

Development of an inversion method to extract information on fault geometry from teleseismic data

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SUMMARY

Teleseismic waveforms contain information on fault slip evolution during an earthquake, as well as on the fault geometry. A linear finite-fault inversion method is a tool for solving the slip-rate function distribution under an assumption of fault geometry as a single or multiple-fault-plane model. An inappropriate assumption of fault geometry would tend to distort the solution due to Green's function modelling errors. We developed a new inversion method to extract information on fault geometry along with the slip-rate function from observed teleseismic waveforms. In this method, as in most previous studies, we assumed a flat fault plane, but we allowed arbitrary directions of slip not necessarily parallel to the assumed fault plane. More precisely, the method represents fault slip on the assumed fault by the superposition of five basis components of potency-density tensor, which can express arbitrary fault slip that occurs underground. We tested the developed method by applying it to real teleseismic P waveforms of the M_W 7.7 2013 Balochistan, Pakistan, earthquake, which is thought to have occurred along a curved fault system. The obtained spatiotemporal distribution of potency-density tensors showed that the focal mechanism at each source knot was dominated by a strike-slip component with successive strike angle rotation from 205° to 240° as the rupture propagated unilaterally towards the south-west from the epicentre. This result is consistent with Earth's surface deformation observed in optical satellite images. The success of the developed method is attributable to the fact that teleseismic body waves are not very sensitive to the spatial location of fault slip, whereas they are very sensitive to the direction of fault slip. The method may be a powerful tool to extract information on fault geometry along with the slip-rate function without requiring detailed assumptions about fault geometry.

Key words: Image processing; Inverse theory; Time-series analysis; Earthquake dynamics; Earthquake source observations.

1 INTRODUCTION

In general, the seismic waveform contains information on the spatiotemporal moment density tensor distribution. This means that information on the slip-rate function, rupture extent and fault geometry is included in the seismic waveform. Finite-fault inversion of seismic waveforms, a tool that can be used to estimate the spatiotemporal slip-rate distribution of an earthquake, has been used to analyse seismic-source processes since the 1980s (e.g. Olson & Apsel 1982; Hartzell & Heaton 1983). However, the non-uniqueness of the source models for the same earthquake obtained by different inversion schemes has been an issue in seismic source inversion studies (e.g. Beresnev 2003; Mai *et al.* 2016). Simplified assumptions of Earth's structure and fault geometry are dominant sources of modelling errors. Modelling errors tend to distort solutions if

we neglect it (e.g. Yagi & Fukahata 2008, 2011), which can be one of the causes of the non-uniqueness and may critically affect the interpretation of rupture process.

The major source of modelling errors is the uncertainty in Green's functions; that is, a discrepancy between the true and the calculated Green's function due to inaccuracy of the Earth's structural model (e.g. Yagi & Fukahata 2011). Several inversion schemes have been developed to mitigate the effects of uncertainty in Green's functions by introducing the uncertainty into the data covariance matrix (Yagi & Fukahata 2011; Minson *et al.* 2013; Duputel *et al.* 2014). Yagi & Fukahata (2011) showed that their formulation made it possible to stably estimate the slip-rate function without forcing non-negative slip, which had been the common constraint to obtain stable estimate of slip (e.g. Du *et al.* 1992; Freymueller *et al.* 1994). Minson *et al.* (2013) also proposed a stable inversion method without a

smoothing constraint as a prior probability-density function, although an implicit smoothing constraint that forces constant slip on each source knot is still imposed in their method (e.g. Bodin *et al.* 2009; Nocquet 2018). These studies have opened a window into a stable and flexible seismic source inversion method that requires fewer prior constraints by more properly evaluating errors included in the seismic source inversion.

Another major source of the modelling errors is uncertainty in fault geometry. A simplified, inappropriate approximation of fault geometry by a flat or curved model plane may lead to inappropriate source models even if the uncertainty in Green's functions is taken into account, especially for an earthquake that occurs in a geometrically complex fault system (e.g. Gombert *et al.* 2018; Ragon *et al.* 2018). Ragon *et al.* (2018) introduced the uncertainty in fault geometry into the data covariance matrix based on the framework of Duputel *et al.* (2014), making it possible to robustly estimate the slip distribution even under an assumption that the approximated fault geometry is slightly fluctuated from the true one. The introduction of the uncertainty in fault geometry into an inversion scheme may contribute to constructing an unbiased source model, but is not designed to extract the information on fault geometry, which is included in the observed data (e.g. Fukahata & Wright 2008; Minson *et al.* 2014; Chen *et al.* 2014; Zhang & Wang 2015; Duputel & Rivera 2017; Shi *et al.* 2018). Information on the fault geometry is valuable for understanding the details of rupture propagation because discontinuities or spatial variations in the fault geometry play an important role in the rupture propagation (e.g. Aki 1979; Okuwaki & Yagi 2018), especially for earthquakes such as the 2008 Wenchuan, China (M_W 7.9), the 2013 Balochistan, Pakistan (M_W 7.7) and the 2016 Kaikoura, New Zealand (M_W 7.8) earthquakes, which occurred in geometrically complex fault systems. In order to reveal the relationship between the geometrical complexity of a fault and the rupture propagation, the finite-fault inversion framework adopting finely discretized model-source space is preferable, rather than the multiple-point-source inversions (e.g. Kikuchi & Kanamori 1991; Duputel *et al.* 2012).

In this study, we propose a new finite-fault inversion framework of teleseismic waveforms to estimate both the fault geometry and the shear-slip evolution by introducing the Green's function uncertainty into the data covariance matrix (Yagi & Fukahata 2011) and by increasing the number of shear-slip components from two to five, which can represent arbitrary direction of shear-slip direction (i.e. the direction of shear-slip is not confined on the assumed fault plane). It is generally difficult to estimate both a complex fault geometry and the shear-slip evolution simultaneously because of the non-linearity of the inverse problem to be solved (e.g. Fukahata & Wright 2008; Ragon *et al.* 2018). However, the teleseismic waveforms are insensitive to an assumed spatial location of shear-slip but sensitive to an assumed direction of shear-slip. Hence, we can estimate the spatiotemporal variation of focal mechanisms along the model plane, which represents the information on the true fault geometry. Possible bias in the seismic source image due to an inappropriate assumption of fault geometry can be significantly reduced in our proposed inversion framework that can represent the shear-slip direction independently from the assumed model fault plane.

