

Central peak and acoustic anomalies of the relaxor ferroelectric lead magnesium niobate-lead titanate single crystal studied by the micro-Brillouin scattering

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Quasielastic central peak and acoustic anomalies in a [001] lead magnesium niobate-lead titanate single crystal have been investigated by the high-resolution micro-Brillouin scattering. Two acoustic anomalies were observed at temperatures, $T_{C-T} \sim 415$ K and $T_{T-R} \sim 348$ K. The damping showed a significant dispersion at $T \sim 550$ K, associated with the enhanced dynamics of local strain fields created by the polar microregions. The temperature evolution of the quasielastic central peak exhibited anomalies in the overall intensity and relaxation time that were associated with the line shape interferences of the central peak and the acoustic mode. © 2007 American Institute of Physics. [DOI: 10.1063/1.2768306]

The complex perovskite relaxor ferroelectric (FE) system $\text{Pb}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3$ (PMN) and its solid solutions with PbTiO_3 (PT) such as, $(1-x)\text{Pb}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3-x\text{PbTiO}_3$ (PMN- x PT) and $(1-x)\text{Pb}(\text{Zn}_{1/3}\text{Nb}_{2/3})\text{O}_3-x\text{PbTiO}_3$ (PZN- x PT) have been a subject of extensive investigations owing to both their exceptionally strong piezoelectric properties and quest for a comprehensive understanding of the nature of relaxor FE materials.^{1,2} In the phase diagram, compositions of the most practical interest lie in the low symmetry monoclinic region called the morphotropic phase boundary (MPB); separating rhombohedral ($R3m$) and tetragonal ($P4mm$) FE phases. Relaxor FE single crystals near this boundary exhibit unusual high piezoelectric constant, electromechanical coupling coefficient, and ultrahigh strains^{3,4} when poled along the pseudocubic [001] direction which is different from the polarization direction ([111]).

The most probable scenario toward understanding these extraordinary piezoelectric properties of the relaxor FE state stems in the well established existence of unique entities the so-called polar microregions (PMRs) with randomly oriented local polarization, appearing at/above Burns temperature T_B (Ref. 5) (~ 620 K for PMN) in a disordered (site and/or charge) matrix that give rise to local quenched random fields, and hence transition to a long range ordered state is prevented. The relaxor FE materials are also fascinating because they resemble glasses in the sense that they exhibit aging, memory, rejuvenation,⁶⁻⁹ and Vogel-Fulcher-type relaxation^{10,11} phenomena. However, in contrast to glasses, these characteristics have been observed in relaxors not only in the frozen regime but also even above the susceptibility maximum temperature T_m (~ 265 K for PMN). At the same time the relaxor state can be switched to FE long range ordered state by applying sufficiently high electric fields, adding PT (a classic soft mode FE) to varying concentrations, and exhibiting the presence of some criticality. This has led to

recent findings that the presence of critical behavior in the electric field-temperature-composition phase diagram of PMN- x PT (and possibly in other relaxors) is at the heart of this giant electromechanical response near the MPB.¹² Coupling between the optical and acoustic phonons in PMN (Ref. 13) and PMRs and TA phonons in PZN- x PT (Ref. 14) has been associated with the dynamics of ion off-centering and appearance of PMRs, respectively. In the Raman study of the PZN crystal, both relaxation and coupled phonon dynamics have been identified.¹⁵ Recent light scattering studies on PZN- x PT crystal have indicated the existence of an intermediate temperature at which the coupling between the PMRs and acoustic modes becomes more visible.¹⁶ In this letter we present a direct evidence of mode coupling that gives rise to line shape anomalies in the light scattering spectra of the PMN-33%PT [001]-oriented single crystal using the high-resolution micro-Brillouin scattering.

