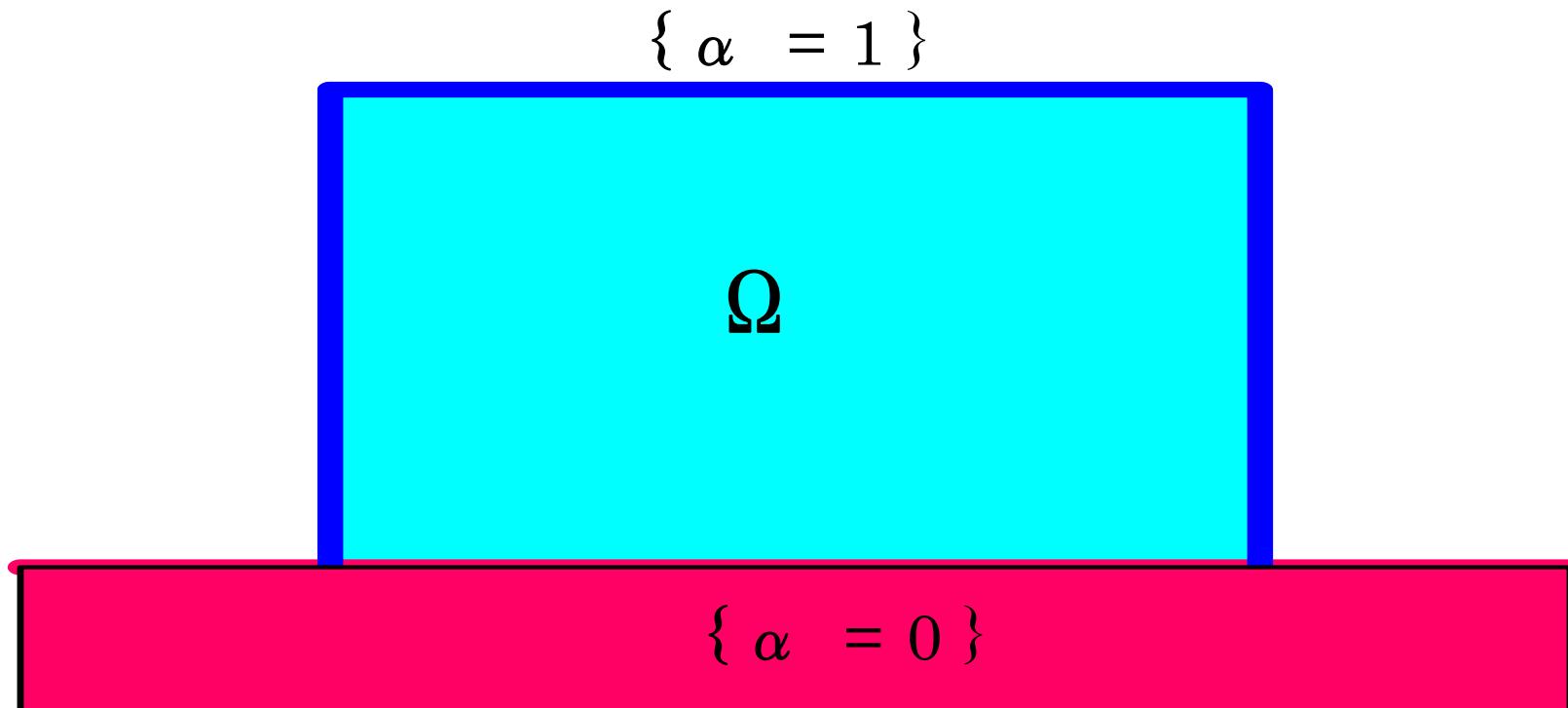


Boundary Value Problems of Nonlinear Elastostatics

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Domain with Corner Singularities



Abstract

- This talk is devoted to a **semigroup approach** to an initial-boundary value problem of **nonlinear elastodynamics** in the case where the boundary condition is a **regularization** of the genuine mixed displacement-traction boundary condition.

My Works

- **Taira:** On boundary value problems of nonlinear elastostatics, *Osaka Journal of Mathematics*, 33 (1996), 555-585.
<http://projecteuclid.org/euclid.ojm/1200786927>
- **Taira:** Introduction to boundary value problems of nonlinear elastostatics, *Tsukuba Journal of Mathematics*, 32 (2008), 67-138
- **Taira :** A mixed problem of linear elastodynamics, *Journal of Evolution Equations*, 13 (2013), 481-507

DOI: 10.1007/s00028-013-0187-1

References (1)

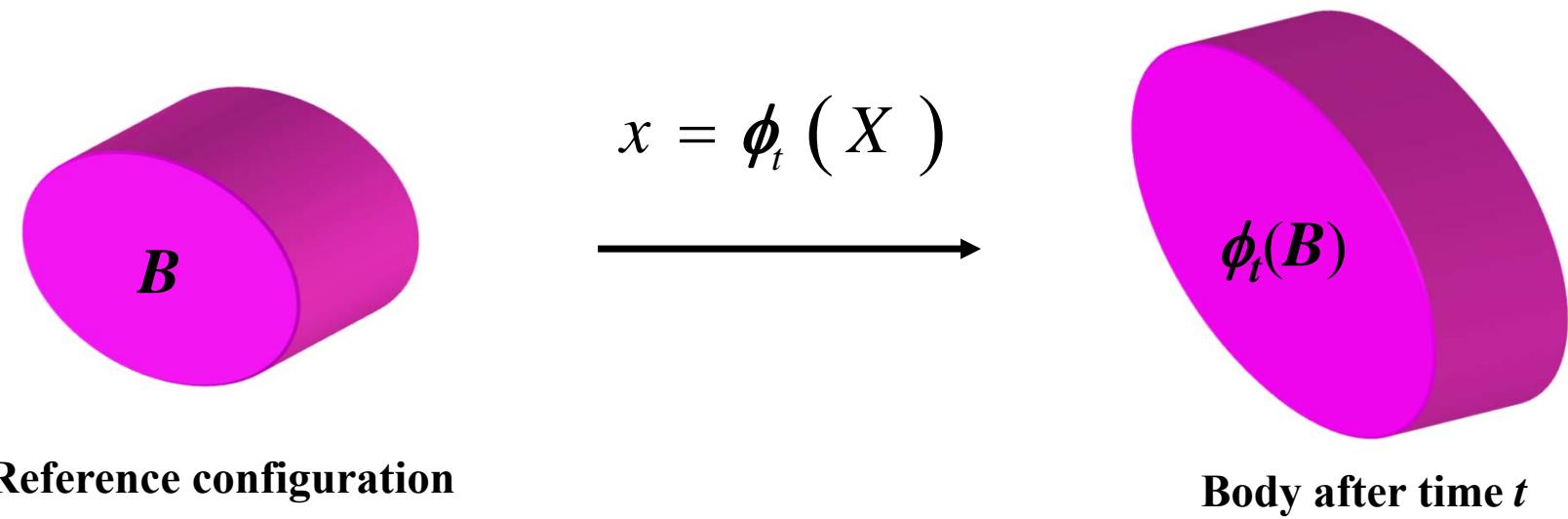
- **Ciarlet**: Mathematical elasticity, North-Holland, 1988
- **Duvaut et Lions**: Les inéqations en mécanique et en physique, Dunod, 1972
- **Marsden and Hughes**: Mathematical foundations of elasticity, Prentice-Hall, 1983
- **Valent**: Boundary value problems of finite elasticity, Springer-Verlag, 1988

References (2)

- **Goldstein**: Semigroups of linear operators and applications, Oxford University Press, 1985
- **Weiss**: Abstract vibrating systems, J. Math. Mech. 17 (1967), 241—255.