We applied the developed method to the M_W 7.7 2013 Balochistan, Pakistan, earthquake. Detailed analyses of the optical satellite images following this earthquake had revealed a curved fault geometry consistent with the tectonic fabric around the potential source area (Avouac *et al.* 2014; Barnhart *et al.* 2014; Jolivet *et al.* 2014; Zinke *et al.* 2014). We expected that the complex

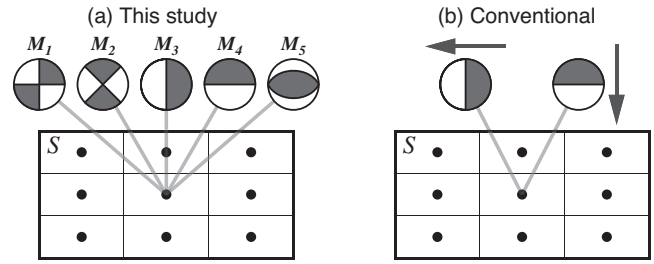


Figure 1. Representations of shear-slip at each source knot along a model plane S by (a) our proposed method and (b) conventional method. Each dot represents a source knot located at the centre of each subfault on the model plane S . (a) The five beach balls denote the five basis components of double-couple moment tensor (Kikuchi & Kanamori 1991); each component corresponds to a shear dislocation source, but is not needed to be parallel to the model plane S . The moment tensors, plotted by using a lower-hemisphere stereographic projection, are not rotated according to the model plane setting. (b) The two beach balls denote the two basis components of double-couple moment tensor that correspond to the shear dislocation sources along the model plane S . Arrows denote the slip directions of the two double-couple components whose nodal planes correspond to model plane S . The two moment tensors, plotted by using a lower-hemisphere stereographic projection, are rotated so that their one nodal planes correspond to model plane S ; therefore, they do not correspond to M_3 and M_4 (projected in (a)).

curved fault geometry of the 2013 Balochistan earthquake would help us to evaluate our new method that solves the spatiotemporal evolution of the potency-density tensor (e.g. Ampuero & Dahlen 2005).

2 MATHEMATICAL FORMULATION

Earthquake sources can be generally represented by volume density of moment-rate tensors $\dot{\mathbf{M}}(t, \xi)$ (e.g. Backus & Mulcahy 1976), which is expressed by a linear combination of five basis double-couple components \mathbf{M}_q (Kikuchi & Kanamori 1991, Fig. 1a):

$$\dot{\mathbf{M}}(t, \xi) = \sum_{q=1}^5 \dot{m}_q(t, \xi) \mathbf{M}_q, \quad (1)$$

where ‘five’ corresponds to the degree of freedom of the moment-rate tensor without isotropic expansion sources, $\dot{m}_q(t, \xi)$ is a spatiotemporal moment-rate volume-density function of q th basis component of moment tensor, and ξ represents a location in the source area. In a conventional scheme of finite-fault inversion, we solve a slip-rate function at each source knot with one or two directions of shear-slip on the assumed fault plane (e.g. Olson & Apsel 1982; Ide *et al.* 1996; Ji *et al.* 2002), which is equivalent to solving the potency-rate areal-density function for two double-couple components (Fig. 1b), where the potency rate corresponds to the moment rate divided by the rigidity of the medium. As shown in Fig. 1(b), one of the two nodal planes is constrained by the model fault plane in the conventional scheme. In this study, however, we propose a formulation for finite-fault inversion of teleseismic waveforms that solves the potency-rate areal-density function for five double-couple components (Fig. 1a), based on the framework of Yagi & Fukahata (2011).

We represent a seismic waveform u_j for a far-field term observed at a station j by a linear combination of moment-rate volume-density functions of five basis components of double-couple moment tensor

(Fig. 1a):

$$u_j(t) = \sum_{q=1}^5 \int_V G_{qj}^0(t, \xi) * \dot{m}_q(t, \xi) d\xi + e_{bj}(t), \quad (2)$$

where V is a 3-D source area, G_{qj}^0 is a Green's function, e_{bj} is a background and instrumental noise and $*$ denotes the convolution operator in the time domain. Eq. (2) requires the model source space to have three dimensions so as to represent the spatiotemporal moment-rate volume-density distribution. In order to further represent the 3-D source area as a simple finite-fault model in two dimensions, we assume that a finite-fault model obtained by a teleseismic body wave inversion is not seriously affected by a slight change of source location of the Green's function:

$$G_{qj}^0(t, \xi + \delta\xi) \approx G_{qj}^0(t, \xi), \quad (3)$$

where $\xi + \delta\xi$ is a neighbouring location around ξ . In fact, however, Green's function can change with source depth to some degree (e.g. Yagi *et al.* 2012). Through the synthetic tests, we evaluated the sensitivity to a slight change of source depth and confirmed that the change of source depth does not seriously affect the estimated distribution of focal mechanisms and overall features in the source-time function (Supporting Information Text. S1 and Fig. S9).

The approximation of eq. (3) makes it possible to change the volume integral to a surface integral on a model plane S embedded in the source area V . Then, eq. (2) becomes

$$u_j(t) = \sum_{q=1}^5 \int_S G_{qj}^0(t, \xi) * \dot{m}'_q(t, \xi) d\xi + e_{bj}(t), \quad (4)$$

where $\dot{m}'_q(t, \xi)$ is the moment-rate areal-density function, which is the projection of the volume density of the moment-rate function \dot{m}_q on the flat-model plane S .

The moment-rate areal-density function is represented by the potency-rate areal-density function \dot{D}_q :

$$\dot{m}'_q(t, \xi) = \mu(\xi) \dot{D}_q(t, \xi), \quad (5)$$

where $\mu(\xi)$ is the rigidity.

By substituting eq. (5) into eq. (4), we represent the observed waveform u_j by the potency-rate areal-density function \dot{D}_q :

$$u_j(t) = \sum_{q=1}^5 \int_S \tilde{G}_{qj}(t, \xi) * \dot{D}_q(t, \xi) d\xi + e_{bj}(t), \quad (6)$$

with

$$\tilde{G}_{qj}(t, \xi) = \mu(\xi) G_{qj}^0(t, \xi). \quad (7)$$

Eq. (6) means that the distribution of potency-rate areal-density tensors can be estimated in the framework of the finite-fault inversion method. The slip rate corresponds to potency-rate areal density where a shear plane and its area are explicitly given, and the potency-rate areal density is equal to the slip rate if the slip direction of the potency-rate areal-density tensor lies on the model plane S .

In the conventional formulation, one of the two nodal planes determined by the two basis components of double-couple potency tensor corresponds to the model plane S , which means that the fault geometry is approximated by the model plane S (Fig. 1b).

While in this study, we adopt $q = 5$ for the potency-rate areal-density function in eq. (6), which means that the shear plane is resolved with the five basis double-couple components of the potency-rate areal-density tensor, and it is therefore no longer required to be consistent with an arbitrarily presumed model plane geometry (Fig. 1a).

Following the multiple-time-window inversion method (e.g. Olson & Apsel 1982), we represent the potency-rate areal-density function \dot{D}_q by a linear combination of a finite number of basis functions:

$$\dot{D}_q \cong \sum_{k=1}^K \sum_{l=1}^L a_{qkl} X_k(\xi) T_l(t - t_k), \quad (8)$$

where X_k and T_l are the basis functions for space and time, respectively, and a_{qkl} are the expansion coefficients (model parameters) to be estimated from the observed waveform data. t_k is the time at which rupture can start on a space knot k .