PMN-33%PT single crystal, grown by the Bridgman method, was put inside a cryostat cell (THMS 600) placed on the stage of an optical microscope (Olympus BH-2). The phonon modes were excited by a single frequency Ar⁺-ion laser at a wavelength of 514.5 nm and a power of ~ 100 mW. The micro-Brillouin light scattering experiments were performed in the backward scattering geometry by using a six pass tandem Fabry-Pérot interferometer from room temperature to $T \sim 650$ K with a stability of ± 0.1 K and an accuracy of ± 1 K. For accurate detection of the acoustic anomaly the Brillouin spectra were recorded by employing a smaller free spectral range (FSR) of 30 GHz while a relatively larger FSR of the order of 100 GHz was used to observe the quasielastic central peak (CP).

Figure 1 shows the typical Brillouin spectra of PMN-33%PT [001] single crystal at some selected temperatures. These spectra were measured in the scattering geometry, $c(a, a+b)\bar{c}$, where a , b , and c represent [100], [010], and [001] directions, respectively. One longitudinal acoustic (LA) mode and an additional transverse acoustic (TA) mode coupled with the quasielastic CP were observed in this geometry. The Brillouin frequency shift (ν_{LA}) and the hyper-

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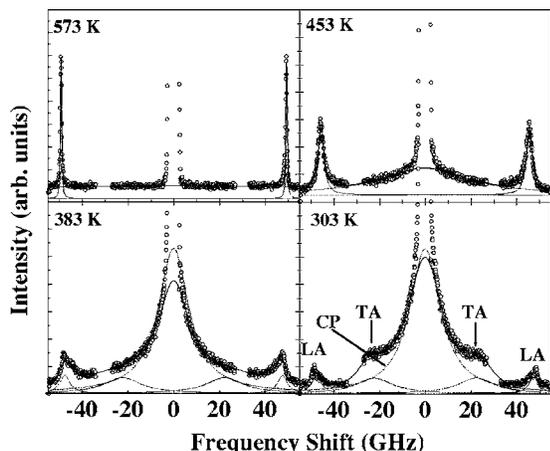


FIG. 1. Brillouin spectra of a PMN-33%PT single crystal at some selected temperatures. The circles show the experimental data and solid line between the circles is the fitting result.

sonic damping (Γ_{LA}) of the LA mode obtained from these spectra have been plotted in Fig. 2 as a function of temperature. The acoustic velocities of LA and TA modes are controlled by the respective elastic constants, e.g., c_{11} and c_{44} , in the cubic phase. The measured values of c_{11} and c_{44} for $T \geq T_{C-T}$, say, at ~ 453 K, were 163 and 77 GPa, respectively. The existence of different acoustic anomalies can be clearly differentiated from these data. It is also important to mention that by analyzing the spectra presented in Fig. 1 with the damped harmonic oscillator (DHO) couple mode scheme (described below) or uncoupled DHO or Voigt function reproduced similar results. The reason lies in the fact that LA mode is hardly affected by the CP and/or the TA mode.

It can be observed from Fig. 2 that with decreasing temperature, Γ deviates from linearity and begins to increase at $T \leq 550$ K, whereas ν_{LA} shows softening comparatively at a much higher temperature (beyond the measured temperature). It is remarkable to note that this dispersion in Γ and ν_{LA} exhibits similar temperature dependencies to that of PLZT relaxor FE ceramics.¹⁰ The two subsequent acoustic anomalies occurring at temperatures $T_{C-T} \sim 415$ K and $T_{T-R} \sim 348$ K show the paraelectric (PE) cubic-tetragonal and then tetragonal-rhombohedral FE phase transformations, respectively. This depicts that the structural changes have sig-

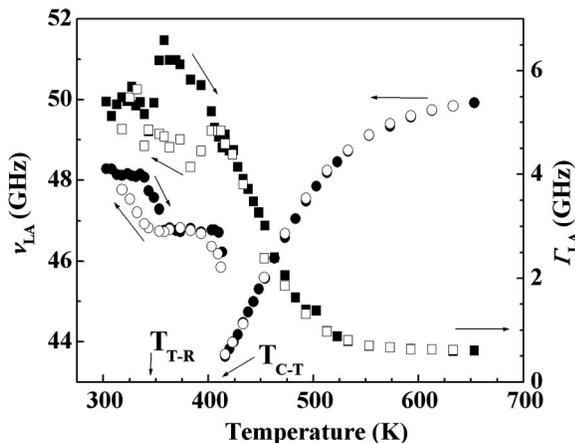


FIG. 2. Temperature dependencies of the Brillouin frequency shift (ν_{LA}) (circles) and damping factor (Γ) (squares) of the longitudinal acoustic mode of a PMN-33%PT single crystal.

