

# Dependence of crystal orientation in Al-induced crystallized poly-Si layers on SiO<sub>2</sub> insertion layer thickness

著者別名	都甲 薫, 末益 崇
journal or publication title	Journal of crystal growth
volume	356
page range	65-69
year	2012-10
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URL	<a href="http://hdl.handle.net/2241/117630">http://hdl.handle.net/2241/117630</a>

doi: 10.1016/j.jcrysgr.2012.07.015

1 **Dependence of crystal orientation in Al-induced crystallized poly-Si layers**  
2 **on SiO<sub>2</sub> insertion layer thickness**

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5 Atsushi Okada<sup>a</sup>, Kaoru Toko<sup>a</sup>, Kosuke O. Hara<sup>b</sup>, Noritaka Usami<sup>b,c</sup>, and Takashi Suemasu<sup>a,b</sup>

6 *<sup>a</sup>Institute of Applied Physics, University of Tsukuba, 1-1-1 Tennodai, Tsukuba, Ibaraki*  
7 *305-8573, Japan*

8 *<sup>b</sup>Japan Science and Technology Agency, CREST, 5, Sanbancho, Chiyoda-ku, Tokyo 102-0075,*  
9 *Japan*

10 *<sup>c</sup>Institute for Materials Research, Tohoku University, 2-1-1, Katahira, Aoba-ku, Miyagi*  
11 *980-8577, Japan*

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13 We have fabricated poly-Si thin films on fused silica substrates by the Al-induced  
14 crystallization (AIC) method with SiO<sub>2</sub> insertion layers of various thicknesses (0-20 nm). The  
15 growth morphologies of poly-Si layers were dramatically changed by the SiO<sub>2</sub> thickness, *i.e.*,  
16 thin layers (2 nm) provided high growth rates and (100) orientations, and thick layers (10 nm)  
17 provided low growth rates and (111) orientations. These results showed that the crystal  
18 orientation of AIC-Si significantly depends on the diffusion rate of Si atoms into the Al layer.

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20 PACS: 61.05.cp; 61.05.jh; 68.55.ag

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24 *Keywords:* A1. Growth models; A3. Solid phase epitaxy; B2. Semiconducting silicon; B3.

25 Solar cells

26

## 1. Introduction

The Al-induced crystallization (AIC) method is known as a method for obtaining thin poly-Si layers on glass substrates, where amorphous Si (a-Si) layers on Al are transformed into a crystalline phase via exchange between the Al and Si layers during annealing [1]. This method enables us to fabricate seed layers for epitaxially grown absorbers for thin-film solar cells on inexpensive SiO<sub>2</sub> substrates [2-7]. Conventional solid phase crystallization (SPC) requires temperatures higher than 600°C, and the Si grains obtained are as small as a few μm [8]. In the AIC method, Si crystallizes at relatively low temperatures in the range of 400-500°C, and large grains of tens of μm in size can be obtained even on non-lattice-matched substrates [9,10]. Therefore, the Si seed crystals fabricated by the AIC method seem to enable the growth of high-quality devices such as thin-film solar cells and thin-film transistors on low-cost glass substrates [11-15]. The crystal orientation is critical for device performances [16]. Thus, the control of crystal orientations in the poly-Si layers is required.

However, factors determining the preferential crystal orientation of AIC-Si layers remain an open question. Preferential (100) orientation of the AIC-Si layers has been reported [9,13,17]. On the other hand, growth of (111)-oriented Si layers has been also reported [15,18]. Kurosawa *et al.* showed that the preferential orientation of AIC-Si on fused silica can be controlled by changing the exposure time of Al in air, and they proposed

46 a model showing that the orientation of AIC-Si is determined by the crystal phase of the  
47 native Al oxide [19]. In contrast, Jung *et al.* considered that the preferential orientation  
48 depends on the annealing temperature, and proposed a model showing that the growth rate  
49 of AIC-Si determines the preferential orientation [20]. In this study, our aim was to  
50 investigate the influence of Si diffusion rate on the preferential orientation of AIC-Si. We  
51 also aimed to develop a method to control the crystal orientation of AIC-Si. For this purpose,  
52 we introduced SiO<sub>2</sub> layers with various thicknesses as intermediate layers between the a-Si  
53 and Al layers and carried out the AIC method. The SiO<sub>2</sub> layer thickness is easier to control  
54 compared to the thickness of a native Al oxide. In addition, the effect of the intermediate  
55 layer crystallinity on the AIC-Si is negligible when we discuss the crystal orientation of  
56 AIC-Si layers.