Formulation of the Problem

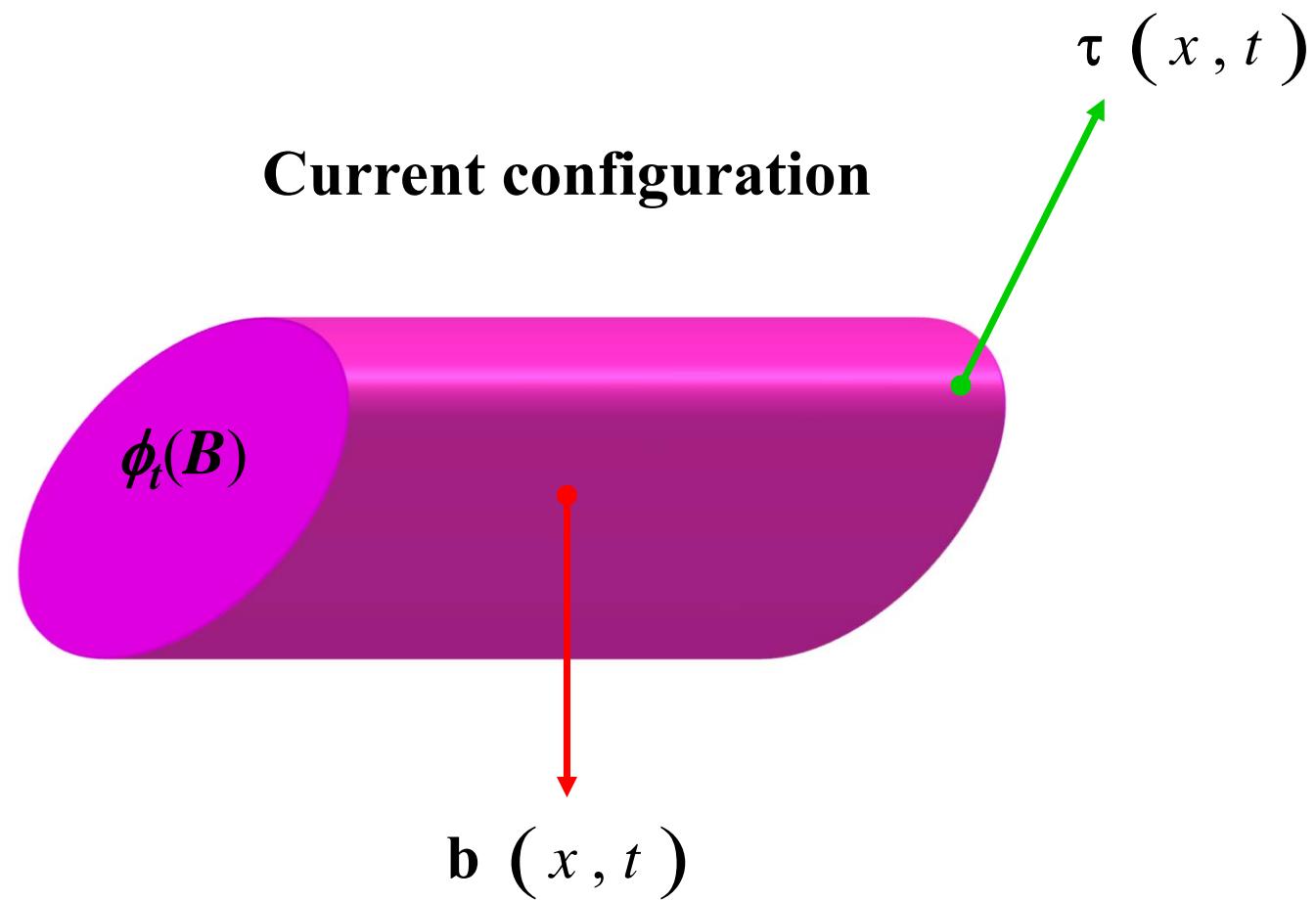
Nonlinear Elasticity



Notation

Notation	Meaning
Ω	Domain in \mathbf{R}^3
$B = \overline{\Omega}$	Reference configuration
ϕ_t	Deformation of B
$X = (X_1, X_2, X_3)$	Material point
$x = (x_1, x_2, x_3)$	Spatial point

Eulerian Description (1)

