Hydrogen induced structural relaxation in bulk metallic glasses

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Abstract

In order to clarify the hydrogen induced structural relaxation (HISR) of hydrogenated bulk and marginal metallic glasses (BMG and MMG), the hydrogen concentration ($C_H$) dependence of the peak temperature ($T_p$) and the peak height ($Q^{-1}_p$) of the hydrogen internal friction peak (HIFP) in $a$-$Zr_{55}Cu_{30}Al_{10}Ni_{5}$ (numbers indicate at.%) (BMG), $a$-$Zr_{54}Cu_{30}Al_{10}Ni_{5}Si_{1}$ (MMG) and $a$-$Zr_{50}Cu_{49}Al_{10}Si_{1}$ (MMG) were studied. It is found that the $T_p$ vs. $C_H$ data and the $Q^{-1}_p$ vs. $C_H$ data are well explained by the relationships of $T_p = \Delta T_p \exp(-C_H/\tau_H) + T_{p,0}$ and $Q^{-1}_p \propto \ln(C_H/\tau_H)$, respectively, not only for all the present metallic glasses but also for various Zr-base BMG and MMG reported. The characteristic $C_H$ dependence of $T_p$ and $Q^{-1}_p$ was attributed to effects of the HISR. The detailed features of the $C_H$ dependence of $T_p$ and $Q^{-1}_p$ were discussed in the light of material parameters and HISR.

Keywords: Metallic glass; Hydrogen; Internal friction; Structural relaxation

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1. Introduction

A recent feasibility study of hydrogenated Zr-base metallic glasses as the high-damping and high-strength materials [1-3] has evoked the internal friction works on hydrogenated bulk metallic glasses [4-9]. In these materials, the peak height, $Q^{-1}_p$, and the peak temperature, $T_p$, of the hydrogen induced internal friction peak (HIFP) and the tensile strength are indicative of the high-damping and high-strength performance. Although hydrogenated metallic glasses are very brittle after heavy hydrogenation, the tensile strength of metallic glasses tend to increase with increasing hydrogen concentration below a few tens percent [1-3]. It is known that $T_p$ shows a decrease with increasing hydrogen concentration ($C_H$) for various metallic glasses [10]. Such a decrease in $T_p$ has been explained by a change in the chemical potential of hydrogen in metallic glasses with increasing $C_H$ [11,12]. On the other hand, recent intensive study on the HIFP for low $C_H$ for various metallic glasses revealed that the decrease in $T_p$ observed for low $C_H$ was much stronger than that expected from the chemical potential of hydrogen [1-10], indicating that the hydrogen induced structural relaxation (HISR) [13] may play an important role. In the present work, the hydrogen concentration dependence of $T_p$ and $Q^{-1}_p$ for low $C_H$ were studied in the light of HISR. Specimens prepared here were amorphous ($a$) $Zr_{55}Cu_{30}Al_{10}Ni_{5}$ (numbers indicate at.%) as an example of bulk metallic glasses (BMG)
and \(a\)-Zr\(_{55}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\) as a silicon-modified BMG. As a reference material, \(a\)-Zr\(_{40}\)Cu\(_{49}\)Al\(_{10}\)Si\(_{1}\) was revisited as a silicon-modified marginal metallic glass (MMG) which showed the highest \(T_p\) in the as hydrogen charged state among various \(a\)-alloys [2]. The present results will be discussed together with the previous data.

2. Experimental

Alloy ingots were prepared by arc-melting. Amorphous ribbon specimens of \(a\)-Zr\(_{55}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\), \(a\)-Zr\(_{54}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\)Si\(_{1}\) and \(a\)-Zr\(_{40}\)Cu\(_{49}\)Al\(_{10}\)Si\(_{1}\) were prepared by melt spinning in a high-purity Ar gas atmosphere, where the thickness and the width of ribbons were about 30 \(\mu\)m and 1 mm, respectively. The specimen surfaces were polished mechanically in water avoiding heating up during polishing to remove a surface layer and to smoothen. Hydrogen charging was made electrolytically in 0.1N H\(_2\)SO\(_4\) solution at room temperature (RT). The hydrogen charged specimen was aged for a few days at RT before measurements to homogenize the hydrogen distribution in the specimen. Specimen \(a\)-Zr\(_{40}\)Cu\(_{49}\)Al\(_{10}\)Si\(_{1}\) was annealed at 400 K in order to complete the probable thermal structural relaxation prior to the measurements of internal friction, of length change and of X-ray diffraction (XRD). The internal friction, \(Q^{-1}\), and the resonant frequency, \(f\), were measured by means of the vibrating reed method at about 300 Hz and the strain amplitude of \(10^{-6}\). The internal friction measurements were conducted in the temperature range between 80 K and 360 K. A change in the specimen length, \(\Delta L/L\), due to the hydrogenation was measured by an optical microscope with a micrometer stage. The XRD measurements were made by the conventional \(\theta\)-2\(\theta\) scan using the Cu K\(\alpha\) radiation. The hydrogen concentration, \(C_H\), was determined by the thermal desorption method similar to that mentioned in [13].

