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Measurements of the chiral magnetic effect with background isolation in 200 GeV Au+Au collisions at STAR

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Abstract

Using two novel methods, pair invariant mass (m_{inv}) and comparative measurements with respect to reaction plane (Ψ_{RP}) and participant plane (Ψ_{PP}), we isolate the possible chiral magnetic effect (CME) from backgrounds in 200 GeV Au+Au collisions at STAR. The invariant mass method identifies the resonance background contributions, coupled with the elliptic flow (v_2), to the charge correlator CME observable ($\Delta\gamma$). At high mass ($m_{inv} > 1.5$ GeV/ c^2) where resonance contribution is small, we obtain the average $\Delta\gamma$ magnitude. In the low mass region ($m_{inv} < 1.5$ GeV/ c^2), resonance peaks are observed in $\Delta\gamma(m_{inv})$. An event shape engineering (ESE) method is used to model the background shape in m_{inv} to extract the potential CME signal at low m_{inv} . In the comparative method, the Ψ_{RP} is assessed by spectator neutrons measured by the ZDCs, and the Ψ_{PP} by the 2nd-harmonic event plane measured by the TPC. The v_2 is stronger along Ψ_{PP} and weaker along Ψ_{RP} ; in contrast, the magnetic field, mainly from spectator protons, is weaker along Ψ_{PP} and stronger along Ψ_{RP} . As a result, the $\Delta\gamma$ measured with respect to Ψ_{RP} and Ψ_{PP} contain different amounts of CME and background, and can thus determine these two contributions. It is found that the possible CME signals with background isolation by these two novel methods are small, on the order of a few percent of the inclusive $\Delta\gamma$ measurements.

Keywords: QCD, heavy-ion collisions, chiral magnetic effect, invariant mass, reaction plane, participant plane

1. Introduction

Quark interactions with topological gluon fields can induce chirality imbalance and local parity violation in quantum chromodynamics (QCD) [1]. In relativistic heavy-ion collisions, this can lead to observable electric charge separation along the strong magnetic field, \vec{B} , produced mainly by spectator protons [2]. This is called the chiral magnetic effect (CME). The commonly used observable to search for the CME-induced charge separation is the three-point azimuthal correlator difference [3], $\Delta\gamma \equiv \gamma_{OS} - \gamma_{SS}$; $\gamma = \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle \approx \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_c) \rangle / v_2$, where ϕ_α and ϕ_β are the azimuthal angles of two charged particles, of opposite electric charge sign (OS) or same sign (SS), and Ψ_{RP} is that of the reaction plane (span by the impact parameter direction and the beam) to which \vec{B} is perpendicular on average. The latter is often surrogated by the azimuthal angle of a third particle, ϕ_c , with a resolution correction factor given by the particle's elliptic anisotropy (v_2). Significant $\Delta\gamma$ has indeed been observed in heavy-ion collisions [4]. One of the difficulties in its CME interpretation is a major background contribution arising from the coupling of resonance decay correlations and the v_2 stemming from the participant geometry [5].

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2. Invariant mass dependence of the $\Delta\gamma$ correlator

The main backgrounds for the $\Delta\gamma$ are from the resonance decays coupled with v_2 [5]. A new analysis approach exploiting the particle pair invariant mass, m_{inv} , to identify the backgrounds and, hence, to extract the possible CME signal is proposed [6]. Figure 1 shows the m_{inv} dependences of (left panel) the relative excess of OS over SS charged π pairs, $r = (N_{OS} - N_{SS})/N_{OS}$, and (middle panel) the three-point correlator difference, $\Delta\gamma = \gamma_{OS} - \gamma_{SS}$. A lower cut on m_{inv} was used to suppress the resonance contributions. Figure 1 (right panel) shows the inclusive $\Delta\gamma$ over all mass (black) and the $\Delta\gamma$ at $m_{inv} > 1.5$ GeV/ c^2 (red) as a function of centrality in Au+Au collisions at 200 GeV in Run-11. In 20-50% collisions centrality, combining results from Run-11 (~0.5 billion minimum-bias events, year 2011), Run-14 (~0.8 billion, year 2014) and Run-16 (~1.2 billion, year 2016), the $\Delta\gamma$ at $m_{inv} > 1.5$ GeV/ c^2 is $(5 \pm 2 \pm 4)\%$ of the inclusive $\Delta\gamma$. The systematic uncertainty (the second error) is currently estimated by the differences among the three runs.