57

## 58 **2. Experimental procedures**

59 Fused silica wafers were used as substrates in this work. A 100-nm-thick Al film was  
60 sputtered at room temperature (RT) on the substrate. Next, an amorphous SiO<sub>2</sub> layer was  
61 subsequently sputtered at RT on the Al layer, followed by sputtering a 100-nm-thick a-Si  
62 film at RT without breaking the vacuum. The thickness of the SiO<sub>2</sub> layer was varied from 0  
63 to 20 nm. All the depositions were carried out by radio-frequency (RF) magnetron  
64 sputtering. For comparison, conventional AIC method was also performed, that is,

65 deposition of Al layers, followed by breaking the vacuum to form a native Al layer for 48 h,  
66 and the subsequent deposition of a-Si layers. The argon pressure during the sputtering was  
67 0.2 Pa. The RF power was set at 100 W. The AIC method was carried out by annealing the  
68 samples in N<sub>2</sub> atmosphere at 500 °C for 10 h. Sample preparation was summarized in Table  
69 1. The surface morphologies of some of the samples were observed during annealing by  
70 optical microscopy. After annealing, the Al and oxide layers were etched away using a HF  
71 solution (HF: 2%). The crystal orientation of AIC-Si was characterized by electron  
72 backscatter diffraction (EBSD) measurement.

73

### 74 **3. Results and discussion**

75 The time evolution of optical microscope images for sample D, prepared with a  
76 2-nm-thick SiO<sub>2</sub> intermediate layer, is shown in Figs. 1(b)-1(e), and those for sample H,  
77 prepared with a 10-nm-thick SiO<sub>2</sub> layer, are shown in Figs. 1(f)-1(i). Figure 1(a) shows the  
78 expected schematic cross-sectional diagram of each crystallization stage. Upon heating,  
79 interdiffusion of Al and Si atoms begins. After the diffused atoms become supersaturated,  
80 they begin to nucleate. The time necessary for Si nuclei to be observed with an optical  
81 microscope is defined as the incubation time. After that, these nuclei were grown laterally.  
82 The change of Si crystal radius per unit time is defined as growth rate. In growth stage, few  
83 nuclei were generated. The dependence of the incubation time and growth rate on the

84 thickness of the SiO<sub>2</sub> intermediate layer is shown in Fig. 2. The results for sample I were  
85 excluded as described later. This graph shows that the thicker the SiO<sub>2</sub> intermediate layer,  
86 the longer the incubation time and the smaller the growth rate. This means that the SiO<sub>2</sub>  
87 intermediate layer works as a diffusion barrier for Si atoms, which is supported by the fact  
88 that the diffusion coefficient of Si in SiO<sub>2</sub> is approximately 1000 times smaller than that in  
89 Al [21,22]. On the basis of the results shown in Fig. 2, it can safely be stated that the  
90 diffusion rate of Si depends on the SiO<sub>2</sub> intermediate layer thickness.

91 Figure 3 shows the crystal orientation mappings obtained by EBSD for AIC-Si layers  
92 fabricated with various SiO<sub>2</sub> intermediate layer thicknesses. When we employed the  
93 conventional AIC method, highly (111)-oriented AIC-Si layers were formed in sample A as  
94 shown in Fig. 3(a). We next discuss the influence of SiO<sub>2</sub> intermediate layer thickness on  
95 the AIC-Si crystal orientation. As shown in Figs. 3(b) and 3(f), the grain size of AIC-Si is  
96 considerably small, less than 1 μm both when the SiO<sub>2</sub> intermediate layer thickness is zero  
97 (sample B) and when it is very thick (20 nm, sample I). In the case of AIC-Si formed  
98 without the SiO<sub>2</sub> layer as in sample B, the Si atoms were thought to diffuse into the Al  
99 layer very rapidly, causing supersaturation shortly upon heating, and thus the Si atoms  
100 started to nucleate everywhere at the Si/Al interface, resulting in small grains as shown in  
101 Fig 3(b). In the case of very thick SiO<sub>2</sub> layer (20 nm, sample I), the results have not been  
102 completely understood. In general, sputtered SiO<sub>2</sub> films contain a certain amounts of