3. Results

Figures 1(a) and 1(b) show examples of HIFPs observed at \(~300\) Hz in \(a\)-Zr\(_{55}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\) and \(a\)-Zr\(_{54}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\)Si\(_{1}\) specimens with various \(C_H\), respectively. The HIFP was observed as a very broad peak accompanied with the low temperature tail, indicating that the activation energy shows a wide distribution as already reported but for \(a\)-CuTi [15]. As seen in Figs. 1(a) and 1(b) for both the \(a\)-Zr\(_{55}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\) and \(a\)-Zr\(_{54}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\)Si\(_{1}\) specimens: \(T_p\) of HIFP shows a rapid decrease with increasing \(C_H\) below about 15 at\% and then tends to show saturation for higher values: \(Q^{-1}\) also tends to show saturation following a rapid increase with increasing \(C_H\).

Figure 2(a) shows the \(T_p\) vs. \(C_H\) data observed for the reference material (MMG) \(a\)-Zr\(_{40}\)Cu\(_{49}\)Al\(_{10}\)Si\(_{1}\) (after 400 K annealing) and the bulk metallic glass (BMG) \(a\)-Zr\(_{55}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\) together with those reported for (BMG) \(a\)-Zr\(_{60}\)Cu\(_{30}\)Al\(_{10}\) [1]. It was found for various \(a\)-alloys that the \(T_p\) vs. \(C_H\) data can be well explained by,

\[
T_p = \Delta T_p \exp(-C_H/\tau_H) + T_{p,0},
\]

(1)

where \(\Delta T_p\) and \(\tau_H\) have usual meanings and \(T_{p,0}\) is the ultimate value of \(T_p\) found at the high \(C_H\) region. As will be mentioned, \(\Delta T_p\) is indicative of the distance of the glass structure measured from the ultimate glass after the HISR. Figures 2(b) and 2(c) are enlarged redrawing of the \(T_p\) vs. \(C_H\) data of fig. 2(a), but here the \(\ln(T_p-T_{p,0})\) is plotted against \(C_H\) together with the results obtained in the silicon-modified (BMG) \(a\)-Zr\(_{54}\)Cu\(_{30}\)Al\(_{10}\)Ni\(_{5}\)Si\(_{1}\) and in other metallic glasses previously investigated [2]. It is more clearly seen in Figs. 2(b) and 2(c)
that the $T_p$ vs. $C_H$ data observed for various metallic glasses are well explained by Eq. (1). Figure 3 shows the $\Delta T_p$ vs. $\tau_H$ data found for various Zr-Cu base metallic glasses, $a$-Zr$_50$Ni$_{50}$ and $a$-Cu$_{50}$Ti$_{50}$ [1-3,13,15]. In Fig. 3, the $\Delta T_p$ vs. $\tau_H$ data found for all the Zr-Cu base metallic glasses and $a$-Zr$_50$Ni$_{50}$ and $a$-Cu$_{50}$Ti$_{50}$ fall around the dashed line 1 (Group 1) except the data found for $a$-Zr$_5$Cu$_{30}$Al$_{10}$Ni$_5$, $a$-Zr$_{44}$Cu$_{30}$Al$_{10}$Ni$_5$Si$_1$ and $a$-Zr$_5$Cu$_{50}$ classified as Group 2. In this Fig. 3, $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ shows the highest $T_p$ in the as charged state among various Zr-Cu base metallic glasses [2]: after 400 K annealing, $\Delta T_p$ markedly decreases and fits well the dashed line 1. This finding denotes that the strong thermal structural relaxation in the as charged state took place in $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$. It is noted that BMGs $a$-Zr$_{60}$Cu$_{30}$Al$_{10}$ and $a$-Zr$_{55}$Cu$_{50}$Al$_{10}$Ni$_5$ are classified as the Group 1 and the Group 2, respectively, and Zr-Cu base metallic glasses found in the Group 2 are those with the compositional Zr concentration between 50 at% and 55 at%.