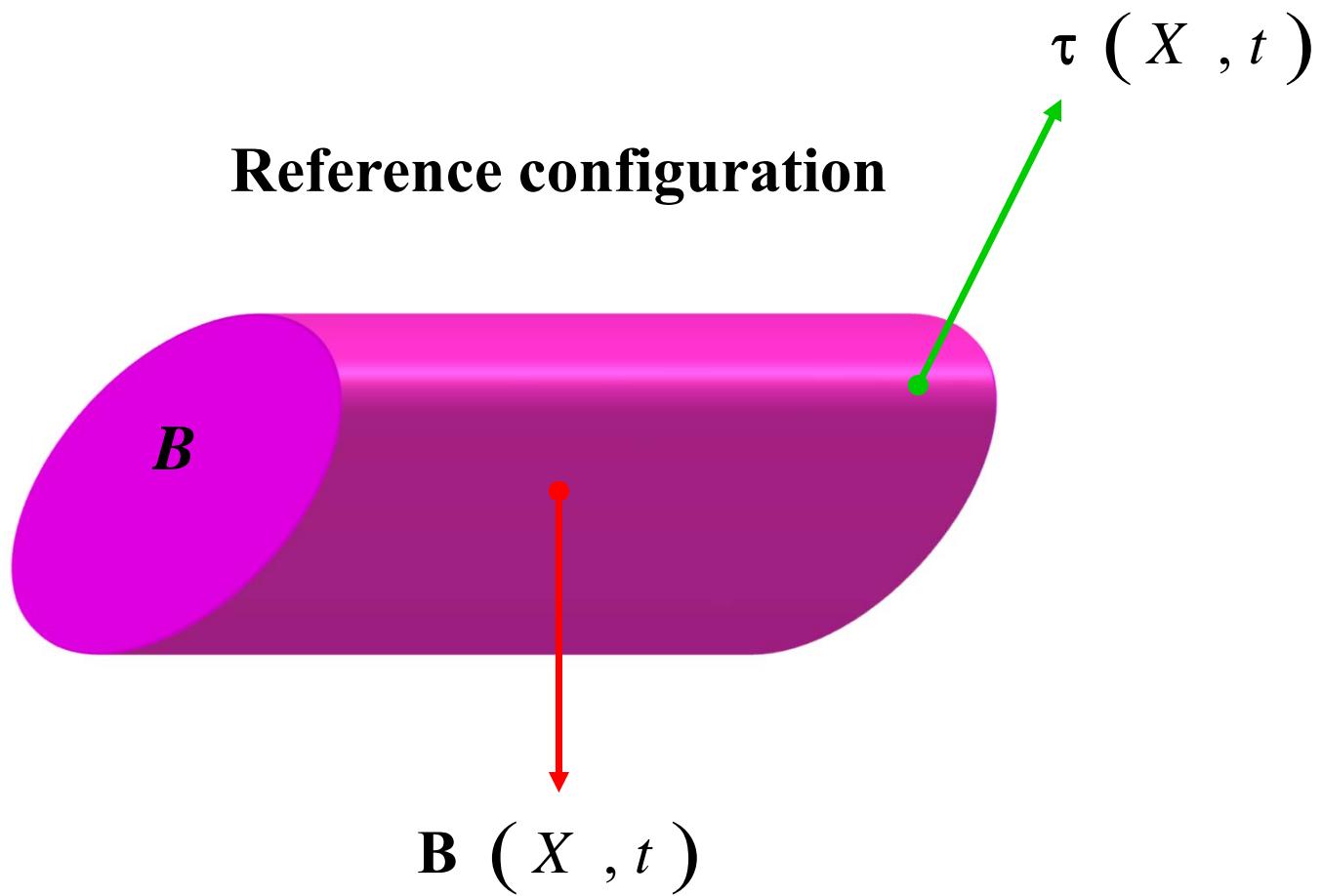


Leonhard Euler

Leonhard Euler (1707-1783)
**Swiss Mathematician, Physicist, Astronomer
and Engineer**

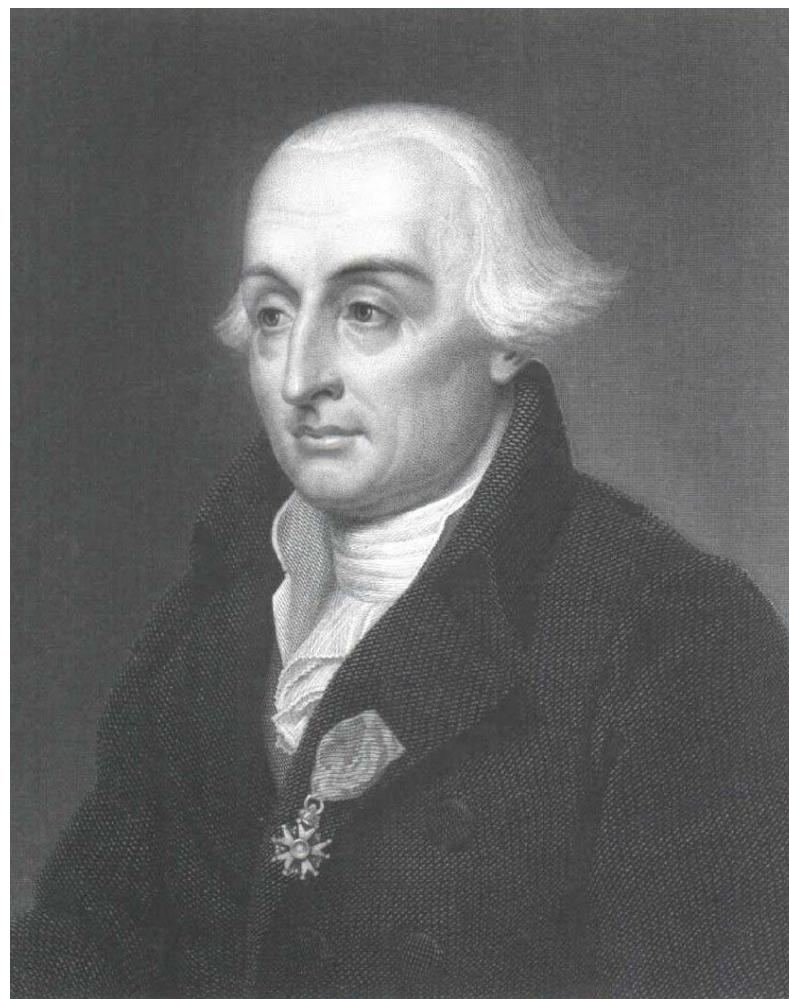


Lagrangian Description



Joseph-Louis Lagrange

Joseph-Louis Lagrange (1736-1813)
Italian Mathematician and Astronomer



Continuum Mechanics (1)

Form	Conservation of mass	Balance of linear momentum
Eulerian form	$\dot{\rho} + \rho \operatorname{div} \mathbf{v} = 0$	$\rho \dot{\mathbf{v}} = \operatorname{div} \boldsymbol{\sigma} + \rho \mathbf{b}$
Lagrangian form	$\rho_0(x) = \rho(\phi_t(x), t) J(x, t)$	$\rho_0 \frac{\partial \mathbf{V}}{\partial t} = \operatorname{Div} \mathbf{P} + \rho_0 \mathbf{B}$

Divergence of a Tensor Filed

$$\mathbf{T} = \begin{pmatrix} T_{11} & T_{12} & T_{13} \\ T_{21} & T_{22} & T_{23} \\ T_{31} & T_{32} & T_{33} \end{pmatrix}$$

\Rightarrow

$$\text{div } \mathbf{T} = \begin{pmatrix} \partial_1 T_{11} + \partial_2 T_{12} + \partial_3 T_{13} \\ \partial_1 T_{21} + \partial_2 T_{22} + \partial_3 T_{23} \\ \partial_1 T_{31} + \partial_2 T_{32} + \partial_3 T_{33} \end{pmatrix}$$

The divergence theorem for a tensor field

$$\int_{\Omega} \operatorname{div} \mathbf{T} \, dx = \int_{\partial\Omega} \mathbf{T} \cdot \mathbf{n} \, d\sigma$$

Notation (1)

Notation	Meaning
$\rho_0(X), \rho(x, t)$	Density
$\mathbf{V}(X, t), \mathbf{v}(x, t)$	Velocity
$\mathbf{B}(X, t), \mathbf{b}(x, t)$	Body force
$\sigma(x, t)$	Cauchy stress tensor
$\mathbf{P}(X, t)$	First Piola-Kirchhoff stress tensor