103 defects [23]. We speculate that these defects might significantly influence the diffusion of  
104 Si atoms and thereby generation of Si nuclei at the Al/SiO<sub>2</sub> interface, making it difficult to  
105 understand the AIC process by the simple model shown in Fig. 1. In contrast, when the  
106 SiO<sub>2</sub> intermediate layer was 2-10 nm in thickness (samples C-H), the AIC-Si grains  
107 exceeded tens of μm in size as shown in Figs. 3(c)-3(e). We think that the Si diffusion was  
108 suppressed moderately by this SiO<sub>2</sub> layer, and thus only a small number of Si nuclei were  
109 generated at the Al/SiO<sub>2</sub> interface, leading to lateral growth of Si nuclei as shown in Figs.  
110 3(c)-3(e). Figure 4 shows the SiO<sub>2</sub> intermediate layer thickness dependence of crystal  
111 orientation fractions of Si. The results obtained for sample I were not included. When the  
112 SiO<sub>2</sub> intermediate layer was 1 nm thick, the AIC-Si became highly (100) oriented. On the  
113 other hand, highly (111)-oriented AIC-Si layers were formed when the SiO<sub>2</sub> intermediate  
114 layer was 10 nm thick. The crystal orientation was determined from grains whose  
115 misorientation from the ideal (100) and (111) planes was less than 10°. The areas other  
116 than the (100) and (111) planes were much smaller, and thus were excluded in Fig. 4.  
117 These results suggest that the orientation of a poly-Si thin film fabricated by the AIC  
118 method depends significantly on the SiO<sub>2</sub> layer thickness, thereby the diffusion rate of Si.

119 Next, we discuss the mechanism of preferential AIC-Si crystal orientation dependence  
120 on the SiO<sub>2</sub> intermediate layer thickness, and thus the diffusion rate of Si atoms to the Al  
121 layer. The surface energy in crystalline Si is the lowest for the (111) plane, followed by the



122 (100) plane [24]. We think that the crystal plane dependence of the interface energy at the  
123 crystalline Si/amorphous SiO<sub>2</sub> interface is mostly the same as that in the crystalline Si,  
124 because SiO<sub>2</sub> has no anisotropy. The growth rate in Si is the largest on the (100) plane [25].  
125 The nucleation of Si occurs at the Al/SiO<sub>2</sub> interface in AIC, because the concentration of Si  
126 is the largest there within the Al layer. When the SiO<sub>2</sub> intermediate layer is thick and the Si  
127 diffusion rate is small, the concentration of Si in the Al increases gradually. Therefore, in  
128 the nucleation stage, only the most stable nuclei, (111)-oriented Si nuclei, are generated. In  
129 the crystal growth stage, these (111)-oriented nuclei grow laterally. As a result, the poly-Si  
130 film becomes (111) oriented. This (111) domination appears to contradict the theory  
131 presented by Schneider *et al* [26]. This contradict is possibly originated from the difference  
132 in the Al and Si thicknesses, because these thicknesses significantly affect the orientation  
133 of the AIC-Si layer [27]. In contrast, when the SiO<sub>2</sub> intermediate layer is thin and the Si  
134 diffusion rate is large, the degree of supersaturation is likely to become large. Therefore, in  
135 the nucleation stage, (100)-oriented nuclei, the second most stable nuclei, are generated in  
136 addition to the (111)-oriented nuclei. In the crystal growth stage, (100)-oriented nuclei  
137 grow rapidly because these grains have (100) facets. Therefore, (100)-oriented nuclei  
138 finally become dominant.