Figure 4 (a) shows the $Q_1^p$ vs. $C_H$ data observed for $a$-Zr$_{55}$Cu$_{50}$Al$_{10}$Ni$_5$ and $a$-Zr$_{54}$Cu$_{50}$Al$_{10}$Ni$_5$Si$_1$, where $Q_1^p$ increased monotonously showing the tendency of saturation. Figure 4(c) is redrawing of Fig. 4(a) but here $Q_1^p$ is plotted against $\ln(C_H/\tau_H)$. Figure 4(b) shows the relaxation strength, $S_t$, vs. $\ln(C_H/\tau_H)$ data, where $S_t$ was evaluated as the modulus defect observed in the $f$ data between 80 K and 360 K. In Fig. 4(b), $S_t$ shows a linear increase with increasing $\ln(C_H/\tau_H)$ in the whole range. The $Q_1^p$ vs. $\ln(C_H/\tau_H)$ data shown in Fig. 4(c) is very similar to those for $S_t$ except for $\ln(C_H/\tau_H)$ below -1 or $C_H$ below 5.5 at%. It is noted that $Q_1^p$ and $S_t$ measure the constituent relaxation strength at around $T_p$ and the total relaxation strength including the low temperature tail of the HIFP, respectively. Figures 5(a) and (b) are similar to Fig. 4(c) but display the data of: $a$-Zr$_{40}$Cu$_{50}$Al$_{10}$; $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ in the as charged state; $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ after 400 K annealing; $a$-Zr$_{50}$Cu$_{30}$Al$_{10}$ and $a$-Zr$_{59}$Cu$_{50}$Al$_{10}$Si$_1$. The linear increase in $Q_1^p$ with increasing $\ln(C_H/\tau_H)$ can be commonly seen except that the deviatory high $Q_1^p$ data were observed in the as charged state for $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ here and $a$-Zr$_{40}$Cu$_{50}$Al$_{10}$ [2] and $a$-Zr$_{50}$Cu$_{50}$Al$_{10}$ [1].

Figure 6 shows changes in the specimen length, $\Delta L/L_0$, due to hydrogenation observed for $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ in the as charged state and those measured after 400 K annealing, where $\Delta L/L_0$ showed the almost linear increase with increasing $C_H$. The hydrogen induced volume expansion, $(\Delta V/V)_H (\approx 3(\Delta L/L_0))$, per unit $C_H$ seen in Fig. 6 is very similar to that reported for various metallic glasses [16,17]. The $\Delta L/L_0$ data shown in Fig. 6 indicate that $(\Delta V/V)_H$ remained unchanged after 400 K annealing in the present experimental accuracy.

4. Discussion

As already mentioned, the decrease in $T_p$ due to a change in the chemical potential of hydrogen in metallic glasses with increasing $C_H$ [11,12] is expected to be a function of the constituent concentration of alloy elements with hydrogen affinity. On the other hand, as it can be evaluated from the slopes (Eq. 1) of the plots of Figs. 2(b) and 2(c), $\tau_H \sim 14$ at.% found in $a$-Zr$_{40}$Cu$_{50}$Al$_{10}$ or $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ is much larger than $\tau_H \sim 5$ at.% found in $a$-Zr$_{50}$Cu$_{50}$Al$_{10}$, indicating that the decrease in $T_p$ with increasing $C_H$ is mainly associated with the HISR. Then the $\Delta T_p$ vs. $\tau_H$ data shown in Fig. 3 may be indicative of the HISR in various metallic glasses. As seen in Fig. 2 or estimated from the data reported in Refs. [1-7,9,10], $T_{p,0}$ in Zr-base metallic glasses is found near 200 K, which is $T_p$ of the HIFP in both Zr-base BMGs and MMGs after completion of the HISR, the ultimate glass after the HISR hereafter. In other words, $\Delta T_p$ is indicative of the distance of the glass structure measured from the ultimate glass after the HISR, where an increase in the HISR is expected with increasing $\Delta T_p$ as observed for $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$. For the Group 1 metallic glasses shown in Fig. 3, $\tau_H$ increases with
increasing $\Delta T_p$, where $\Delta T_p$ and $\tau_H$ are the lowest for $a$-Zr$_{60}$Cu$_{30}$Al$_{10}$ (BMG) and $a$-Cu$_{50}$Ti$_{50}$ (MMG) among those for the Group 1 metallic glasses. On the other hand, the crystallization volume, $(\Delta V/V)_x$, is 0.3 % for $a$-Zr$_{60}$Cu$_{30}$Al$_{10}$ [18] and 0.8 % for $a$-Cu$_{50}$Ti$_{50}$ [19], respectively. For the Group 2 metallic glasses shown in Fig. 3, $(\Delta V/V)_x$ is 0.4 % for $a$-Zr$_{55}$Cu$_{30}$Al$_{10}$Ni$_{15}$ [18] and 2.1 % for $a$-Zr$_{50}$Cu$_{50}$ [20], respectively. This is indicative that the correlation between $(\Delta V/V)_x$ and $\Delta T_p$ and/or $\tau_H$ is weak. No changes in $\Delta L/L_0$ observed for the thermal structural relaxation of $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_{1}$ shown in Fig. 6 suggest that the structural relaxation responsible for changes in $T_p$ is hardly associated with a change in the specimen volume. The linear relationship found between $Q^{-1}_p$ and $\ln(C_H/\tau_H)$ cannot be explained by the conventional model for the HIFP [11] and may be one of characteristics of the HISR. One can speculate that the HISR brings about the shear deformation of the local glass structures resulting in a decrease in the elastic anisotropy around hydrogen atoms.