By substituting eq. (8) into eq. (6), we can represent the observed waveform u_j by the basis functions of space and time:

$$u_j(t) = \sum_{k=1}^K \sum_{l=1}^L \sum_{q=1}^5 a_{qkl} \int_S \tilde{G}_{qj}(t, \xi) * (X_k(\xi) T_l(t - t_k)) d\xi + e_{bj}(t). \quad (9)$$

Following Yagi & Fukahata (2011), we introduce the modelling error of Green's function δG_{qj} as

$$\tilde{G}_{qj}(t, \xi) = G_{qj}(t, \xi) + \delta G_{qj}(t, \xi). \quad (10)$$

We can rewrite eq. (9) by using the modelling error of Green's function δG_{qj} :

$$u_j(t) = \sum_{k=1}^K \sum_{l=1}^L \sum_{q=1}^5 a_{qkl} \int_S (G_{qj}(t, \xi) + \delta G_{qj}(t, \xi)) * (X_k(\xi) T_l(t - t_k)) d\xi + e_{bj}(t). \quad (11)$$

The observation equation at the j th station can be expressed in vector form as

$$\mathbf{d}_j = \mathbf{H}_j \mathbf{a} + \mathbf{e}_{gj}(\mathbf{a}) + \mathbf{e}_{bj}, \quad (12)$$

where \mathbf{d}_j is the N_j (number of data point at the j th station) dimensional data vector that represents the observed waveform u_j , \mathbf{a} is the $M (=5KL)$ dimensional model parameter vector, \mathbf{H}_j is the $N_j \times M$ dimensional coefficient matrix that represents convolutions of the calculated Green's functions and time-domain basis functions, \mathbf{e}_{gj} is the N_j dimensional vector that represents the modelling error of Green's function and \mathbf{e}_{bj} is the N_j dimensional vector that represents the background noise. Following Yagi & Fukahata (2011), we derive the data covariance matrix of \mathbf{d} from eq. (12) and represent the stochastic model that relates the data vector \mathbf{d} for all stations by the model parameters \mathbf{a} .

As prior constraints, we use the smoothing constraint on each potency-rate areal-density function \dot{D}_q in space and time, represented as,

$$\nabla^2 \dot{D}_q(t, \xi) + e_s = 0, \quad (13)$$

$$\frac{\partial^2}{\partial t^2} \dot{D}_q(t, \xi) + e_t = 0. \quad (14)$$

As shown in eq. (13), we apply the smoothing constraint to the instantaneous spatial distribution of potency-rate areal-density function \dot{D}_q in space and time, which makes solutions more stable than

the smoothing constraint only on the total spatial distribution (e.g. Ide *et al.* 1996; Ji *et al.* 2002).

We can rewrite eqs (13) and (14) in vector form as,

$$\mathbf{S}_1 \mathbf{a} + \mathbf{e}_s = \mathbf{0}, \quad (15)$$

$$\mathbf{S}_2 \mathbf{a} + \mathbf{e}_t = \mathbf{0}, \quad (16)$$

where \mathbf{S}_1 and \mathbf{S}_2 are $M \times M$ dimensional matrices. We combine these two prior constraints into a single probability-density function, following the method of Fukahata *et al.* (2003, 2004).

The remaining formulation process, which is conducted following Yagi & Fukahata (2011), determines the unknown scaling parameters of the modelling error, background noise and prior constraints on the basis of Akaike's Bayesian Information Criterion (ABIC; Akaike 1980; Yabuki & Matsu'ura 1992) and iteratively estimates model parameter vectors. As shown by Fukuda & Johnson (2008), a problem of over-smoothing and under-smoothing can be avoided in our inversion scheme that adopts the ABIC without the non-negative constraint.

In the following sections, we use 'potency-rate density' and 'potency density' instead of 'potency-rate areal density' and 'potency areal density' for simplicity.

3 APPLICATION

3.1 2013 Balochistan earthquake

To test the formulation developed in this study, we applied it to the 2013 Balochistan, Pakistan, earthquake. The epicentre (26.951°N, 65.501°E) determined by the U.S. Geological Survey, National Earthquake Information Center (USGS NEIC; <https://earthquake.usgs.gov/earthquakes/eventpage/usb000jyiv>, last accessed 19 November 2018) was in the eastern part of the Makran accretionary wedge in the south-western Pakistan. In the eastern part of the Makran accretionary wedge, the Arabia Plate is subducting under the Eurasia Plate at a rate of $\sim 40 \text{ mm yr}^{-1}$ in the southern part of Makran accretionary wedge and the eastern side of the Makran accretionary wedge adjoins the India plate, which is converging northwards towards the Eurasia Plate at a rate of $\sim 40 \text{ mm yr}^{-1}$ (DeMets *et al.* 2010) (Fig. 2a). Analyses of optical satellite images (Avouac *et al.* 2014; Jolivet *et al.* 2014; Zinke *et al.* 2014) and of interferometric synthetic aperture radar (InSAR) data (Barnhart *et al.* 2014) show a curved surface rupture pattern for the Balochistan earthquake, which is consistent with the arcuate topographic features around the epicentre. The global centroid moment tensor (GCMT; <http://www.globalcmt.org/CMTsearch.html>, last accessed 19 November 2018) solution and the W -phase moment tensor solution determined by USGS (<https://earthquake.usgs.gov/earthquake/eventpage/usb000jyiv>, last accessed 19 November 2018) suggest that the Balochistan earthquake was a left-lateral strike-slip earthquake (Fig. 2c). The spatiotemporal distributions of high-frequency radiation sources tracked by backprojection techniques show mostly unilateral rupture propagation towards the south-west from the epicentre (Avouac *et al.* 2014). The possibly curved and complex fault geometry of the Balochistan earthquake makes construction of a reliable finite-fault model problematic, because use of a simplified fault geometry such as a single planar model or even a configuration of multiple rectangles would cause significant modelling errors as we reviewed in the Introduction and as we discuss further in Section 4.3. Thus, the 2013 Balochistan earthquake provides us an opportunity to test whether our method can extract a complex fault

geometry by using the finite-fault inversion framework developed in this study.