139

#### 140 **4. Conclusions**

141 In summary, we have fabricated poly-Si thin films on fused silica substrates by the AIC  
142 method using SiO<sub>2</sub> intermediate layers with various thicknesses (0-20 nm) between the a-Si  
143 (100 nm) and Al (100 nm) layers. The incubation time and growth rate of the Si grains  
144 depended on the SiO<sub>2</sub> layer thickness. With increasing SiO<sub>2</sub> layer thickness from 2 to 10  
145 nm, the crystal orientation of the AIC-Si changed gradually from the (100) to (111)  
146 orientation. This selective formation technique of (100) and (111) oriented poly-Si films  
147 opens up a possibility of high-quality epitaxial templates for thin film devices.

148

#### 149 **Acknowledgments**

150 This work was financially supported by Japan Science and Technology Agency,  
151 CREST and Grant-in-Aid for Research Activity Start-up (No. 23860011 MEXT). A part of  
152 this work was conducted by the AIST Nano-Processing Facility, supported by  
153 “Nanotechnology Network Japan” of MEXT, Japan.

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155 **References**

- 156 [1] O. Nast, T. Puzzer, L. M. Koschier, A. B. Sproul, and S. R. Wenham, *Appl. Phys. Lett.* 73  
157 (1998) 3214.
- 158 [2] Z. Shi and M. A. Green, *Prog. Photovoltaics* 6 (1998) 247.
- 159 [3] P. I. Widenborg and A. G. Aberle, *J. Cryst. Growth* 242 (2002) 270.
- 160 [4] E. Pihan, A. Slaoui, and C. Maurice, *J. Cryst. Growth* 305 (2007) 88.
- 161 [5] H. Kuraseko, N. Orita, H. Koaizawa, and M. Kondo, *Appl. Phys. Express* 2 (2009)  
162 015501.
- 163 [6] I. Gordon, L. Carnel, D. Van Gestel, G. Beaucarne, and J. Poortmans, *Thin Solid Films*  
164 516 (2008) 6984.
- 165 [7] B. R. Wu, S. Y. Lo, D. S. Wu, S. L. Ou, H. Y. Mao, J. H. Wang, and R. H. Horng, *Thin*  
166 *Solid Films* 520 (2012) 5860.
- 167 [8] R.B. Bergmann, G. Oswald, M. Albrecht, and V. Gross, *Sol. Energy Mater. Sol. Cells* 46  
168 (1997) 147.
- 169 [9] H. Kim, D. Kim, G. Lee, D. Kim, and S. H. Lee, *Sol. Energy Mater. Sol. Cells* 74 (2002)  
170 323.
- 171 [10] S. Gall, M. Muske, I. Sieber, O. Nast, and W. Fuhs, *J. Non-Cryst. Solids* 299-302 (2002)  
172 741.
- 173 [11] Y. Ishikawa, A. Nakamura, Y. Uraoka, and T. Fuyuki, *Jpn. J. Appl. Phys.* 43 (2004) 877.

- 174 [12] I. Gordon, L. Carnel, D. Gestel, G. Beaucarne, and J. Poortmans, *Prog. Photovoltaic* 15  
175 (2007) 575.
- 176 [13] P. Prathap, O. Tuzun, D. Madi, and A. Slaoui, *Sol. Energy Mater. Sol. Cells* 95 (2011)  
177 S44.
- 178 [14] K. Toko, M. Kurosawa, H. Yokoyama, N. Kawabata, T. Sakane, Y. Ohta, T. Tanaka, T.  
179 Sadoh, and M. Miyao, *Appl. Phys. Express* 3 (2010) 075603.
- 180 [15] D. Tsukada, Y. Matsumoto, R. Sasaki, M. Takeishi, T. Saito, N. Usami, and T. Suemasu,  
181 *Appl. Phys. Express* 2 (2009) 051601.
- 182 [16] S. M. Sze, *Physics of Semiconductor Devices*, p. 386, 2nd ed. (Wiley, New York, 1981)
- 183 [17] S. Gall, J. Schneider, J. Klein, K. Hübener, M. Muske, B. Rau, E. Conrad, I. Sieber, K.  
184 Petter, K. Lips, M. Stöger-Pollach, P. Schattschneider, and W. Fuhs, *Thin Solid Films*  
185 511-512 (2006) 7.
- 186 [18] Y. Sugimoto, N. Takata, T. Hirota, K. Ikeda, F. Yoshida, H. Nakashima, and H.  
187 Nakashima, *Jpn. J. Appl. Phys.* 44 (2005) 4770.
- 188 [19] M. Kurosawa, N. Kawabata, T. Sadoh, and M. Miyao, *Appl. Phys. Lett.* 95 (2009)  
189 132103.
- 190 [20] M. Jung, A. Okada, T. Saito, T. Suemasu, and N. Usami, *Appl. Phys. Express* 3 (2010)  
191 095803.
- 192 [21] S. Fujikawa, K. Hirano, and Y. Fukushima, *Metall. Trans. A* 9A (1978) 1811.