Louis Augustin Cauchy

**Louis Augustin Cauchy (1789-1857)
French Mathematician and Physicist**



Notation (2)

Notation	Meaning
$\phi_t(X)$	Deformation
$\mathbf{F}(X, t) = D\phi_t(X)$	Deformation gradient
$J(X, t) = \det \mathbf{F}(X, t)$	Jacobian
$\bullet = \frac{D}{Dt} = \frac{\partial}{\partial t} + \mathbf{v} \bullet \nabla$	Material derivative

Deformation Gradient

$$\boldsymbol{\varphi} = \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \varphi_3 \end{pmatrix} : \overline{\Omega} \rightarrow \mathbf{R}^3$$

\Rightarrow

$$\nabla \boldsymbol{\varphi} = \begin{pmatrix} \partial_1 \varphi_1 & \partial_2 \varphi_1 & \partial_3 \varphi_1 \\ \partial_1 \varphi_2 & \partial_2 \varphi_2 & \partial_3 \varphi_2 \\ \partial_1 \varphi_3 & \partial_2 \varphi_3 & \partial_3 \varphi_3 \end{pmatrix}$$

Displacement gradient

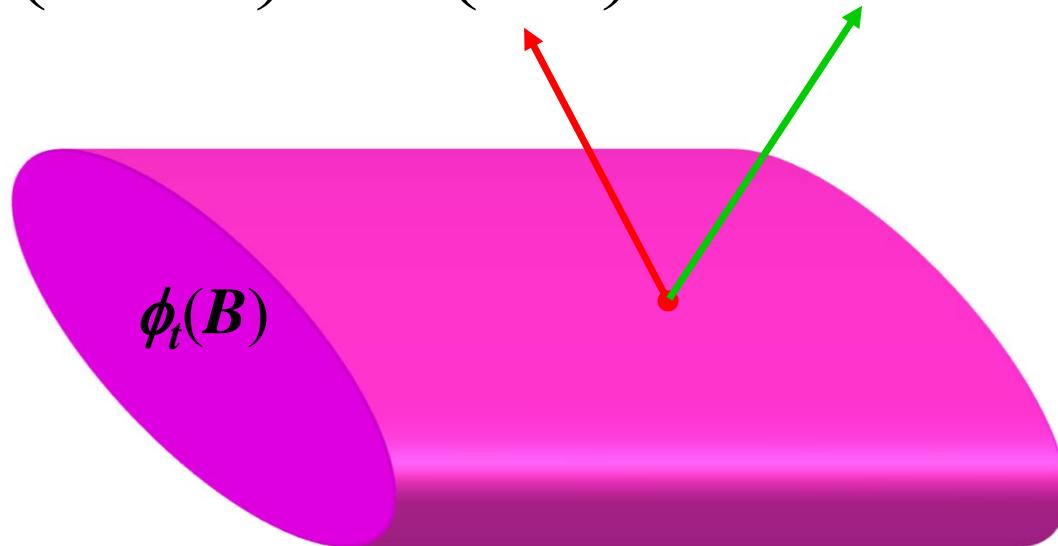
$$\boldsymbol{u} = \begin{pmatrix} u_1 \\ u_2 \\ u_3 \end{pmatrix} = \begin{pmatrix} \varphi_1 - x_1 \\ \varphi_2 - x_2 \\ \varphi_3 - x_3 \end{pmatrix} : \overline{\Omega} \rightarrow \mathbf{R}^3$$

⇒

$$\nabla \boldsymbol{u} = \begin{pmatrix} \partial_1 u_1 & \partial_2 u_1 & \partial_3 u_1 \\ \partial_1 u_2 & \partial_2 u_2 & \partial_3 u_2 \\ \partial_1 u_3 & \partial_2 u_3 & \partial_3 u_3 \end{pmatrix} = \nabla \boldsymbol{\varphi} - \mathbf{I}$$

Eulerian Description (2)

$$\mathbf{t} (x, t, \mathbf{n}) = \boldsymbol{\sigma} (x, t) \bullet \mathbf{n}$$



Current configuration

Continuum Mechanics (2)

Form	Cauchy's Theorem
Eulerian form	$\mathbf{t} (x, t, \mathbf{n}) = \boldsymbol{\sigma} (x, t) \bullet \mathbf{n}$
Lagrangian form	$\mathbf{T} (X, t, \mathbf{N}) = \mathbf{P} (X, t) \bullet \mathbf{N}$

Notation (3)

Notation	Meaning
$t(x, t, n)$	Cauchy stress vector
$\sigma(x, t)$	Cauchy stress tensor
$T(X, t, N)$	First Piola-Kirchhoff stress vector
$P(X, t)$	First Piola-Kirchhoff stress tensor
n, N	Outward unit normal vector

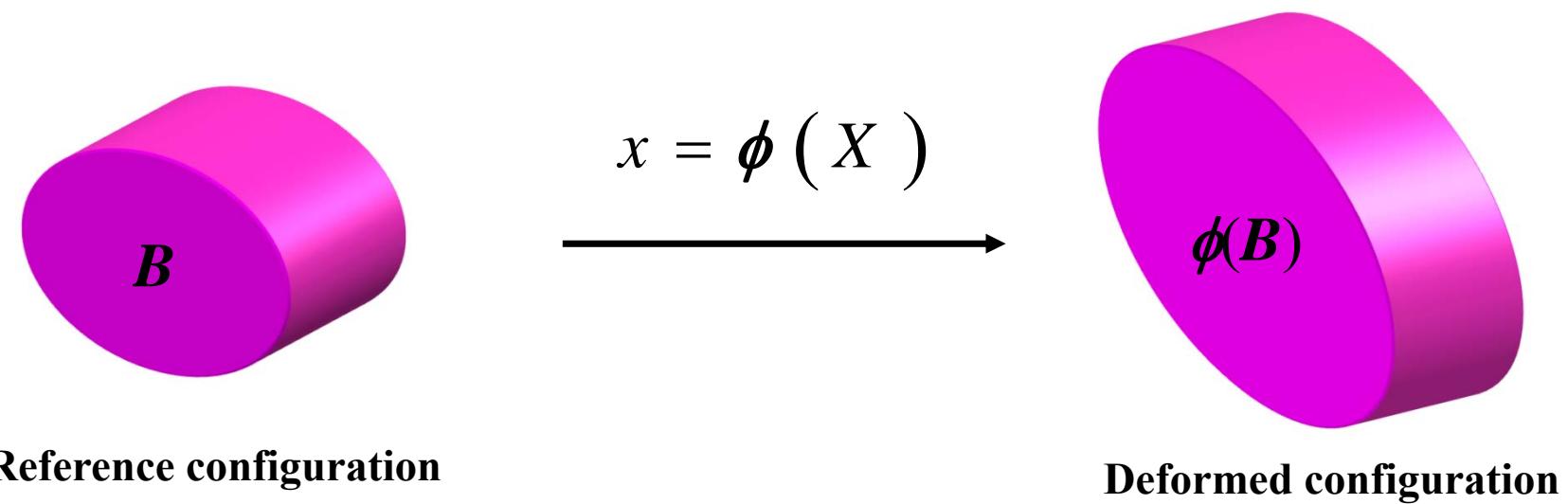
Continuum Mechanics (3)