3.2 Data processing and model setting

The waveform data were downloaded from the Incorporated Research Institutions for Seismology Data Management Center (IRIS-DMC). We selected 36 vertical components of globally observed teleseismic P waveforms, ensuring a high signal-to-noise ratio and a good azimuthal distribution (Fig. 2b). We manually picked the first motion of the teleseismic P waveforms, removed the seismograph response from the original waveforms to convert into velocity waveforms, and then re-sampled the waveform data at 0.8-s intervals. Theoretical Green's functions were calculated following Kikuchi & Kanamori (1991) at 0.1-s intervals, and the attenuation time constant t^* for the P wave was taken to be 1.0 s. We used the 1-D near-source velocity structure (Table 1) used in Avouac *et al.* (2014) to calculate the theoretical Green's functions. We do not apply any filter to both the observed waveforms and the theoretical Green's functions. We set a single, flat-model plane, 200 km long and 45 km wide, with a strike of 230° and a dip of 0° (Fig. 2b) that entirely covered the possible surface rupture area inferred by Zinke *et al.* (2014). As described in Section 2, our inversion method enables us to estimate the spatiotemporal distribution of potency-rate density. Hence, the strike and dip angles of the assumed model plane are independent of the focal parameters on each space knot that we solve. We horizontally placed the model plane at 7.5 km depth, which minimized the ABIC value in our test of different hypocentral depth configurations. The model plane was expanded into bilinear B-spline functions with an interval of 5 km and adopted the epicentral location of 26.951°N, 65.501°E for the initial-rupture point. The potency-rate-density function at each source knot was represented as a linear combination of B-spline functions with an interval of 0.8 s, assuming a total duration of 31 s in order to flexibly capture the variation of rupture behaviours. We also assumed the potency-rate density to be zero from 60 s after the initiation of rupture on the basis of the estimated source-time function in Avouac *et al.* (2014) and Barnhart *et al.* (2014). We assumed the maximum rupture-front velocity at to be 4 km s^{-1} , taking into account the possibility of supershear rupture reported by Avouac *et al.* (2014).

3.3 Results

We obtained the distribution of potency-density tensors by integrating the potency-rate-density function for each basis component of potency tensor with respect to time at each space knot (Fig. 3). We represent the potency density at each space knot as an average of absolute value of eigenvalues for T - and P -axes of a diagonalised potency-density tensor at each space knot. The potency density is larger around the epicentre and extend to south-west from the epicentre about 150 km. It is difficult to uniquely determine fault parameters from a focal mechanism solution by our method alone. Here, we take the fault plane whose strike was in the range between 180° and 270°, based on the mapping of the surface rupture (e.g. Zinke *et al.* 2014). The potency-density tensors show strike-slip faulting with a strike of 216° and a dip of 77° around the epicentre (Fig. 3). We can also see strike-slip faulting with a strike of 243° and a dip of 45° around 80 km south-west from the epicentre, and reverse or oblique-slip faulting with a strike of 244° and a dip of 53° around 140 km south-west from the epicentre (Fig. 3). The total

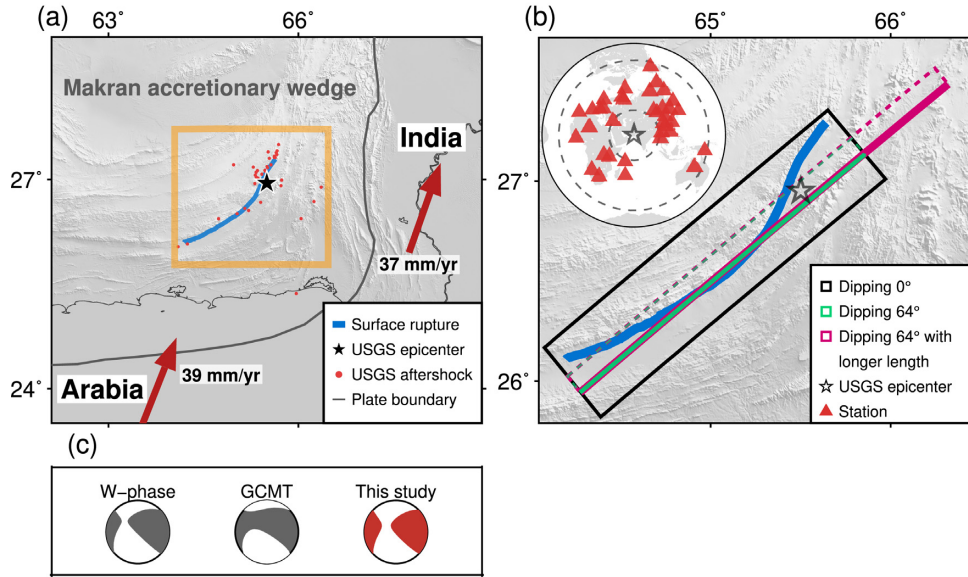


Figure 2. Summary of the tectonic setting and the surface-rupture trace of the 2013 Balochistan earthquake, the model plane settings, and the moment tensor solutions from USGS NEIC, GCMT and this study. (a) Surface rupture trace (blue line) was derived from Zinke *et al.* (2014). Red dots are aftershocks with $M_W \geq 3$ occurring within 1 week of the main shock (USGS NEIC; <https://earthquake.usgs.gov/earthquakes>, last accessed 19 November 2018). Arrows indicate the relative motions of the Arabia and India plates against the Eurasia plate by using the UNAVCO (<https://www.unavco.org>) plate motion calculator, adopting the MORVEL model (DeMets *et al.* 2010). The orange rectangle denotes the area shown in Fig. 2(b). Topography was drawn by using 2010 Global Multi-resolution Terrain Elevation Data 7.5 arcsecond median elevations (GMTED; https://topotools.cr.usgs.gov/gmted_viewer/) (Danielson & Gesch 2011) and the intensity of shade was derived from the directional derivative in 45° azimuthal direction. (b) Model planes used for the analyses of this study. The top of each model fault plane is shown by a bold line, and the dashed lines outline the dipping model planes. The star denotes the initial rupture point. The blue line represents the surface-rupture trace measured by Zinke *et al.* (2014). The inset map shows the distribution of stations (red triangles) around the epicentre (star) in an azimuthal equidistant projection. The dotted circles show the teleseismic distances at 30° and 90° . Topography was drawn by using GMTED and the intensity of shade was derived from the directional derivative in 45° azimuthal direction. (c) Moment tensor solutions. The grey beach balls show the *W*-phase CMT solution provided by USGS (<https://earthquake.usgs.gov/earthquakes/eventpage/usb000jyiv/moment-tensor>, last accessed 19 November 2018) and the GCMT solution from the Global CMT Catalogue Search (<http://www.globalcmt.org/CMTsearch.html>, last accessed 19 November 2018). The red beach ball represents the total moment tensor estimated by this study for a model plane dipping 0° .

Table 1. Near-source velocity model (from Avouac *et al.* 2014) used for calculating Green's functions.