Figures 7(a) and 7(b) are redrawings of the $\Delta T_p$ vs. $\tau_H$ data shown in Fig. 3 but here $\Delta T_p$ and $\tau_H$ are plotted against the mean heat of solution, $\Delta H$, estimated from the chemical composition of specimens and the heat of solution for elementary materials [21-23]. Except for $a$-Zr$_{50}$Ni$_{50}$, specimens can be classified into two groups, the low-$\Delta H$ group with $\Delta H < -10$ kJ/mol and the high-$\Delta H$ group with $\Delta H > -6$ kJ/mol. For the low-$\Delta H$ group, $\Delta T_p$ remains constant and $\tau_H$ shows an increase with increasing $\Delta H$. For the high-$\Delta H$ group, both $\Delta T_p$ and $\tau_H$ shows an increase with increasing $\Delta H$. The underlying mechanism for the $\Delta H$ dependence of $\Delta T_p$ and $\tau_H$ is not known at present, however, the present work demonstrates the new aspect for the HISR.

5. Conclusion

The $C_H$ dependence of $T_p$ and $Q^{-1}_p$ of the HIFP in $a$-Zr$_{55}$Cu$_{30}$Al$_{10}$Ni$_{5}$, $a$-Zr$_{54}$Cu$_{30}$Al$_{10}$Ni$_{5}$Si$_{1}$ and $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_{1}$ were measured. It is found that the $T_p$ vs. $C_H$ data and the $Q^{-1}_p$ vs. $C_H$ data are well explained by the relationships of $T_p = \Delta T_p \exp(-C_H/\tau_H) + T_{p,0}$ and $Q^{-1}_p \propto \ln(C_H/\tau_H)$, respectively, not only for all the present metallic glasses but also for various Zr-base metallic glasses reported. The characteristic $C_H$ dependence of $T_p$ and $Q^{-1}_p$ was attributed to effects of the hydrogen induced structural relaxation (HISR). The detailed features of the $C_H$ dependence of $T_p$ and $Q^{-1}_p$ were discussed in the light of material parameters and HISR.