Form	Balance of angular momentum	Conservation of energy
Eulerian form	$\sigma = {}^t \sigma$	$\rho \dot{e} + \text{div } \mathbf{q} = \text{tr}(\sigma \mathbf{d}) + \rho r$
Lagrangian form	$\mathbf{S} = {}^t \mathbf{S}$	$\rho_0 \frac{\partial E}{\partial t} + \text{D iv } \mathbf{Q} = \text{tr}(\mathbf{S} \mathbf{D}) + \rho_0 R$

Notation (4)

Notation	Meaning
$\mathbf{S}(X, t)$	Second Piola-Kirchhoff stress tensor
$E(X, t), e(x, t)$	Internal energy function
$\mathbf{D}(X, t), \mathbf{d}(x, t)$	Material rate of deformation tensor
$R(X, t), r(x, t)$	Heat supply
$\mathbf{Q}(X, t), \mathbf{q}(x, t)$	Heat flux vector

Nonlinear Elastostatics



Notation

Notation	Meaning
Ω	Bounded domain in \mathbf{R}^3 with smooth boundary
$B = \overline{\Omega}$	Reference configuration
ϕ	Deformation of B
$X = (X_1, X_2, X_3)$	Material point
$x = (x_1, x_2, x_3)$	Spatial point

Elasticity

A material is said to be **elastic**

\Leftrightarrow

\exists a function $\hat{\mathbf{P}}(X, \mathbf{F})$ of points $X \in B$
and 3×3 matrices $\mathbf{F} = (F_{ij})$ with $\det \mathbf{F} > 0$
such that

$$\boxed{\mathbf{P}(X) = \hat{\mathbf{P}}(X, D\phi(X))}$$

Constitutive Function

$\hat{P}(X, F)$ is called a **constitutive function**

Hyper-Elasticity

A material is said to be **hyperelastic**

\Leftrightarrow

\exists A function $W(X, \mathbf{F})$ of points $X \in B$
and 3×3 matrices $\mathbf{F} = (F_{ij})$ with $\det \mathbf{F} > 0$
such that

$$\widehat{\mathbf{P}}(X, \mathbf{F}) = \frac{\partial W}{\partial \mathbf{F}}(X, \mathbf{F})$$

$$\boxed{\widehat{P}_{ij}(X, \mathbf{F}) = \frac{\partial W}{\partial F_{ij}}(X, \mathbf{F})}$$

Stored Energy Function

$W(X, F)$ is called a **stored energy function**

First Elasticity Tensor

The **four - index tensor**
defined by the formula

$$\mathbf{A} = \frac{\partial \widehat{\mathbf{P}}}{\partial \mathbf{F}} = \frac{\partial^2 W}{\partial \mathbf{F} \partial \mathbf{F}},$$

$$\begin{aligned} \mathbf{A}_{ij\ell m}(X, \mathbf{F}) &= \frac{\partial \widehat{P}_{ij}}{\partial F_{\ell m}}(X, \mathbf{F}) \\ &= \frac{\partial^2 W}{\partial F_{ij} \partial F_{\ell m}}(X, \mathbf{F}) \end{aligned}$$

is called the **first elasticity tensor**.

Typical Examples

Hencky-Nadai Elasto-Plastic Material

The **stored energy function** has the form

$$W(X, \mathbf{F}) = \frac{3}{4} \int_0^{\Gamma(\mathbf{F})} g(\xi) d\xi + \frac{K}{2} \left(\sum_{k=1}^3 F_{kk} - 3 \right)^2$$

$$g(\xi) \in C^\infty([0, \infty), \mathbf{R})$$

$$\Gamma(\mathbf{F}) = \frac{4}{3} \sum_{i,j=1}^3 \left(\frac{1}{2} (F_{ij} + F_{ji}) - \frac{1}{3} \left(\sum_{k=1}^3 F_{kk} \right) \delta_{ij} \right)^2$$

The Modulus of Compression

K is called the **modulus of compression**

Saint Venant-Kirchhoff Isotropic Material

The **stored energy function** has the form

$$W(X, \mathbf{F})$$

$$= \frac{\lambda(X)}{8} \left(\sum_{k=1}^3 C_{kk}(\mathbf{F}) - 3 \right)^2 + \frac{\mu(X)}{4} \sum_{i,j=1}^3 (C_{ij}(\mathbf{F}) - \delta_{ij})^2$$

$$C_{ij}(\mathbf{F}) = \sum_{k=1}^3 F_{ki} F_{kj}$$

Saint Venant

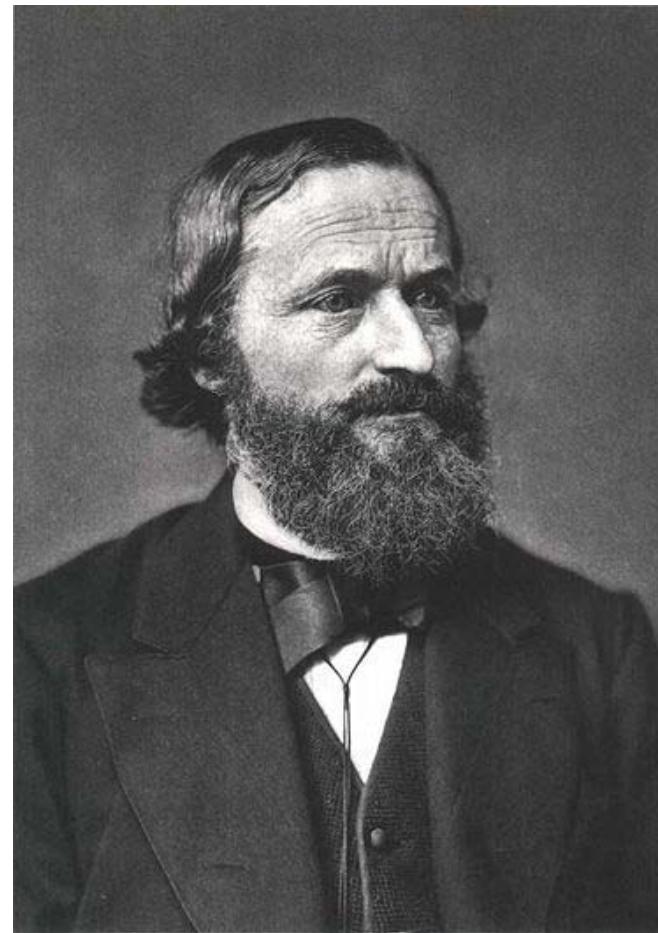
**Adhémar Jean Claude Barré de Saint Venant
(1797-1886)**