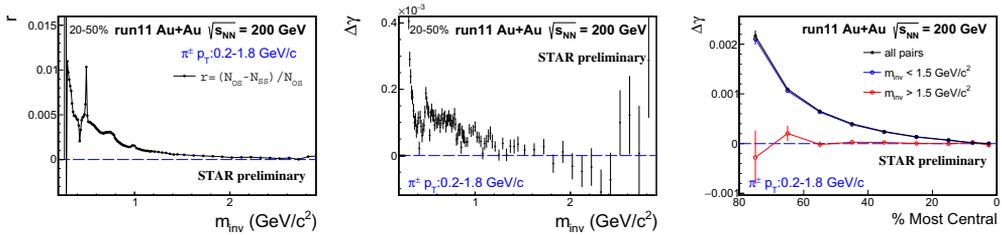


Fig. 1. Pion pair invariant mass, m_{inv} , dependences of (left panel) the relative excess of opposite-sign (OS) over same-sign (SS) pion pairs, $r = (N_{OS} - N_{SS})/N_{OS}$, and (middle panel) the three-point correlator difference, $\Delta\gamma = \gamma_{OS} - \gamma_{SS}$. (Right panel) The inclusive $\Delta\gamma$ over all mass (black points) and the $\Delta\gamma$ at $m_{inv} > 1.5$ GeV/ c^2 (red points) as a function of centrality. The pions are identified by STAR TPC and TOF with $0.2 < p_T < 1.8$ GeV/ c . Data from Run-11. Errors are statistical.

The CME is expected to be a low p_T phenomenon [7]; its contribution to high mass may be small. To extract CME at low mass, resonance contributions need to be subtracted. The inclusive $\Delta\gamma$ can be expressed as $\Delta\gamma(m_{inv}) = r(m_{inv}) \times \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{reso.}) \rangle \times v_{2,reso.} + \Delta\gamma_{CME}$ [6]. The event shape engineering (ESE) [8] method provides a tool to select events with different v_2 values by cutting on the q_2 ($q_2 = |\sum e^{i2\phi}|/\sqrt{N}$). The difference of the $\Delta\gamma(m_{inv})$ from different q_2 classes can be regarded as the background $\Delta\gamma(m_{inv})$ shape [9], assuming the CME are the same for events from different q_2 classes.

Figure 2 shows the m_{inv} dependence of (left top panel) the $r = (N_{OS} - N_{SS})/N_{OS}$, and (left middle panel) the $\Delta\gamma_A$ and $\Delta\gamma_B$ from ESE selected event samples A (large 50% q_2) and B (small 50% q_2), respectively. The correlators are calculated by $\gamma = \cos(\phi_\alpha + \phi_\beta - 2\phi_c)/v_{2,c}$. The TPC full event is divided into east and west sub-events, with α, β and q_2 from one sub-event and c from the other. Figure 2 (left bottom panel) shows the inclusive (0-100% q_2) $\Delta\gamma$ compared with $\Delta\gamma_A - \Delta\gamma_B$. A linear function, $\Delta\gamma_A = b \times \Delta\gamma_B + (1 - b) \times \Delta\gamma_{CME}$, is used to extract the CME. Figure 2 (right panel) shows the fit result in Run-16. Combining Runs 11, 14 and 16, the fit parameter $\Delta\gamma_{CME}$ is $(2 \pm 4 \pm 6)\%$ of the inclusive $\Delta\gamma$ in 20-50% centrality Au+Au collisions at 200 GeV. The systematic uncertainty is currently estimated by the differences among the three runs.