V_P (km s^{-1})	V_S (km s^{-1})	Density (10^3 kg m^{-3})	Thickness (km)
5.44	3.00	2.50	4.00
6.25	3.45	2.60	12.00
6.53	3.60	2.70	14.00
6.80	3.90	2.90	12.00
7.50	4.30	2.90	3.00
8.11	4.49	3.30	0.00

moment tensor, which was obtained by taking a spatial integration of the potency-density tensors shown in Fig. 3, shows strike-slip faulting with a strike of 226° and a dip of 61° , and it includes 8 per cent of the non-double-couple components (Fig. 2c). The total released seismic moment was $7.53 \times 10^{20} \text{ Nm}$ (M_W 7.8), comparable to the GCMT solution of $5.59 \times 10^{20} \text{ Nm}$ (M_W 7.8). The source-time function, obtained by calculating the seismic moment at each time step, has a major peak at around 17 s and three minor peaks at around 29, 43 and 53 s (Fig. 3 and Supporting Information Fig. S14). To clearly visualize the spatiotemporal variation of focal mechanisms, we showed a potency-density tensor on its centroid location within 10-s time window by taking the spatial and temporal integration of the potency-rate density along the model plane

(Fig. 4). In this representation, the centroid location propagates unilaterally towards the south-west from the epicentre, with a longer centroid migration distance in the period from 20 to 50 s compared with that before 20 s or after 50 s. The potency-density tensor in each time window shows a successive transition during 0 to 60 s from strike-slip faulting with strike and dip angles of 205° and 68° , respectively, to oblique-slip faulting with strike and dip angles of 241° and 39° , respectively, with clockwise north-to-south rotation of the strike angle (Fig. 4). This transition of the focal mechanism in time is consistent with what can be seen in the static, spatial variation of the distribution of potency-density tensors (Fig. 3). This consistency can be seen even if we took 3-s time window instead of 10-s time window (Supporting Information Fig. S13). We also show the temporal change of focal mechanisms at all space knots by integrating the potency-rate-density function with respect to time in each time window and by plotting the strike and dip angles in the strike angle range between 180° and 270° (Fig. 5a). As can be seen in the histogram snapshots (Fig. 5a), the strike angle changes successively from 200° to 240° during 0 to 50 s, showing the clockwise rotation of the fault plane during rupture evolution. The dip angles are steep, around 80° , during 0 to 20 s, and then become shallow, around 60° , during 20 to 50 s. Note that during the fast migration of the centroid locations from 20 to 50 s (Fig. 4), the distribution of dip angles shows more scatter than the distribution during 0 to 20 s. During 50 to 60 s, the distributions of strike angle and dip angle are scattered. We also plotted the strike and dip angles of the two nodal

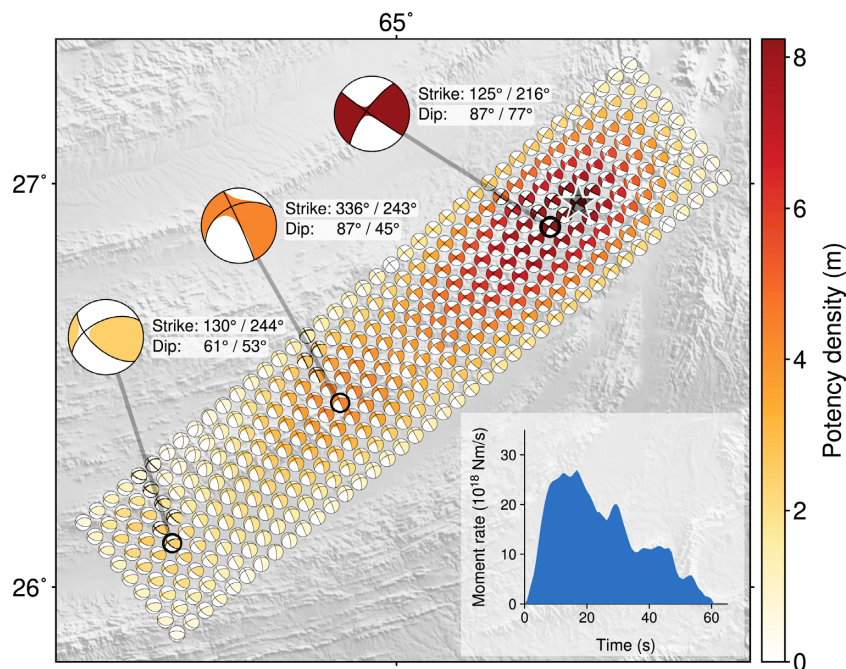


Figure 3. Spatial distribution of the potency-density tensors and moment-rate function (inset). The small beach ball shows the focal mechanism at each source knot, plotted using a lower-hemisphere stereographic projection. The star denotes the epicentre. The beach ball colour shows the potency density. Some beach balls (with black outline) have been selected and enlarged to highlight the spatial variation of the focal mechanism. Strike and dip angles are shown at right side of the enlarged beach balls.

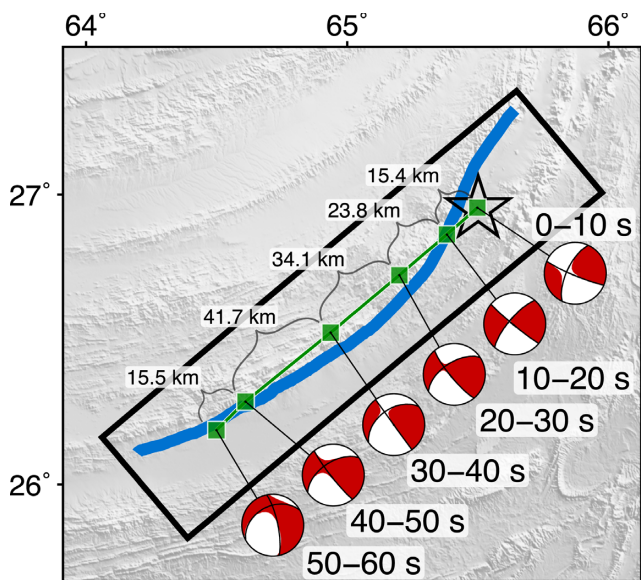


Figure 4. The evolution of the centroid of the potency-rate-density tensors. The centroid locations (green squares) of the potency-rate density and its potency-density tensor is shown at each time window, along with the distance between the neighbouring centroid locations and the corresponding time window for estimation of the centroid location. The blue line represents the surface-rupture trace measured by Zinke *et al.* (2014). The rectangle outlines the model plane. The star denotes the epicentre.

planes for all space knots without taking a certain range of strike angles for all time windows (Fig. 5b). Two strike angle modes at around 130° and 230° (Fig. 5b) reflect the two possible conjugated fault planes.

4 DISCUSSION

4.1 Rupture propagation and fault geometry

The unilateral rupture propagation towards the south-west from the epicentre (Fig. 4) is consistent with the spatiotemporal distribution of high-frequency radiation sources tracked by the backprojection technique (Avouac *et al.* 2014). The estimated average centroid location migration velocity during 20 to 50 s (Fig. 4) was 3.8 km s^{-1} , which is faster than the local shear wave velocity at the depth of the model plane that we placed (Table 1). Although we can not directly compare the migration velocity of the centroid location and the rupture front velocity inferred by the backprojection (Avouac *et al.* 2014), the period of 20 to 50 s, when the centroid migration velocity exceeds the shear-wave velocity, is consistent with the timing of a rupture velocity exceeding the local shear wave velocity observed by Avouac *et al.* (2014).