French Mechanician and Mathematician



Gustav Robert Kirchhoff

Gustav Robert Kirchhoff (1824-1887)
Prussian/German Physicist and Chemist



Lamé functions

$\lambda(X), \mu(X)$ are called **Lame functions**

The Right Cauchy-Green tensor

$$C_{ij}(\mathbf{F}) = \sum_{k=1}^3 F_{ki} F_{kj}$$

Green-Saint Venant strain tensor

$$\mathbf{E} = \frac{1}{2} (\mathbf{C} - \mathbf{I})$$

$$\mathbf{E} = (E_{ij})$$

$$= \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} + \sum_{k=1}^3 \frac{\partial u_k}{\partial x_i} \frac{\partial u_k}{\partial x_j} \right)$$

Saint Venant-Kirchhoff Isotropic Material

$$W(X, F) = \frac{\lambda(X)}{2} (\operatorname{tr} E)^2 + \mu(X) \operatorname{tr}(E^2)$$

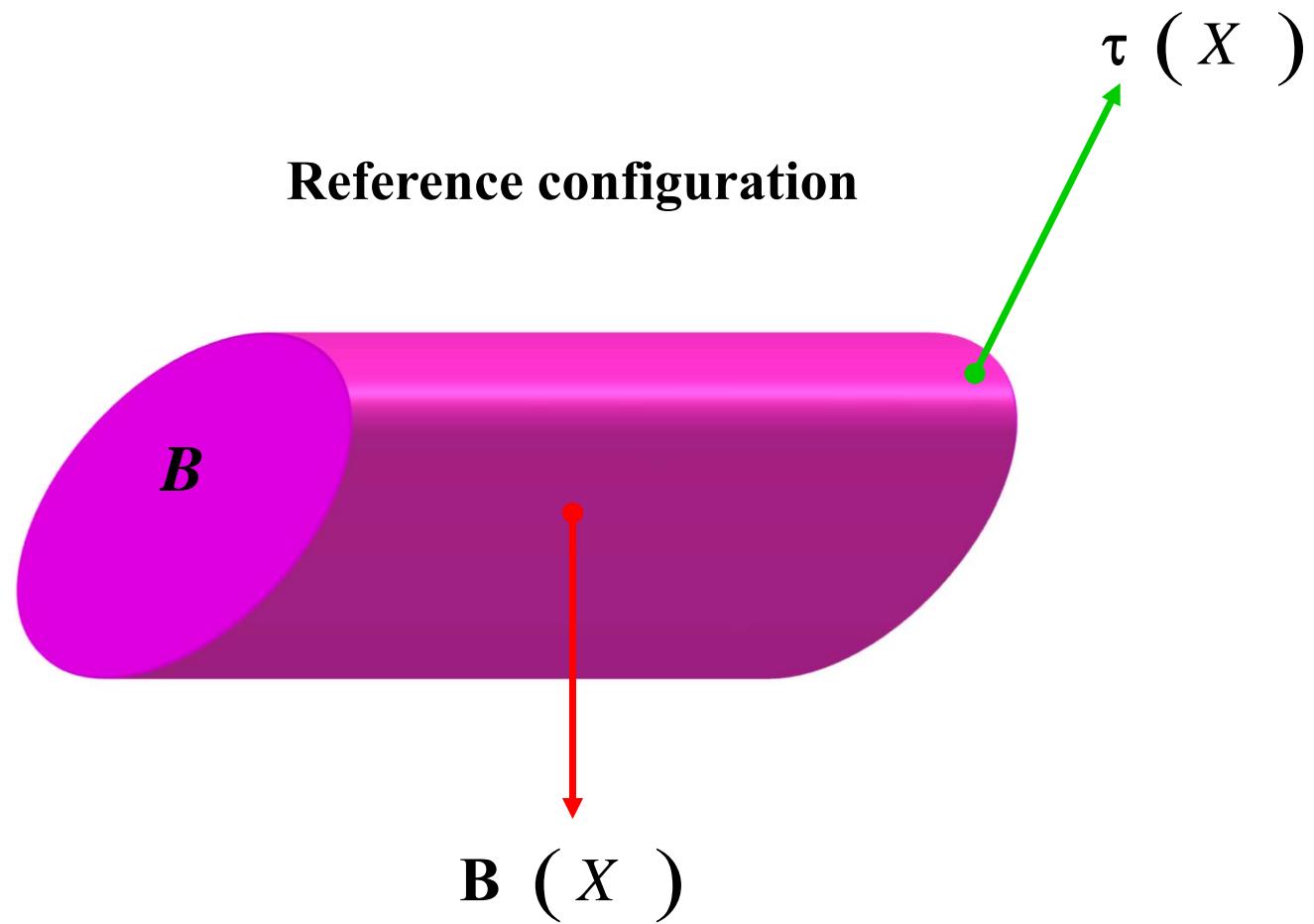
Formulation of the Results

Nonlinear Boundary Value Problem (Lagrange)

This talk is devoted to an L^p approach to the following equilibrium equations for the unknown deformation $\phi(X)$

$$\begin{aligned} \text{Div } \widehat{\mathbf{P}}(X, D\phi(X)) + \mathbf{B}(X) &= \mathbf{0} \quad \text{in } \Omega, \\ \alpha(X)\widehat{\mathbf{P}}(X, D\phi(X)) \cdot \mathbf{N}(X) + (1 - \alpha(X))\phi(X) \\ &= \tau(X) \quad \text{on } \partial\Omega. \end{aligned}$$

Lagrangian Description



Notation

Notation	Meaning
$\widehat{\mathbf{P}}(X, \mathbf{F})$	Constitutive function
$\mathbf{B}(X)$	Body force
$\tau(X)$	Surface force
$\mathbf{N}(X)$	Outward unit normal
$\alpha(X)$	$\alpha(X) \in C^\infty(\partial\Omega)$
	$0 \leq \alpha(X) \leq 1 \text{ on } \partial\Omega$

Boundary Condition

$\{\alpha = 1\}$ (traction)

A diagram showing a light blue rectangular domain labeled Ω . The top and right edges of the rectangle are highlighted in blue, representing traction boundary conditions. The bottom and left edges are highlighted in red, representing displacement boundary conditions.

