A GAP THEOREM FOR COMPLETE FOUR-DIMENSIONAL MANIFOLDS WITH $\delta W^+ = 0$

By

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Abstract. Let M^4 be a complete noncompact oriented fourdimensional Riemannian manifold satisfying $\delta W^+ = 0$, where W^+ is the self-dual part of the Weyl curvature tensor. Suppose its scalar curvature is nonnegative and Sobolev's inequality holds. We show that if the L^2 norm of W^+ is sufficiently small, then $W^+ \equiv 0$.

1. Introduction

Let M^4 be a complete oriented four-dimensional Riemannian manifold. By the Hodge star operator *, the bundle of 2-forms Λ^2 splits into the sum of the bundle of self-dual and anti-self-dual 2-forms $\Lambda^2 = \Lambda_+^2 \oplus \Lambda_-^2$. According to this decomposition, Weyl curvature tensor W splits as $W = W^+ + W^-$. W^+ is called the self-dual part of the Weyl curvature tensor. We consider the equation $\delta W^+ = 0$. Here δ is the formal divergence defined as

$$\delta W^+(X_1,X_2,X_3) = -\sum_{i=1}^4 (
abla_{e_i}W^+)(e_i,X_1,X_2,X_3),$$

where $\{e_i\}$ is an orthonormal basis of TM with positive orientation. When the Ricci tensor is parallel, $\delta W^+ = 0$ ([2], [3]). Therefore, manifolds satisfying $\delta W^+ = 0$ are natural generalizations of Einstein manifolds or symmetric spaces.

Recently Gursky [4], Itoh-Satoh [5] proved L^2 or pointwise isolation theorem of W^+ for compact oriented four-dimensional Riemannian manifolds with $\delta W^+ = 0$.

In this note we give a gap theorem for noncompact manifolds.

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THEOREM 1.1. Let M^4 be a complete noncompact oriented four-dimensional Riemannian manifold with $\delta W^+ = 0$. Suppose its scalar curvature is nonnegative and Sobolev's inequality holds on M^4 . Then there is a constant C > 0 depending only on the Sobolev constant such that if $\int_M |W^+|^2 dv < C$, then $W^+ \equiv 0$.

2. Proof

The method of proof is a standard way in proving such an isolation theorem (for example, see [1] for minimal submanifold case, [6] for harmonic map case).

By the Weitzenböck formula, we have (cf. (3.11) and (3.12) in [4])

(1)
$$\Delta |W^+|^2 \ge 2|\nabla W^+|^2 + R|W^+|^2 - \sqrt{6}|W^+|^3,$$

where R is the scalar curvature of M. By using the following Kato's inequality ((2.1) in [4])

$$|\nabla|W^+| \le \sqrt{\frac{3}{5}} |\nabla W^+|$$

and the assumption $R \ge 0$ to (1), we obtain

$$\frac{2}{3} |\nabla| \mathcal{W}^+| |^2 \le |\mathcal{W}^+|\Delta| \mathcal{W}^+| + \frac{\sqrt{6}}{2} |\mathcal{W}^+|^3.$$

For simplicity we set $a = |W^+|$ and rewrite the above inequality as

$$\frac{2}{3}|\nabla a|^2 \le a\Delta a + \frac{\sqrt{6}}{2}a^3.$$

Let $\lambda \in C_0^1(M)$. Multiplying λ^2 and integrating over M, we get

(2)
$$\frac{2}{3}\int \lambda^2 |\nabla a|^2 \leq \int \lambda^2 a \Delta a + \frac{\sqrt{6}}{2}\int \lambda^2 a^3 dx$$

Since

$$\int 2\lambda a \nabla \lambda \cdot \nabla a + \int \lambda^2 |\nabla a|^2 + \int \lambda^2 a \Delta a = 0,$$

we get from (2)

(3)
$$\frac{5}{3}\int \lambda^2 |\nabla a|^2 + 2\int \lambda a \nabla \lambda \cdot \nabla a \leq \frac{\sqrt{6}}{2}\int \lambda^2 a^3.$$

We have

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$$2\int |\lambda a \nabla \lambda \cdot \nabla a| \leq \varepsilon \int \lambda^2 |\nabla a|^2 + \frac{1}{\varepsilon} \int a^2 |\nabla \lambda|^2$$

for any $\varepsilon > 0$. Plugging this into (3), we get

(4)
$$\left(\frac{5}{3}-\varepsilon\right)\int\lambda^2|\nabla a|^2\leq\frac{\sqrt{6}}{2}\int\lambda^2a^3+\frac{1}{\varepsilon}\int a^2|\nabla\lambda|^2.$$

On the other hand we have Sobolev's inequality

$$\left(\int |f|^{4/3}\right)^{3/4} \leq C_S \int |\nabla f| \quad \text{for } \forall f \in C_0^1(M).$$

Substituting f^3 into f, we obtain

$$\left(\int |f|^4\right)^{1/4} \leq 3C_{\mathcal{S}}\left(\int |\nabla f|^2\right)^{1/2}.$$

We apply this inequality for λa and get

(5)
$$\left(\int \lambda^4 a^4\right)^{1/2} \le 18C_S^2 \left\{\int \lambda^2 |\nabla a|^2 + \int a^2 |\nabla \lambda|^2\right\}.$$

From (4) and (5), we get

(6)
$$\left(\int \lambda^4 a^4\right)^{1/2} \leq \frac{18C_s^2}{\frac{5}{3} - \varepsilon} \left\{ \frac{\sqrt{6}}{2} \int \lambda^2 a^3 + \frac{1}{\varepsilon} \int a^2 |\nabla \lambda|^2 \right\} + 18C_s^2 \int a^2 |\nabla \lambda|^2.$$

Applying Hölder's inequality

$$\int \lambda^2 a^3 \le \left(\int \lambda^4 a^4 \right)^{1/2} \left(\int a^2 \right)^{1/2}$$

to (6), we get

$$\left\{1 - \frac{9\sqrt{6}}{\frac{5}{3} - \varepsilon} C_S^2 \left(\int a^2\right)^{1/2}\right\} \left(\int \lambda^4 a^4\right)^{1/2} \le \left\{\frac{18C_S^2}{\varepsilon(\frac{5}{3} - \varepsilon)} + 18C_S^2\right\} \int a^2 |\nabla \lambda|^2.$$

If we have

(7)
$$1 - \frac{9\sqrt{6}}{\frac{5}{3} - \varepsilon} C_s^2 \left(\int a^2 \right)^{1/2} > 0,$$

then by using a standard cut-off function argument, we conclude $a = |W^+| \equiv 0$. Since we can choose ε arbitrarily small, (7) is satisfied if

$$\int_{M} |W^{+}|^{2} < \frac{25}{4374} \cdot \frac{1}{C_{S}^{4}}.$$

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