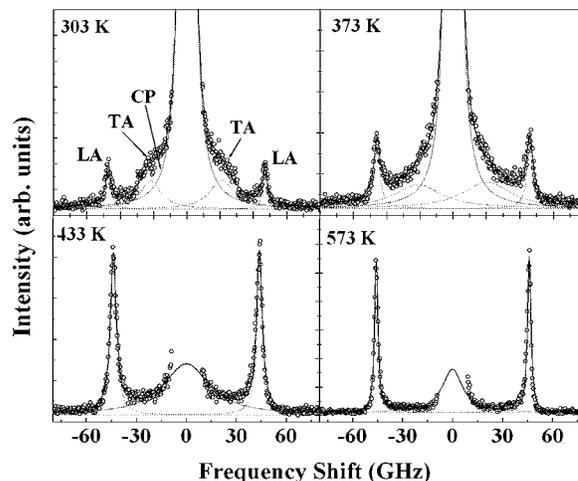


FIG. 3. Brillouin spectra of PMN-33%PT crystal measured with a FSR of 100 GHz in the rhombohedral (303 K), tetragonal (373 K), cubic (433 K), and cubic (573 K) with no CP, regions, respectively.

nificant influence on the elastic behavior of the crystal. The gradual softening of ν_{LA} and rise in Γ can be attributed to the appearance (and subsequent growth) of the PMRs at/above the Burns temperature T_B , far from T_m , where the optical index of refraction deviates from the high temperature linear behavior.⁵

The temperature dependent Brillouin spectra of the CP in the temperature regimes of interest are shown in Fig. 3. From these spectra two temperature regions can be identified: (i) $T \leq T_{C-T}$ and (ii) $T \geq T_{C-T}$, where the spectra consists of LA mode and TA mode coupled with the CP, and LA mode and the CP, respectively. Accordingly, we have analyzed the spectra for $T \leq T_{C-T}$ by a damped harmonic oscillator (describing LA mode, $\omega_{LA} = 2\pi\nu_{LA}$),

$$\chi_1(\omega, T) \propto [\omega_{LA}^2(T) - \omega^2 - i\omega\Gamma_{LA}(T)]^{-1},$$

and a DHO (describing TA mode) coupled with the Debye-type CP by the response function with relaxing self-energies related to the imaginary part of the dynamic susceptibility,¹⁷

$$\chi_2(\omega, T) \propto \left[\omega_{TA}^2(T) - \omega^2 - i\omega\Gamma_{TA}(T) - \frac{\delta^2(T)}{1 - i\omega\tau} \right]^{-1},$$

where τ is the Debye relaxation time and δ^2 is the coupling constant, i.e., the measure of interaction strength between the quasielastic CP and the coupled TA mode of frequency $\omega_{TA} (= 2\pi\nu_{TA})$ and damping factor Γ_{TA} .

The TA mode became very weak approaching $T \sim 370$ K and ultimately disappeared in the vicinity of T_{C-T} . Damping (Γ_{TA}) and frequency shift (ω_{TA}) of the TA mode remained unchanged for $T \leq T_{T-R}$ and were of the order of ~ 4.5 and ~ 163 GHz, respectively; however, the same could not be obtained reliably for $T \geq T_{T-R}$ due to complete merging of TA mode into the CP. In the PE phase ($T \geq T_{C-T}$) the fitting process was carried out with an uncoupled CP and the LA mode. It is important to note that the repeated fitting process yielded reproducible behavior of the fitting parameters in the two temperature regimes.

