. Although the mechanism behind the large increase of  $J_0$  at  $t_{\text{air}} = 150$  h is not clear at present, it is safe to state that  $\eta$  is improved with increasing  $t_{\text{air}}$  when the a-Si layer thickness is small.

We next discuss the influence of  $d_{\text{a-Si}}$  upon the solar cell performance when  $t_{\text{air}} = 150$  h. Figure 2 shows typical examples of the  $J$ - $V$  characteristics under AM1.5 illumination where  $d_{\text{a-Si}}$  varies as 0.7, 2, 3, and 5 nm. It can be seen that the  $J$ - $V$  curves significantly depend upon  $d_{\text{a-Si}}$  and, in particular, the sample with  $d_{\text{a-Si}} = 0.7$  nm exhibits  $J$ - $V$  characteristics similar to those obtained without an a-Si capping layer (results in Ref. 27). To understand this result, we compared the solar cell properties of p-BaSi<sub>2</sub>/n-Si heterojunction solar cells fabricated with varying  $d_{\text{a-Si}}$ , as shown in Fig. 3. Some variation can be expected in the solar cell parameters, but it can be seen in Fig. 3 that there is a definite dependence of the solar cell parameters upon  $d_{\text{a-Si}}$ . When the a-Si layer was thinnest (0.7 nm), the  $R_{\text{S}}$  value was relatively large. This large  $R_{\text{S}}$  value suggests that it is possible to oxidize BaSi<sub>2</sub> with a sufficiently thin a-Si layer,<sup>25</sup> which could also explain the similar  $J$ - $V$  curves

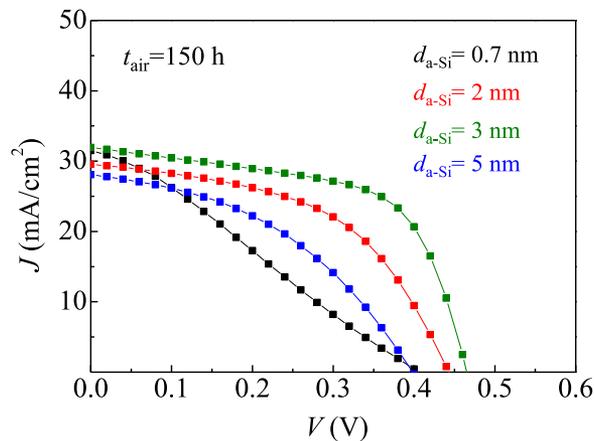


FIG. 2.  $J$ - $V$  characteristics under AM1.5 illumination measured for samples whose  $d_{\text{a-Si}}$  varies as 0.7, 2, 3, or 5 nm. The  $t_{\text{air}}$  was fixed at 150 h.

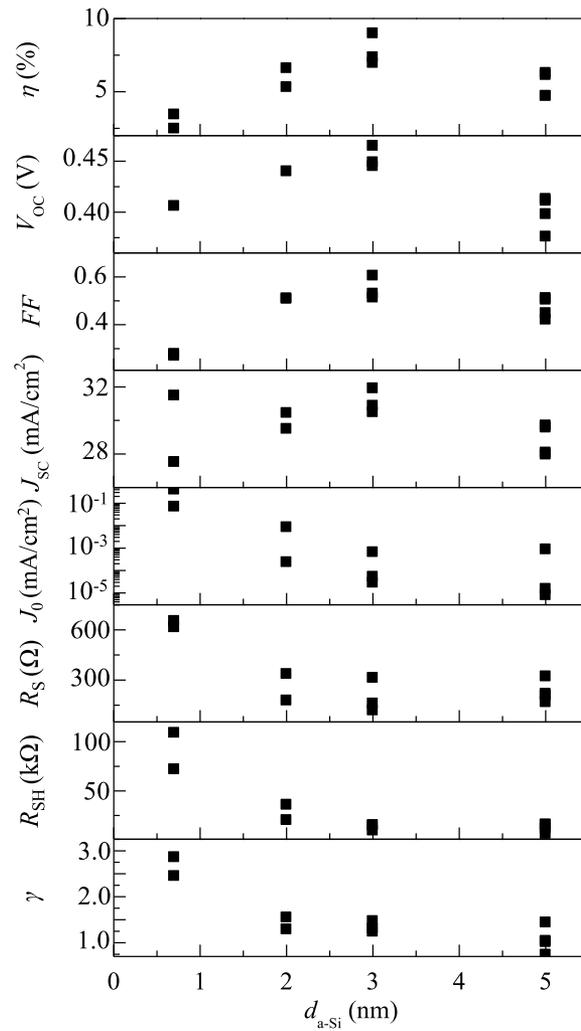


FIG. 3. Solar cell parameters of the p-BaSi<sub>2</sub>/n-Si heterojunction solar cells, plotted as a function of  $d_{a-Si}$ . The  $t_{air}$  was fixed at 150 h.