$\{\alpha = 0\}$ (displacement)

Example of Regularization (Real axis)

$$\alpha(x_1) = \begin{cases} e^{-\frac{1}{x_1}} & \text{for } x_1 > 0, \\ 0 & \text{for } x_1 \leq 0 \end{cases}$$

Example of Regularization (Unit Circle)

$$a(x_1, x_2) = a(\cos \theta, \sin \theta) =$$
$$\begin{cases} \exp[2/\pi + 1/(\theta + \pi/2)](1 - \exp[2/\pi - 1/(\theta + \pi)]), \\ \theta \in (-\pi, -\pi/2) \\ 0, \theta \in [-\pi/2, 0] \\ \exp[2/\pi - 1/\theta](1 - \exp[2/\pi + 1/(\theta - \pi/2)]), \\ \theta \in (0, \pi/2), \\ 1, \theta \in [\pi/2, \pi] \end{cases}$$

Main Results

1. If the linearized problem has unique solutions, then so does the nonlinear one, nearby. This is done by using the **linear L^p theory** and the **inverse mapping theorem**.
2. Our results can be applied to the **Saint Venant-Kirchhoff** elastic material and the **Hencky-Nadai** elasto-plastic material.

Crucial Point

- The crucial point is how to find a **function space** associated with the **degenerate** boundary condition in which the linearized problem has unique solutions.

Function Spaces (1)

For all real s , we define **Sobolev spaces**

$$H^{s,p}(\Omega, \mathbf{R}^3) = H^{s,p}(\mathbf{R}^3, \mathbf{R}^3)|_{\Omega}$$

For all $s > 1/p$, we define **Besov spaces**

$$B^{s-1/p,p}(\partial\Omega, \mathbf{R}^3) = H^{s,p}(\Omega, \mathbf{R}^3)|_{\partial\Omega}.$$

Function Spaces (2)

$$s > 1 + 1/p$$

$$\mathbf{B}_{(\alpha)}^{s-1-1/p, p} (\partial\Omega, \mathbf{R}^3)$$

$$= \alpha(X) \mathbf{B}^{s-1-1/p, p} (\partial\Omega, \mathbf{R}^3) + (1 - \alpha(X)) \mathbf{B}^{s-1/p, p} (\partial\Omega, \mathbf{R}^3)$$

$$(A) \quad 0 \leq \alpha(X) \leq 1 \text{ and } \alpha(X) \neq 1 \text{ on } \partial\Omega.$$

Definition of a norm

$$\begin{aligned} & |\varphi|_{\mathbf{B}_{(\alpha)}^{s-1-1/p, p}(\partial\Omega, \mathbf{R}^3)} \\ &= \inf \left\{ |\varphi_1|_{s-1-1/p, p} + |\varphi_2|_{s-1/p, p} : \varphi = \alpha(X)\varphi_1 + (1 - \alpha(X))\varphi_2 \right\} \end{aligned}$$

$$(A) \quad 0 \leq \alpha(X) \leq 1 \text{ and } \alpha(X) \not\equiv 1 \text{ on } \partial\Omega.$$

Function Spaces (3)

$$\mathbf{B}_{(\alpha)}^{s-1-1/p, p}(\partial\Omega, \mathbf{R}^3) = \mathbf{B}^{s-1/p, p}(\partial\Omega, \mathbf{R}^3)$$

if $\alpha(X) \equiv 0$ on $\partial\Omega$ (**Dirichlet**)

$$\mathbf{B}_{(\alpha)}^{s-1-1/p, p}(\partial\Omega, \mathbf{R}^3) = \mathbf{B}^{s-1-1/p, p}(\partial\Omega, \mathbf{R}^3)$$

if $\alpha(X) \equiv 1$ on $\partial\Omega$ (**Neumann**)

Function Spaces (4)

$$X = \mathbf{H}^{s,p}(\Omega, \mathbf{R}^3),$$

$$Y = \mathbf{H}^{s-2,p}(\Omega, \mathbf{R}^3) \times \mathbf{B}_{(\alpha)}^{s-1-1/p, p}(\partial\Omega, \mathbf{R}^3).$$

Nonlinear Map

$$F : X \rightarrow Y$$

$$\boxed{F(\phi) = \left(-\operatorname{Div} \hat{\mathbf{P}}(D\phi), \alpha \hat{\mathbf{P}}(D\phi) \cdot \mathbf{N} + (1-\alpha)\phi \Big|_{\partial\Omega} \right)}$$

Fundamental Assumptions (1)

(A) $0 \leq \alpha(X) \leq 1$ and $\alpha(X) \not\equiv 1$ on $\partial\Omega$.

(P) $\overset{o}{\mathbf{P}} \left(D \overset{o}{\phi} \right) = 0$ when $\overset{o}{\phi} = \mathbf{I}_\Omega$.

Remarks

1. Condition (P) implies that the reference configuration is a **natural state**.
2. Condition (A) implies that our boundary condition is not equal to the **pure traction** boundary condition

Fundamental Assumptions (2)

$$(H) \quad \overset{\circ}{A} = \frac{\partial^2 W}{\partial F \partial F} \quad \text{at } \overset{\circ}{\phi} = I_{\Omega}$$

enjoys the property of **symmetry**

$$\overset{\circ}{A}_{ij\ell m} = \overset{\circ}{A}_{\ell m ij} = \overset{\circ}{A}_{j i \ell m}$$

and is **uniformly pointwise stable**

$$\frac{1}{2} \mathbf{e} \cdot \overset{\circ}{A} \cdot \mathbf{e} \geq \exists \eta \ \| \mathbf{e} \|$$

Main Theorem (1)

Main Theorem states that if the linearized problem is uniformly pointwise stable, then, for **slight perturbations** of the load or boundary conditions from their values at the natural state, then the nonlinear problem has a **unique solution**.

Main Theorem (2)

$$1 < p < \infty, \quad s > 3/p + 1.$$

\exists A neighborhood U of $\overset{\circ}{\phi} = \mathbf{I}_\Omega$

and a neighborhood V of the point

$$\left(-\operatorname{Div} \widehat{\mathbf{P}}\left(D\overset{\circ}{\phi}\right), \alpha \widehat{\mathbf{P}}\left(D\overset{\circ}{\phi}\right) \bullet \mathbf{N} + (1 - \alpha)\overset{\circ}{\phi} \Big|_{\partial\Omega} \right)$$

such that the **nonlinear map**

$$F : U \rightarrow V$$

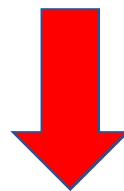
is **bijective**.

Typical Examples

Hencky-Nadai Elasto-Plastic Material

- (A) $0 \leq \alpha(X) \leq 1$ and $\alpha(X) \not\equiv 1$ on $\partial\Omega$.

(G) $g(0) > 0$ and $K > 0$



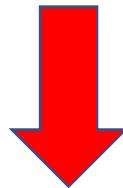
**Condition (H) is satisfied and
so Main Theorem applies.**

Saint Venant-Kirchhoff Isotropic Material

(A) $0 \leq \alpha(X) \leq 1$ and $\alpha(X) \not\equiv 1$ on $\partial\Omega$.