- 193 [22] S. Fukatsu, T. Takahashi, and K. M. Itoh, *Appl. Phys. Lett.* 83 (2003) 3897.
- 194 [23] Y. Urabe, T. Sameshima, K. Motai, and K. Ichimura, *Jpn. J. Appl. Phys.* 47 (2008) 8003.
- 195 [24] A. A. Stekolnikov, J. Furthmuller, and F. Bechstedt, *Phys. Rev. B* 65 (2002) 115318.
- 196 [25] L. Csepregi, E. F. Kennedy, and J. W. Mayer, *J. Appl. Phys.* 49 (1978) 3906.
- 197 [26] J. Schneider, A. Sarikov, J. Klein, M. Muske, I. Sieber, T. Quinn, H.S. Reehal, S. Gall,  
198 and W. Fuhs, *J. Crystal Growth* 287 (2006) 423.
- 199 [27] M. Kurosawa, K. Toko, N. Kawabata, T. Sadoh, and M. Miyao, *Solid-State Electron.* 60  
200 (2011) 7.
- 201

202 Fig. 1 (a) Schematic cross-sectional diagrams of each crystallization stage, optical  
203 micrographs of surface morphologies of sample D, prepared with a 2-nm-thick SiO<sub>2</sub>  
204 intermediate layer, after annealing for (b) 40, (c) 50, (d) 55, and (e) 60 min, and those for  
205 sample H, prepared with a 10-nm-thick SiO<sub>2</sub> intermediate layer after annealing for (f) 120, (g)  
206 150, (h) 180, and (i) 210 min.

207

208 Fig. 2 Dependence of the incubation time and growth rate on the SiO<sub>2</sub> intermediate layer  
209 thickness.

210

211 Fig. 3 Dependence of crystal orientation mappings of AIC-Si on the SiO<sub>2</sub> intermediate layer  
212 thickness for (a) sample A (native Al-oxide), (b) sample B (0 nm), (c) sample D (2 nm), (d)  
213 sample G (5 nm), (e) sample H (10 nm), and (f) sample I (20 nm).