3. $\Delta\gamma$ with respect to Ψ_{RP} (ZDC) and Ψ_{PP} (TPC)

The CME-driven charge separation is along the magnetic field direction (Ψ_B). The major background to the CME is related to the elliptic flow anisotropy (v_2), determined by the participant geometry. A novel idea of differential measurements with respect to the reaction plane (Ψ_{RP}) and participant plane (Ψ_{PP}) is proposed [10], where the Ψ_{RP} could be assessed by spectator neutrons measured by the zero-degree calorimeters (ZDC) [11]. The v_2 is stronger along Ψ_{PP} and weaker along Ψ_{RP} ; in contrast, the magnetic field, being mainly from spectator protons, is weaker along Ψ_{PP} and stronger along Ψ_{RP} . The $\Delta\gamma$ values measured with respect to Ψ_{RP} and Ψ_{PP} contain different amounts of CME and background, and can thus determine these

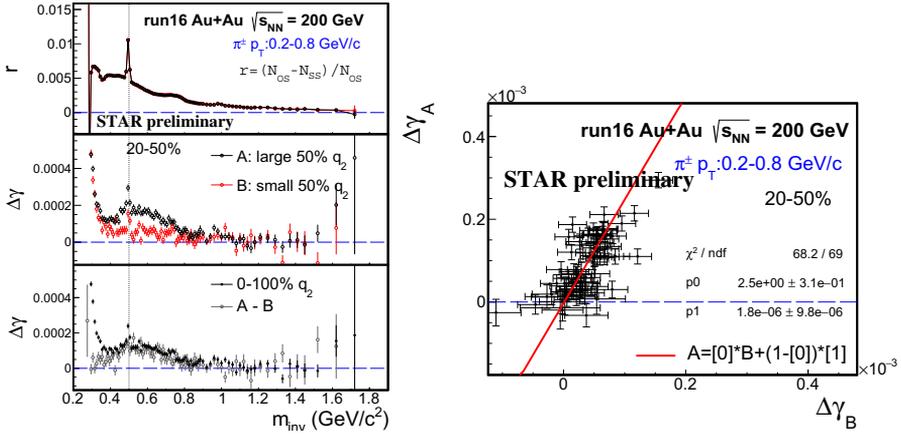


Fig. 2. Pion pair invariant mass, m_{inv} , dependences of (left top panel) the $r = (N_{OS} - N_{SS})/N_{OS}$, (left middle panel) the $\Delta\gamma_A$ and $\Delta\gamma_B$ from ESE selected event samples A (large 50% q_2) and B (small 50% q_2), and (left bottom panel) the inclusive (0-100% q_2) $\Delta\gamma$ compared with $\Delta\gamma_A - \Delta\gamma_B$. (Right panel) $\Delta\gamma_A$ vs. $\Delta\gamma_B$ fitted by a linear function (see text for explanations). The pions are identified by STAR TPC with $0.2 < p_T < 0.8$ GeV/c. Data from Run-16. Errors are statistical.

two contributions assuming that the CME is proportional to the magnetic field squared and the background is proportional to v_2 , as follows [10]:

$$\begin{aligned}
 \Delta\gamma\{\Psi_{\text{TPC}}\} &= \Delta\gamma_{\text{CME}}\{\Psi_{\text{TPC}}\} + \Delta\gamma_{\text{Bkg}}\{\Psi_{\text{TPC}}\}, & \Delta\gamma\{\Psi_{\text{ZDC}}\} &= \Delta\gamma_{\text{CME}}\{\Psi_{\text{ZDC}}\} + \Delta\gamma_{\text{Bkg}}\{\Psi_{\text{ZDC}}\}, \\
 \Delta\gamma_{\text{CME}}\{\Psi_{\text{TPC}}\} &= a\Delta\gamma_{\text{CME}}\{\Psi_{\text{ZDC}}\}, & \Delta\gamma_{\text{Bkg}}\{\Psi_{\text{ZDC}}\} &= a\Delta\gamma_{\text{Bkg}}\{\Psi_{\text{TPC}}\}, \\
 a &= v_2\{\Psi_{\text{ZDC}}\}/v_2\{\Psi_{\text{TPC}}\}, & A &= \Delta\gamma\{\Psi_{\text{ZDC}}\}/\Delta\gamma\{\Psi_{\text{TPC}}\}, \\
 f_{\text{CME}}^{\text{EP}} &= \Delta\gamma_{\text{CME}}\{\Psi_{\text{TPC}}\}/\Delta\gamma\{\Psi_{\text{TPC}}\} = (A/a - 1)/(1/a^2 - 1).
 \end{aligned} \tag{1}$$