Overall intensity factor and relaxation time (τ) of CP as obtained from the coupled/uncoupled mode analysis of the measured spectra are shown in Fig. 4. Several things can be inferred from these data. (i) With decreasing temperature, relaxation time of CP increases approaching a sharp maxi-

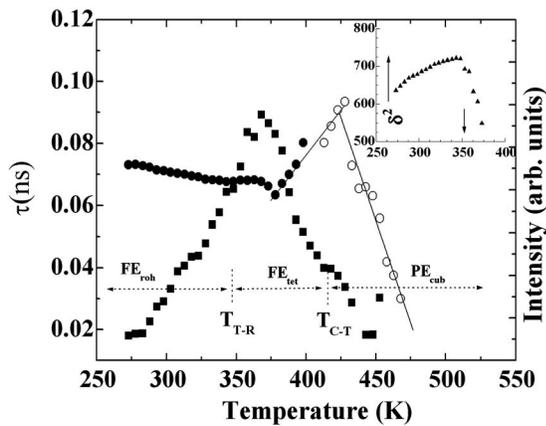


FIG. 4. Temperature dependencies of the overall intensity factor (filled squares) and relaxation time (τ) (filled circles are for coupled mode and open circles for uncoupled mode analysis) of the quasielastic central peak. Solid lines are mere guides for the eye. The inset shows the coupling constant δ^2 as a function of temperature.

imum around T_{C-T} , where overall intensity factor gives a small inflection; it then decreases showing a dip at $T \sim 378$ K with a corresponding maximum in the overall intensity factor at $T \sim 370$ K (slowing down of relaxation dynamics and growth of CP). (ii) With further cooling the overall intensity factor lowers and τ shows only minor rise. (iii) The coupling constant δ^2 is minimum, where τ has a dip ($T \sim 378$ K) and becomes highest at T_{T-R} , it then softens slowly with further cooling. It is also noted that the TA mode does affect the CP line shape but has no appreciable contribution to the LA mode line profile; hence the two acoustic anomalies remain unaffected by its appearance. The anomalous behavior of CP parameters (reflected in overall intensity factor and τ) for $T \leq T_{C-T}$ demonstrate that it comes from the line shape interferences due to interaction of TA phonon with the quasielastic CP. The reason behind this argument lies in the fact that when we analyzed the same data with uncoupled DHO and Debye relaxation mode, the anomalies in the τ decreased but the intensity of CP showed two peaks. The PMR dynamics is no more effective in this temperature range because the crystal has already been transformed to long-range ordered state. The other factor that may influence CP dynamics in the present case is realignment/movement of domain walls in the FE matrix resulting from the reorientation of polarization direction from [001] to [111].¹⁸ This explanation is supported by similar remarks about PZN-8%PT (Ref. 19) and PMN-60%PT (Ref. 20) crystals, where the interference of acoustic and optical phonon branches gives rise to “waterfall effect” observed in PMN and PZN relaxors. Fourier transform infrared spectra²¹ between 20 and 900 K revealed the existence of an underdamped soft mode in PMN crystal. The appearance of CP below T_B and its slowing down character on cooling down to frequencies below the microwave range has led to the conclusion that the waterfall effect may be associated with the overlapping of the soft mode with the CP.²¹ From the temperature dependence of the acoustic anomalies in PMN-33%PT crystal, it can be argued

that in this crystal, a certain temperature point ($T \sim 550$ K) below T_B indeed exists where the dynamics of PMRs produce sufficient local strain fields in the crystal that result in appreciable dispersion in frequency and damping of the LA phonon mode. This observation is consistent with our previous findings on PLZT (Ref. 22) and recent light scattering studies on PZN-9%PT crystal.¹⁶ It is about 100 K below this temperature where the quasielastic CP becomes more prominent as observed in PLZT relaxor²³ and in the present case also. A similar trend of the CP has also been reported for PMN-32%PT crystal by neutron scattering experiments²⁴ and associated with the dynamical evolution of PMRs.