The trend of the surface-rupture trace measured by Zinke *et al.* (2014) has an azimuthal angle of 200° around the epicentre, 230° around 80 km south-west of the epicentre and 240° around 140 km south-west of the epicentre; these angles are in agreement with the strike angles of the total potency-density tensor distribution (Fig. 3) and the strike angle of the potency-density tensor in each 10 s window (Fig. 4). The estimated shallow-dipping fault in the south-western part of the source region (Fig. 3), which can also be seen in the temporal evolution of the rupture (Figs 4 and 5), may be related to an active thrust fault in the Makran accretionary wedge (Haghpor *et al.* 2012). The CLVD component of potency-density tensor at individual source knots is not apparently significant, implying that teleseismic *P* waves can be used to constrain a CLVD component of a shallow strike-slip earthquake if station distribution ensures a good azimuthal coverage. (Fig. 3).

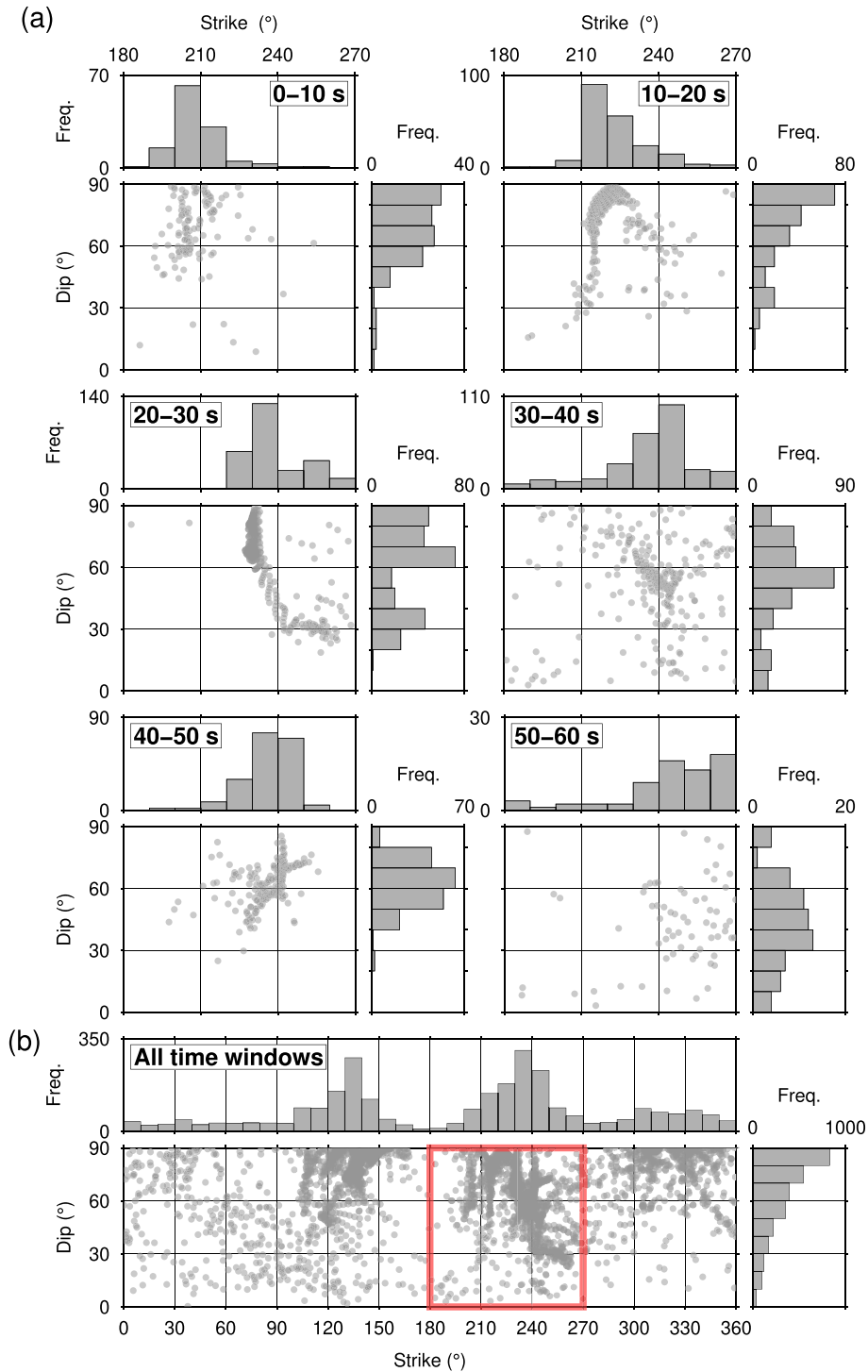


Figure 5. Configurations of strike and dip angles of all source knots in (a) each time window and (b) all time windows. The histograms group the angles in 10° bins. The red rectangle in (b) denotes the range of strike and dip angles shown in (a).

The results of the analysis show that we can estimate the direction of rupture propagation and the curved fault geometry, both of which are consistent with the backprojection analysis and optical satellite images, simply by distributing the source knots along a single model plane dipping 0° . The rupture extent of the Balochistan earthquake seems to be distributed mainly at depths shallower than 15 km (e.g. *Avouac et al. 2014*; *Barnhart et al. 2014*) and uncertainty in Green's function is taken into account, which may account for the realistic

results obtained by using a model plane dipping 0° . It should also be noted that when we apply a method that is the same as our new method but neglecting uncertainty in Green's function, the obtained source model shows the focal mechanisms whose strike angles are inconsistent with the trend of both the north-east and the south-west sides of the surface-rupture trace measured by *Zinke et al. (2014)* (Supporting Information Fig. S12). When we analyse an earthquake with a wide rupture extent in depth, such as a subduction zone

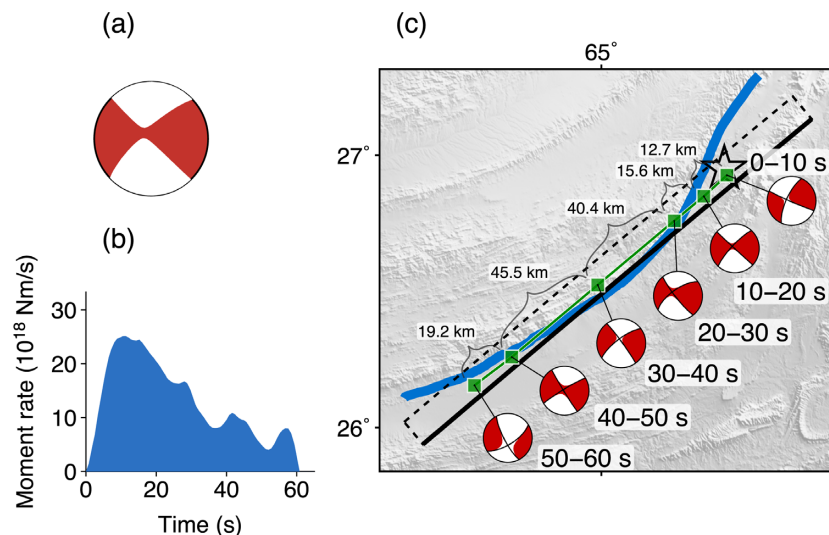


Figure 6. (a) Total moment tensor, (b) moment-rate function and (c) evolution of the centroid of the potency-rate-density tensors obtained by using the model plane dipping 64° . The star denotes the USGS epicentre. The dashed line outlines the model fault, and the bold black line shows the top of the model fault, respectively. The green squares and beach balls represent the centroid location and total potency-density tensor, respectively, each time window, along with the distance between neighbouring centroid locations and the corresponding time window for estimating the centroid location.