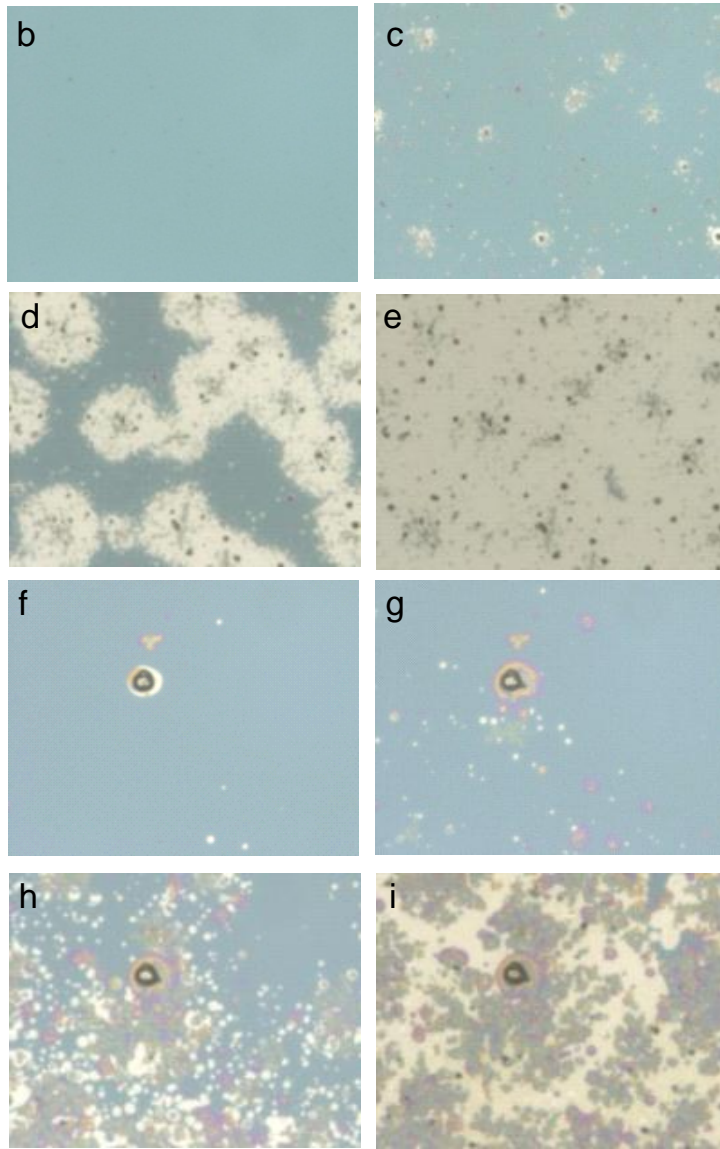
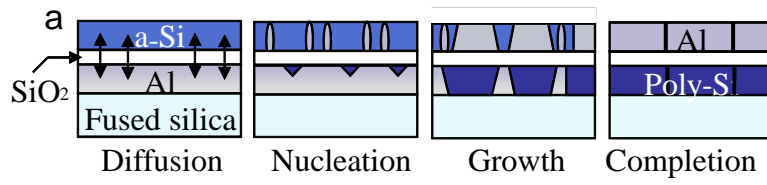
214

215 Fig. 4 Dependence of AIC-Si crystal orientation fractions on the SiO<sub>2</sub> intermediate layer  
216 thickness.

217

Table 1 Preparation of samples A-I. Al layer thickness, air exposure time, SiO<sub>2</sub> and a-Si layer thicknesses are shown.

Sample	Al (nm)	Air exposure (h)	SiO <sub>2</sub> (nm)	Si (nm)
A	100	48	0	100
B	100	0	0	100
C	100	0	1	100
D	100	0	2	100
E	100	0	3	100
F	100	0	4	100
G	100	0	5	100
H	100	0	10	100
I	100	0	20	100