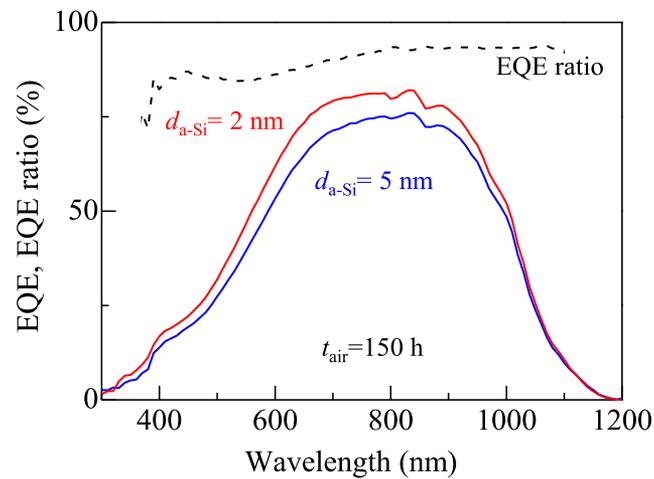


FIG. 4. EQE spectra for samples with  $d_{a-Si} = 2$  and 5 nm. The  $t_{air}$  was fixed at 150 h. The broken line shows the ratio of EQE at  $d_{a-Si} = 5$  nm to that at  $d_{a-Si} = 2$  nm.

for samples with  $d_{\text{a-Si}} = 0.7$  nm and those without a-Si capping layers. In addition, the  $\eta$  value reached a maximum of 9.0% when  $d_{\text{a-Si}}$  was 3 nm and exhibited a decrease at  $d_{\text{a-Si}} = 5$  nm, which was a trend echoed by the  $V_{\text{OC}}$  value. This is likely caused by a reduction of  $J_{\text{SC}}$  for  $d_{\text{a-Si}} = 5$  nm because other parameters such as  $J_0$ ,  $R_{\text{S}}$ ,  $R_{\text{SH}}$  and  $\gamma$  were seen to decrease with increasing  $d_{\text{a-Si}}$ , and to almost saturate at  $d_{\text{a-Si}} = 3$  nm. To clarify this reason, we compared the external quantum efficiency (EQE) spectra for the samples with varying  $d_{\text{a-Si}}$ .

Figure 4 shows the EQE spectra for typical samples with  $d_{\text{a-Si}} = 2$  and 5 nm, where the broken line shows the ratio of EQE at  $d_{\text{a-Si}} = 5$  nm to that at  $d_{\text{a-Si}} = 2$  nm. The EQE spectra are seen to decrease as  $d_{\text{a-Si}}$  increases, especially in the wavelength range shorter than  $\sim 730$  nm, which is equivalent to the bandgap energy of a-Si (1.7 eV). Therefore, it is reasonable to conclude that the reduction in  $J_{\text{SC}}$  arises from the light absorption within the a-Si layer. This result is similar to that reported in HIT solar cells,<sup>28</sup> and we thus attribute the reduction of  $J_{\text{SC}}$  to the same optical loss caused by the absorption in the a-Si layer. Therefore, a further increase in  $d_{\text{a-Si}}$  may be found to reduce  $\eta$ .

#### IV. CONCLUSION

We fabricated a-Si/B-doped p-BaSi<sub>2</sub>(20 nm)/n-Si heterojunction solar cells with various  $t_{\text{air}}$  and  $d_{\text{a-Si}}$  values, and investigated the influence of these varying parameters upon the solar cell performance. Solar cell parameters such as  $\eta$ ,  $V_{\text{OC}}$ ,  $J_{\text{SC}}$ ,  $J_0$ ,  $R_{\text{S}}$ ,  $\gamma$ , and  $R_{\text{SH}}$  were found to depend upon  $t_{\text{air}}$  and  $d_{\text{a-Si}}$ . The  $\eta$  value increased from 5.3 to 6.6% as  $t_{\text{air}}$  increased from 12 to 150 h for samples with  $d_{\text{a-Si}} = 2$  nm. When  $t_{\text{air}}$  was fixed at 150 h, the  $\eta$  value reached a maximum of 9.0% at  $d_{\text{a-Si}} = 3$  nm. These results reveal that the precise control of BaSi<sub>2</sub> oxidation can enhance  $\eta$  much further.

#### ACKNOWLEDGEMENTS

This work was supported in part by the Core Research for Evolutional Science and Technology (CREST) project of the Japan Science and Technology Agency (JST) and by a Grant-in-Aid for Scientific Research A (No. 15H02237) from the Japan Society for the Promotion of Science (JSPS). R.T. was financially supported by a Grant-in-Aid for JSPS Fellows (No. 15J02139).

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