In summary, high-resolution micro-Brillouin scattering technique was applied to investigate the acoustic anomalies and quasielastic CP in a [001]-oriented PMN-33%PT single crystal. Two acoustic anomalies, PE cubic-FE tetragonal ($T_{C-T} \sim 415$ K) and FE tetragonal-FE rhombohedral ($T_{T-R} \sim 348$ K) symmetries, were observed. The LA mode data also exhibited the existence of a temperature point ($T \sim 550$ K), where the dynamics of PMRs becomes more prominent in PMN-33%PT crystal and possibly in other relaxor FE materials. From the analysis of Brillouin spectra with larger FSR coupling between the acoustic mode and the quasielastic, CP was evident which gives rise to temperature dependent line shape anomalies in the relaxation time and overall intensity factor of the CP. The detailed analysis of the CP data will be published elsewhere.

¹L. E. Cross, *Ferroelectrics* **76**, 241 (1987).

²Z.-G. Ye, *Key Eng. Mater.* **155–156**, 81 (1998).

³S.-E. Park and T. R. Shrout, *J. Appl. Phys.* **82**, 1804 (1997).

⁴J. Kuwata, K. Uchino, and S. Nomura, *Jpn. J. Appl. Phys., Part 1* **21**, 1298 (1982).

⁵G. Burns and F. H. Docal, *Phys. Rev. B* **28**, 2527 (1983).

⁶G. Shabbir, J.-H. Ko, and S. Kojima, *Appl. Phys. Lett.* **86**, 012908 (2005).

⁷G. Shabbir and S. Kojima, *J. Korean Phys. Soc.* **46**, 128 (2005).

⁸E. V. Colla, L. K. Chao, and M. B. Weissman, *Phys. Rev. B* **63**, 134107 (2001).

⁹F. Cordero, F. Craciun, A. Franco, D. Piazza, and C. Galassi, *Phys. Rev. Lett.* **93**, 097601 (2004).

¹⁰G. Shabbir and S. Kojima, *Europhys. Lett.* **63**, 388 (2003).

¹¹A. E. Glazounov and A. K. Tagantsev, *Appl. Phys. Lett.* **73**, 856 (1998).

¹²Z. Kutnjak, J. Petzelt, and R. Blinc, *Nature (London)* **441**, 956 (2006).

¹³O. Svitelskiy, J. Toulouse, and Z.-G. Ye, *Phys. Rev. B* **68**, 104107 (2003).

¹⁴D. La-Orautapong, J. Toulouse, Z.-G. Ye, W. Chen, R. Erwin, and J. L. Robertson, *Phys. Rev. B* **67**, 134110 (2003).

¹⁵J. Toulouse, F. Jiang, O. Svitelskiy, W. Chen, and Z.-G. Ye, *Phys. Rev. B* **72**, 184106 (2005).

¹⁶J.-H. Ko, D. H. Kim, and S. Kojima, *Appl. Phys. Lett.* **90**, 112904 (2007).

¹⁷S. M. Shapiro, J. D. Axe, G. Shirane, and A. Riste, *Phys. Rev. B* **6**, 4332 (1972).

¹⁸G. Shabbir, S. Kojima, and C. Feng, *J. Appl. Phys.* **100**, 064107 (2006).

¹⁹J. Hlinka, S. Kamba, J. Petzelt, J. Kulda, C. A. Randall, and S. J. Zhang, *Phys. Rev. Lett.* **91**, 107602 (2003).

²⁰C. Stock, D. Ellis, I. P. Swainson, G. Xu, H. Hiraka, Z. Zhong, H. Luo, X. Zhao, D. Viehland, R. J. Birgeneau, and G. Shirane, *Phys. Rev. B* **73**, 064107 (2006).

²¹S. Kamba, M. Kempa, V. Bovtun, J. Petzelt, K. Brinkman, and N. Setter, *J. Phys.: Condens. Matter* **17**, 3965 (2005).

²²G. Shabbir and S. Kojima, *J. Phys.: Condens. Matter* **15**, 7717 (2003).

²³G. Shabbir and S. Kojima, *Ferroelectrics* **303**, 167 (2004).

²⁴S. N. Gvasaliya, B. Roessli, R. A. Cowley, S. Kojima, and S. G. Lushnikov, *J. Phys.: Condens. Matter* **19**, 016219 (2007).