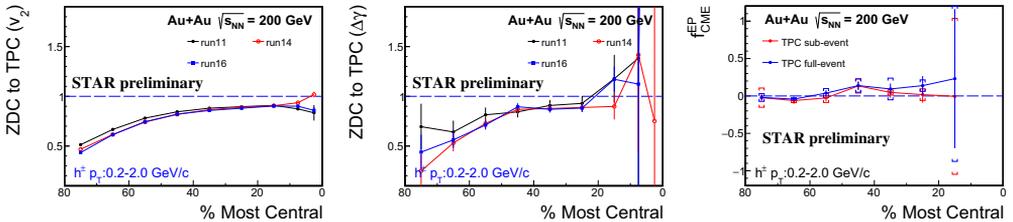


Fig. 3. The centrality dependences of the ratios of the charged hadron v_2 (left panel) and $\Delta\gamma$ (middle panel) measured with respect to the ZDC event plane to those with respect to the TPC event plane. (Right panel) The extracted $f_{\text{CME}}^{\text{EP}}$ as a function of collision centrality. Data from Runs 11, 14, and 16. Error bars are statistical errors. The horizontal caps on the right panel are systematic uncertainties.

Figure 3 shows the ratio of v_2 (left panel) measured with respect to the ZDC event plane and the v_2 with respect to the TPC event plane, $a = v_2\{\Psi_{\text{ZDC}}\}/v_2\{\Psi_{\text{TPC}}\}$ in Eq. (1), and that of $\Delta\gamma$ (middle panel), $A = \Delta\gamma\{\Psi_{\text{ZDC}}\}/\Delta\gamma\{\Psi_{\text{TPC}}\}$ in Eq. (1), as functions of collision centrality. To suppress the non-flow contributions in v_2 and $\Delta\gamma$ measurements, the TPC sub-event method is used, where each TPC event is divided into east and west sub-events, with the Ψ_{TPC} from one sub-event and the particles of interest from the other. Figure 3 (right panel) shows the extracted possible CME fraction ($f_{\text{CME}}^{\text{EP}}$) [10] as function of centrality. For comparison the results from TPC full-event method are also plotted. The extracted $f_{\text{CME}}^{\text{EP}}$ (combined from Runs 11, 14 and 16) are $(9 \pm 4 \pm 7)\%$ and $(12 \pm 4 \pm 11)\%$ from the TPC sub-event and full-event methods in 20-50% Au+Au collisions at 200 GeV, respectively. The systematic uncertainty is currently estimated by the differences among the three runs.

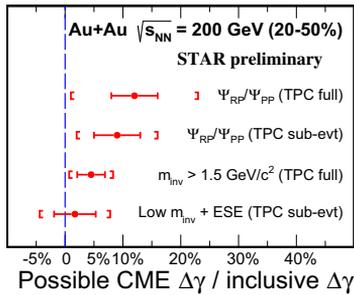


Fig. 4. The possible CME $\Delta\gamma$ over the inclusive $\Delta\gamma$ fraction from different analysis methods in middle central (20-50%) Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. Horizontal bars are statistical errors. Vertical caps are systematic uncertainties.

4. Summary

Charge separation measurements by the three-point azimuthal correlator ($\Delta\gamma$) are contaminated by major backgrounds arising from resonance decay correlations coupled with the elliptical anisotropy (v_2). To reduce/eliminate background contaminations, two novel methods are employed: the $\Delta\gamma$ correlator as a function of the particle pair invariant mass (m_{inv}) and the comparative $\Delta\gamma$ measurements with respect to Ψ_{RP} (estimated by ZDC) and Ψ_{PP} (estimated by TPC).