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References

Fig. 1.
Examples of HIFPs observed at about 300 Hz: (a) $a$-Zr$_{55}$Cu$_{30}$Al$_{10}$Ni$_{5}$
($\bullet$: 1.1 at%, $\bigcirc$: 4.7 at%, $\triangle$: 15.3 at%, $\blacktriangle$: 31.3 at%, $\square$: 59.1 at%).
(b) $a$-Zr$_{54}$Cu$_{30}$Al$_{10}$Ni$_{5}$Si$_{1}$ ($\bullet$: 1.3 at%, $\bigcirc$: 5.1 at%, $\triangle$: 15.8 at%, $\blacktriangle$: 30.4 at%, $\square$: 64.4 at%). The dashed curves denote a theoretical Debye peak with $\tau_0 = 10^{-12}$ s and $E = 0.39$ eV.
Fig. 2.
Examples of the $T_p$ vs. $C_H$ data ((a)) and examples of the $\ln(T_p-T_{p,0})/K$ vs. $C_H$ data ((b), (c)) found for various α-alloys (●: α-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ (after 400 K annealing), □: α-Zr$_{55}$Cu$_{30}$Al$_{10}$Ni$_{5}$, △: α-Zr$_{60}$Cu$_{30}$Al$_{10}$ [1], ■: α-Zr$_{54}$Cu$_{30}$Al$_{10}$Ni$_{5}$Si$_1$, ▲: α-Zr$_{59}$Cu$_{30}$Al$_{10}$Si$_1$ [2], ○: α-Zr$_{40}$Cu$_{50}$Al$_{10}$ [2], ◇: α-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_1$ (as-charged) [2]).
Fig. 3.
The $\Delta T_p$ vs. $\tau_H$ data observed for present specimens together with various metallic glasses reported. Group 1; ◆: $a$-$Zr_{60}Cu_{40}$ [1-3], ▲: $a$-$Zr_{60}Cu_{30}Al_{10}$ [1-3], △: $a$-$Zr_{59}Cu_{30}Al_{10}Si_{1}$ [1-3], ▽: $a$-$Zr_{40}Cu_{50}$ [13], ●: $a$-$Zr_{40}Cu_{50}Al_{10}$ [1-3], ○: $a$-$Zr_{40}Cu_{40}Al_{10}Si_{1}$ (as-charged) [1-3], ◎: $a$-$Zr_{40}Cu_{40}Al_{10}Si_{1}$ (after 400 K annealing), ▼: $a$-$Zr_{30}Cu_{50}Ti_{10}$, ×: $a$-$Zr_{50}Ni_{50}$, +: $a$-$Cu_{50}Ti_{50}$ [15]. Group 2; ■: $a$-$Zr_{55}Cu_{30}Al_{10}Ni_{5}$, □: $a$-$Zr_{54}Cu_{30}Al_{10}Ni_{5}Si_{1}$, ◇: $a$-$Zr_{50}Cu_{50}$ [1-3].
Fig. 4.
(a) The $Q^{-1}_p$ vs. $C_H$ data, (b) the $S_t$ vs. $\ln(C_H/\tau_H)$ data and (c) the $Q^{-1}_p$ vs. $\ln(C_H/\tau_H)$ data observed for $a$-Zr$_{55}$Cu$_{30}$Al$_{10}$Ni$_5$ (●, $\tau_H \approx 14.8$ at.%) and $a$-Zr$_{54}$Cu$_{30}$Al$_{10}$Ni$_5$Si$_1$ (○, $\tau_H \approx 15.5$ at.%). Dashed lines are drawn to guide eyes.
The $S_t$ vs. $\ln(C_H/\tau_H)$ data observed for various metallic glasses: (a) : $a$-Zr$_{40}$Cu$_{50}$Al$_{10}$ ($\bullet$, $\tau_H \approx 14$ at.%); $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_{1}$ (as charged) ($\bigcirc$, $\tau_H \approx 14$ at.%); $a$-Zr$_{40}$Cu$_{49}$Al$_{10}$Si$_{1}$ (after 400 K annealing) ($\triangle$, $\tau_H \approx 13.4$ at.%). (b) : $a$-Zr$_{60}$Cu$_{30}$Al$_{10}$ ($\blacksquare$, $\tau_H \approx 4.9$ at.%); $a$-Zr$_{59}$Cu$_{30}$Al$_{10}$Si$_{1}$ ($\square$, $\tau_H \approx 8$ at.%). Dashed lines are drawn to guide eyes.
Fig. 6.
The $\Delta L/L_0$ vs. $C_H$ data observed for $\alpha$-Zr$_{40}$Cu$_{40}$Al$_{10}$Si$_1$ in the as charged state (●) and those after 400 K annealing (○). The dashed curve is tentatively fitted to the data points.
Fig. 7.
(a) The $\Delta T_p$ and (b) the $\tau_H$ data are plotted against the mean heat of solution.

$\bullet$: $a$-$Zr_{60}Cu_{40}$, $\blacktriangle$: $a$-$Zr_{60}Cu_{30}Al_{10}$, $\triangle$: $a$-$Zr_{59}Cu_{30}Al_{10}Si_1$, $\triangledown$: $a$-$Zr_{40}Cu_{60}$, $\bullet$: $a$-$Zr_{40}Cu_{50}Al_{10}$, $\bigcirc$: $a$-$Zr_{40}Cu_{49}Al_{10}Si_1$ (as-charged), $\bigodot$: $a$-$Zr_{40}Cu_{49}Al_{10}Si_1$ (after 400 K annealing), $\blacktriangledown$: $a$-$Zr_{30}Cu_{60}Ti_{10}$, $\times$: $a$-$Zr_{50}Ni_{50}$, $\blacktriangleleft$: $a$-$Cu_{50}Ti_{50}$, $\blacklozenge$: $a$-$Zr_{35}Cu_{30}Al_{10}Ni_{5}$, $\square$: $a$-$Zr_{34}Cu_{30}Al_{10}Ni_{5}Si_1$, $\bigodot$: $a$-$Zr_{30}Cu_{50}$. See text for details.