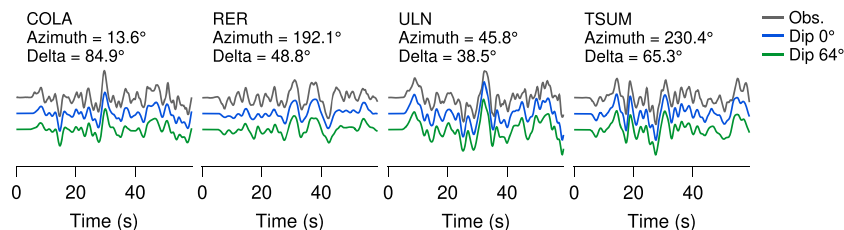


Figure 7. Fitting of waveforms between observed (top grey line) and synthetic waveforms for the model planes dipping 0° (middle blue line) and dipping 64° (bottom green line). The waveforms are plotted with a sampling interval of 0.05 s. The stations were selected to highlight the difference in the results of the fittings. Station code, azimuth and epicentral distance are shown above the traces.

earthquake, it is better to adopt a plausible dip angle of a model plane.

We evaluated the robustness of our new method by performing the synthetic tests, where 4 kinds of input source models were prepared; those were, the source model of the 2013 Balochistan earthquake, a simple reverse faulting model, a source model including drastic change of focal mechanism, and a strike-slip faulting model within a bending fault (Supporting Information Text S1). The results show that our new method can reproduce the drastic change of focal mechanism and the change of strike angle of the bending strike-slip fault during the rupture propagation without *ad hoc* assumptions of fault geometry (Supporting Information Figs S1–S9).

The 2.5 times larger number of model parameters in our new inversion scheme than the conventional one means a freedom from the constraint on shear-slip direction that is one of prior constraints implicitly imposed in conventional methods. Solutions obtained by our new method tend to be spatiotemporally smooth, and inferred focal mechanisms contain CLVD components even if a true source mechanism is a pure double couple, which can be seen in the synthetic tests (Supporting Information Figs S3a, b and S4).

The multiple-point-source inversions (e.g. Kikuchi & Kanamori 1991; Duputel *et al.* 2012) can estimate a single focal mechanism for each subevent with a simplified source-time function, but they can not resolve detailed rupture propagation and a change of focal mechanism within a subevent. A focal mechanism within a subevent

should vary in space and time, and the source-time function should be able to represent azimuth-dependent shape and duration, reflecting the directivity effects. Therefore, applying the multiple-point-source inversion is inappropriate for estimating detailed rupture propagation between subevents and resolving smoothly curved fault geometry, such as those of the 2013 Balochistan earthquake. Comparing with the multiple-point-source inversions, it is relatively difficult for our method to infer an abrupt change of focal mechanism because of the spatiotemporal smoothing constraints, however, our new finite-fault inversion framework can spatiotemporally resolve both the detailed rupture evolution and variation of the focal mechanism without non-negative slip constraint, and this high flexibility is due to the fact that uncertainty in Green's functions is explicitly taken into account.

4.2 Effects of a model plane geometry

For simplicity, we represented the spatiotemporal potency-rate-density distribution in a 2-D model space by assuming eq. (3), which represents the insensitivity of Green's function to a slight change of source location, and then neglected modelling errors originated from uncertainty in the model plane geometry. In this section, we evaluated the effects of the approximation of Green's function in eq. (3) and the robustness of our new method by changing the model plane geometry.

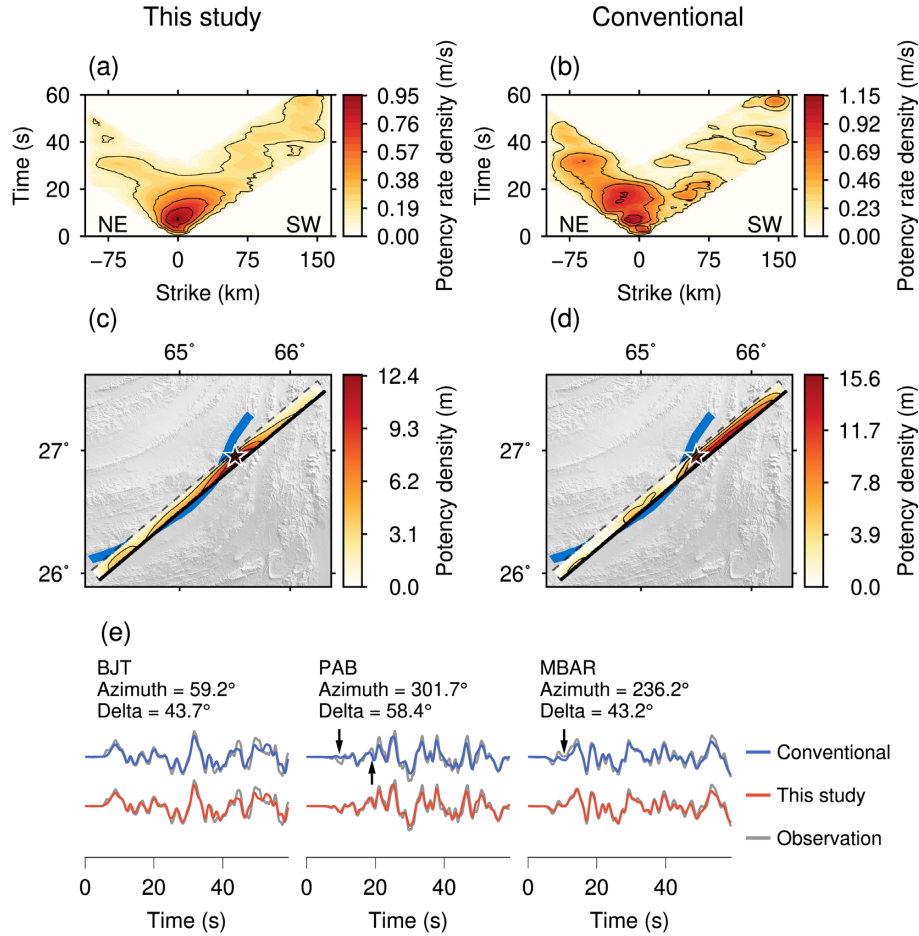


Figure 8. The difference between the models obtained by using the formulation developed in this study and the one used by Yagi & Fukahata (2011). (a, b) Time evolution of potency-rate density along the strike direction and (c, d) total potency-density distribution obtained with the new and conventional formulation, respectively. The star denotes the USGS epicentre. The dashed line outlines the model fault, and the bold black line shows the top of the model fault. Contour intervals in (a), (b), (c) and (d) are 0.19 m s^{-1} , 0.22 m s^{-1} , 3.1 m and 3.7 m , respectively. (e) Comparison of waveform fitting at selected stations. The waveforms are plotted with a sampling interval of 0.05 s . The station code, azimuth, and epicentral distance are shown above the traces. Arrows point out the phases where the observed waveform was not well reproduced by the conventional method.