Resonance structures are observed in $\Delta\gamma$ as function of $\pi\text{-}\pi$ m_{inv} . A lower m_{inv} cut ($m_{inv} > 1.5$ GeV/ c^2) yields a $\Delta\gamma$ fraction of $(5 \pm 2 \pm 4)\%$ of the inclusive $\Delta\gamma$ measurement in 20-50% centrality Au+Au collisions at 200 GeV. In the low mass region, event shape engineering is used to determine the background shape in m_{inv} , and a linear fit to $\Delta\gamma(m_{inv})$ yields a possible CME signal of $(2 \pm 4 \pm 6)\%$ of the inclusive $\Delta\gamma$ measurement in 20-50% centrality Au+Au collisions at 200 GeV.

The $\Delta\gamma$ measurements with respect to Ψ_{RP} and Ψ_{PP} contain different amounts of CME and background. The v_2 is stronger along Ψ_{PP} and weaker along Ψ_{RP} ; and the magnetic field is weaker along Ψ_{PP} and stronger along Ψ_{RP} . By comparing the v_2 and $\Delta\gamma$ with respect to Ψ_{RP} and Ψ_{PP} , the extracted possible CME fractions are $(9 \pm 4 \pm 7)\%$ and $(12 \pm 4 \pm 11)\%$ from the TPC sub-event and full-event methods in 20-50% centrality Au+Au collisions at 200 GeV, respectively.

The extracted potential CME signal fractions (CME $\Delta\gamma$ over the inclusive $\Delta\gamma$) in middle central (20-50%) Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV are summarized in Fig. 4. These data-driven estimates indicate that the possible CME signal is small, within 1-2 σ from zero. Precision can be improved in the future with more Au+Au data and the new isobar run. Possible ZDC upgrades to achieve better Ψ_{RP} determination are being investigated.

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References

- [1] D. Kharzeev, R. D. Pisarski, M. H. G. Tytgat, Phys. Rev. Lett. 81 (1998) 512515.
- [2] K. Fukushima, D. E. Kharzeev, H. J. Warringa, Phys. Rev. D78 (2008) 074033.
- [3] S. A. Voloshin, Phys. Rev. C70 (2004) 057901.
- [4] D. E. Kharzeev, J. Liao, S. A. Voloshin, G. Wang, Prog. Part. Nucl. Phys. 88 (2016) 128.
- [5] F. Wang, Phys. Rev. C81 (2010) 064902; A. Bzdak, V. Koch, J. Liao, Phys. Rev. C81 (2010) 031901; S. Schlichting, S. Pratt, Phys. Rev. C83 (2011) 014913; L. Adamczyk et al. (STAR collaboration), Phys. Rev. C89 (2014) 044908; F. Wang, J. Zhao, Phys. Rev. C95 (2017) 051901(R).
- [6] J. Zhao, H. Li, F. Wang, arXiv:1705.05410; J. Zhao (for STAR), EPJ Web of Conf. 172 (2018) 01005, arXiv:1712.00394; *Id.*, Int. J. Mod. Phys. Conf. Ser. 46, 1860010 (2018), arXiv:1802.03283; H. Li et al., these proceedings, arXiv:1808.03210.
- [7] D. E. Kharzeev, L. D. McLerran, H. J. Warringa, Nucl. Phys. A803 (2008) 227253.
- [8] J. Schukraft, A. Timmins, S. A. Voloshin, Phys. Lett. B719 (2013) 394398.
- [9] J. Zhao, Int. J. Mod. Phys. A33 (2018) 1830010; J. Zhao, Z. Tu, F. Wang, arXiv:1807.05083.
- [10] H. Xu, J. Zhao, X. Wang, H. Li, Z. Lin, C. Shen, F. Wang, Chin. Phys. C42 (2018) 084103; H. Xu, X. Wang, H. Li, J. Zhao, Z. Lin, C. Shen, F. Wang, Phys. Rev. Lett. 121 (2018) 022301; H. Xu et al., these proceedings, arXiv:1808.00133.
- [11] C. Adler, H. Strobele, A. Denisov, E. Garcia, M. Murray, S. White, Nucl. Instrum. Meth. A461 (2001) 337340.