— 50  $\mu\text{m}$

Figure 1



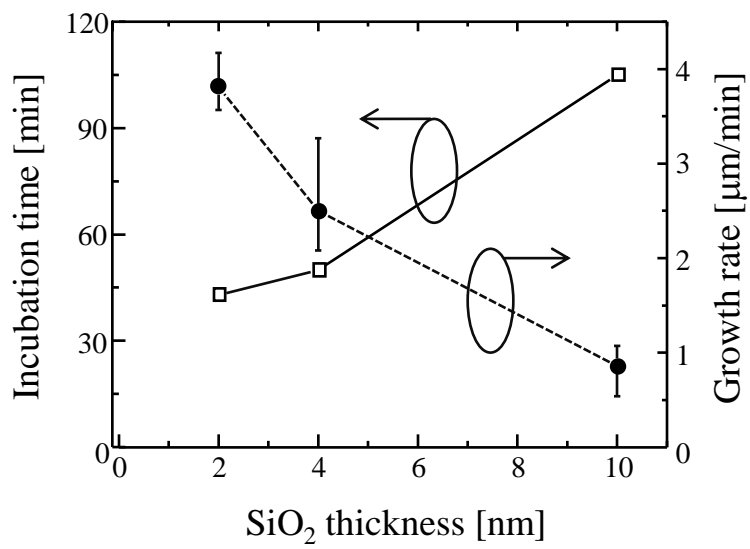


Figure 2

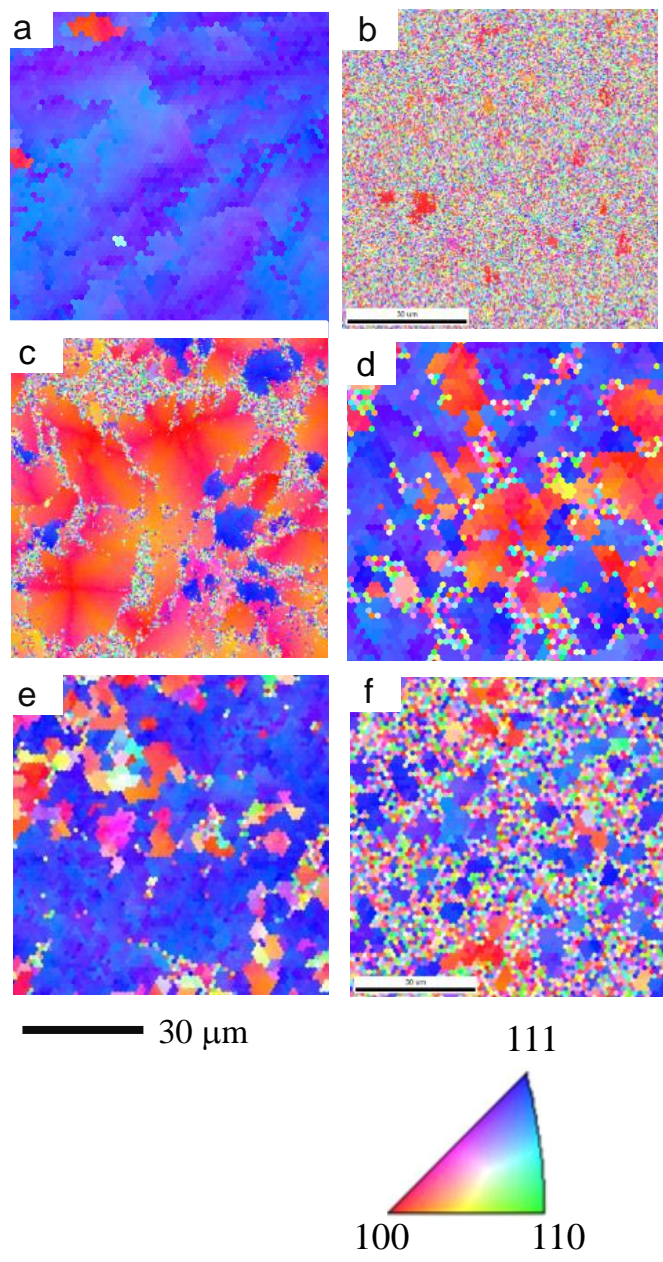


Figure 3

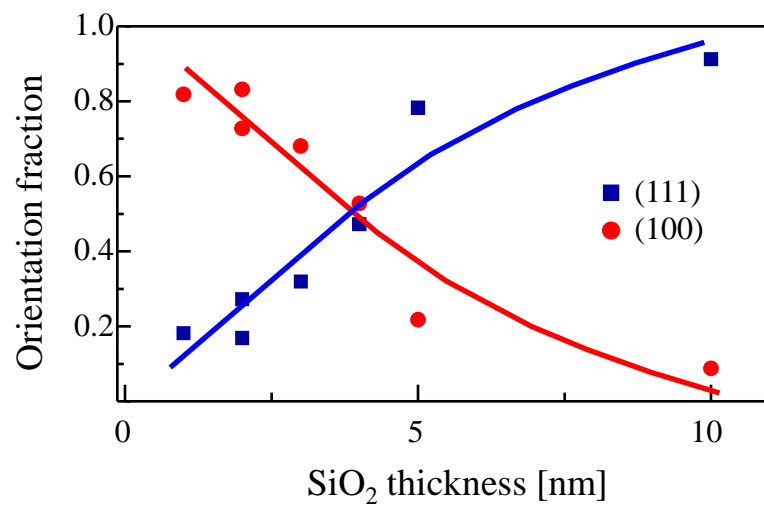


Figure 4