We changed the dip angle of the model plane from 0° to 64° , the average dip angle of model planes used in the finite-fault inversion analysis conducted by USGS NEIC (<https://earthquake.usgs.gov/earthquakes/eventpage/usb000jyiv>, last accessed 19 November 2018), used the hypocentre determined by USGS NEIC as an initial rupture point, and adopted a model plane with a length of 200 km and a width of 25 km (Fig. 2b).

The total moment tensor shows strike-slip faulting with the more realistic fault plane having a 226° strike angle and a 77° dip angle (Fig. 6a); this dip angle is steeper than that of the more realistic fault plane of the total moment tensor obtained by the analysis using a model plane dipping 0° (Fig. 2c). The total released seismic moment was $7.00 \times 10^{20} \text{ Nm}$ (M_w 7.8), which is comparable to the analysis result obtained with the model plane dipping 0° . The source-time function has a major peak at around 11 s and three minor peaks at around 28 , 42 and 57 s (Fig. 6b and Supporting Information Fig. S14). Timings of the peaks in the source-time function are shifted in time within 1 – 5 s earlier than those of the previous result, which is originated from traveltime difference due to the different model plane geometries, but the number and relative amplitudes of the peaks are comparable to the previous result (Fig. 3 and Supporting Information Fig. S14). The focal

mechanisms during 0 to 30 s show overall similarity of the potency-density tensors between the model plane dipping 0° and that dipping 64° (Figs 4 and 6). The average centroid location migration velocity during 20 to 50 s increases to 4.3 km s^{-1} (Fig. 6), which is consistent with the previous result (Fig. 4). The strike angles of the potency-density tensors in the 10 s windows rotate clockwise as the rupture propagates towards the south-west from the epicentre until 50 s .

Although the rupture propagating unilaterally south-estward and the transition of focal mechanism showing the clockwise north-to-south rotation of the strike angle are consistent with the previous results obtained with the model plane dipping 0° , the potency-density tensors during 30 to 60 s (Fig. 6c) show steeper dip angles than those obtained by the analysis adopting with a model plane dipping 0° , and this difference in dip angle seems to become larger as the rupture evolves. The waveform fitting between the observed and synthetic waveforms is comparable between the two analyses (Fig. 7 and Supporting Information Fig. S10), but the variance is 0.04 larger in the analysis with the model plane dipping 64° . Because the dip angle of the model plane prescribes the source spatial location and the model plane dipping 0° horizontally covers the possible source region more broadly than that dipping 64° (Fig.

2b), modelling error due to inadequate assumption of the model-source location may be smaller for a model plane dipping 0° .

4.3 Bias due to limited number of basis double-couple components

To evaluate how the conventional inversion scheme suffers from the modelling error of the fault geometry, we applied the conventional method (Fig. 1b) to the Balochistan earthquake by assuming a flat model plane with a length of 260 km (extending the model plane 60 km farther north-eastward compared with our previous two analyses), and a width of 25 km, a strike of 230° and a dip of 64° (Fig. 2b). In our application of the conventional method, we employed the formulation of Yagi & Fukahata (2011), introducing the uncertainty in Green's functions into the data covariance matrix in the same way as in the method developed in this study, but adopting only the two basis components of potency tensor.

With the conventional inversion scheme, the resultant seismic moment was 7.27×10^{20} Nm (M_w 7.8), which is comparable to that obtained by the new method developed in this study (6.43×10^{20} Nm, M_w 7.8). However, the rupture propagation direction was opposite between the conventional and the new inversion schemes (Fig. 8). The conventional result showed the rupture propagating almost unilaterally north-eastward from the epicentre (Fig. 8b), with the dominant rupture area distributed in the area north-east of the source region (Fig. 8d). This north-eastward rupture propagation is unlikely to be realistic because, as shown in the previous sections, both the surface rupture mapped by the satellite image analyses and the backprojection result suggest that the rupture propagated towards south-west from the epicentre (e.g. Avouac *et al.* 2014). However, it should be noted that the variance of the waveform fittings obtained with the new method was only 0.04 less than that obtained with the conventional method, although the certain phases of the observed waveforms were not well reproduced at some stations in the conventional inversion scheme (Fig. 8e and Supporting Information Fig. S11). Even though the difference of the variance between the two methods seems small, the conventional method using only two basis components of potency tensor could not reproduce the realistic direction of rupture propagation, which was independently inferred from other data and analyses. We thus conclude that, even if we introduce the uncertainty in Green's functions into the covariance matrix, a simplified assumption of the fault geometry that limits the flexibility of slip directions may distort the solution (Fig. 8) and thereby lead to inappropriate interpretations of the source process. When we analyse seismic data with little information on fault geometry of an earthquake, such a single wide model plane, which may be inconsistent with the true one, is commonly adopted for constructing a preliminary source model. In this case, however, inappropriate models may be obtained due to modelling errors originated from uncertainty in fault geometry. Our new method enables us to mitigate such a problem, and to construct a realistic source model without detailed information about the fault geometry.

5 CONCLUSION

In this study, we proposed a new method of finite-fault inversion of teleseismic P waveforms in which shear slip along an assumed flat plane is represented by a superposition of five basis components of potency-density tensor. Because we extended the formulation of Yagi & Fukahata (2011), our method makes it possible not only to

estimate a fault geometry but also to mitigate effects of uncertainties of Green's function and fault geometry, which are two major sources of modelling errors. By applying our method to the 2013 Balochistan earthquake, we confirmed that the estimated direction of rupture propagation and the spatiotemporal distribution of focal mechanisms were consistent with those independently inferred from the backprojection result and the optical satellite images. This result suggests that our method is applicable to an earthquake that involves complex fault geometry and makes it possible to extract the information of the fault geometry, which contributes to a reliable source modelling.

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SUPPORTING INFORMATION

Supplementary data are available at [GJI](https://doi.org/10.1093/gji/ggy1012) online.

Text S1. Synthetic tests.

Figures S1–S9. Results of synthetic tests.

Figure S10. Waveform fitting at all stations between observed and synthetic waveforms for the model planes dipping 0° and dipping 64° .

Figure S11. Waveform fitting at all stations between observed and synthetic waveforms for the source models obtained by a conventional method and our method.

Figure S12. Source models where uncertainty of Green's function is neglected and introduced.

Figure S13. Evolution of the centroid of the potency-rate-density tensors.

Figure S14. Source-time functions for the model planes dipping 0° and 64° .

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