(M) $\exists c_1 > 0, \exists c_2 > 0$ such that

$$\begin{cases} \mu(X) \geq c_1 \text{ on } \Omega, \\ \lambda(X) + \frac{2}{3}\mu(X) \geq c_2 \text{ on } \Omega \end{cases}$$



**Condition (H) is satisfied and
so Main Theorem applies.**

Linearized Problems (Frechet Derivatives)

Mixed Problem of Linear Elastodynamics

We study the following mixed problem of
linear elastodynamics:

$$\frac{\partial^2 \mathbf{u}}{\partial t^2} - \operatorname{div}(\mathbf{a} \cdot \nabla \mathbf{u}) = \mathbf{f} \quad \text{in } \Omega \times (0, \infty),$$

$$\mathbf{u} \Big|_{t=0} = \mathbf{u}_0 \quad \text{in } \Omega,$$

$$\frac{\partial \mathbf{u}}{\partial t} \Big|_{t=0} = \mathbf{u}_1 \quad \text{in } \Omega,$$

$$\alpha(x')(\mathbf{a} \cdot \nabla \mathbf{u} \cdot \mathbf{n}) + (1 - \alpha(x'))\mathbf{u} = \mathbf{0} \quad \text{on } \partial\Omega \times (0, \infty).$$

Example (I)

$$\frac{\partial^2 \mathbf{u}}{\partial t^2} - \operatorname{div} (\lambda \operatorname{tr}(\mathbf{e}(\mathbf{u})) \mathbf{I} + 2\mu \mathbf{e}(\mathbf{u})) = \mathbf{f} \quad \text{in } \Omega \times (0, \infty),$$

$$\mathbf{u} \Big|_{t=0} = \mathbf{u}_0 \quad \text{in } \Omega,$$

$$\left. \frac{\partial \mathbf{u}}{\partial t} \right|_{t=0} = \mathbf{u}_1 \quad \text{in } \Omega,$$

$$\alpha(x')(\boldsymbol{\tau}(\mathbf{u}) \cdot \mathbf{n}) + (1 - \alpha(x'))\mathbf{u} = \mathbf{0} \quad \text{on } \partial\Omega \times (0, \infty).$$

Example (I)

$$\mathbf{e}(\mathbf{u}) = (e_{ij}) = \left(\frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \right), \text{ strain tensor}$$

$$\boldsymbol{\tau}(\mathbf{u}) = (\tau_{ij}) = (\lambda e_{kk} \delta_{ij} + 2 \mu e_{ij}), \text{ stress tensor}$$

Example (Ia)

$$(a_{ij\ell m}) = (\lambda \delta_{ij} \delta_{\ell m} + \mu (\delta_{i\ell} \delta_{jm} + \delta_{im} \delta_{j\ell}))$$

λ, μ : Lame moduli

$$(H) \quad \mu > 0, \quad k = \frac{3\lambda + 2\mu}{3} > 0.$$

k : modulus of compression

Example (Ib)

$$(a_{ij\ell m}) = \left(\left(K - \frac{2}{3} g(0) \right) \delta_{ij} \delta_{\ell m} + g(0) (\delta_{i\ell} \delta_{jm} + \delta_{im} \delta_{j\ell}) \right)$$

K is the **modulus of compression**

$$g(\xi) \in C^\infty([0, \infty), \mathbf{R})$$

$$(G) \quad K > 0, \quad g(0) > 0.$$

Example (II)

$$\frac{\partial^2 u}{\partial t^2} - \Delta u = f \quad \text{in } \Omega \times (0, \infty),$$

$$u \Big|_{t=0} = u_0 \quad \text{in } \Omega,$$

$$\frac{\partial u}{\partial t} \Big|_{t=0} = u_1 \quad \text{in } \Omega,$$

$$\alpha(x') \frac{\partial u}{\partial \mathbf{n}} + (1 - \alpha(x'))u = 0 \quad \text{on } \partial \Omega \times (0, \infty).$$

Example (II)

$$(a_{ij\ell m}) = (\delta_{i\ell} \delta_{jm}).$$

Linear Boundary Value Problem

We are reduced to the study of a problem of **linear elastostatics** for the unknown vector function \mathbf{v} :

$$\mathbf{A}\mathbf{v} = \operatorname{div}(\mathbf{a} \cdot \nabla \mathbf{v}) = \mathbf{f} \quad \text{in } \Omega,$$

$$\mathbf{B}_\alpha \mathbf{v} = \alpha(x') (\mathbf{a} \cdot \nabla \mathbf{v} \cdot \mathbf{n}) + (1 - \alpha(x')) \mathbf{v} = \varphi \quad \text{on } \partial\Omega$$

$\mathbf{a} = (a_{ij\ell m})$: smooth elasticity tensor.

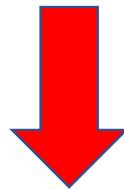
\mathbf{n} : the outward unit normal

Fundamental Existence and Uniqueness Theorem

(A) $0 \leq \alpha(X) \leq 1$ and $\alpha(X) \not\equiv 1$ on $\partial\Omega$.

(H)

$$\begin{aligned} a_{ij\ell m} &= a_{\ell m ij} = a_{ji\ell m} \\ \frac{1}{2} \mathbf{e} \bullet \mathbf{a} \bullet \mathbf{e} &\geq \exists \eta \|\mathbf{e}\|^2 \end{aligned}$$



$(\mathbf{A}, \mathbf{B}_\alpha) : \mathbf{H}^{s,p}(\Omega, \mathbf{R}^3) \rightarrow \mathbf{H}^{s-2,p}(\Omega, \mathbf{R}^3) \times \mathbf{B}_{(\alpha)}^{s-1-1/p, p}(\partial\Omega, \mathbf{R}^3)$

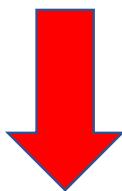
is an algebraic and topological **isomorphism**

Regularity Theorem

Regularity Theorem

(H)

$$\begin{aligned} a_{ij\ell m} &= a_{\ell m ij} = a_{j i \ell m} \\ \frac{1}{2} \mathbf{e} \bullet \mathbf{a} \bullet \mathbf{e} &\geq \exists \eta \|\mathbf{e}\|^2 \end{aligned}$$



$$1 < p < \infty, \quad s > 1/p + 1$$

$$\boxed{\mathbf{u} \in \mathbf{L}^p(\Omega, \mathbf{R}^3), \quad \mathbf{A} \mathbf{u} \in \mathbf{H}^{s-2,p}(\Omega, \mathbf{R}^3),} \\ \boxed{\mathbf{B}_\alpha \mathbf{u} \in \mathbf{B}_{(\alpha)}^{s-1-1/p, p}(\partial\Omega, \mathbf{R}^3) \Rightarrow \mathbf{u} \in \mathbf{H}^{s,p}(\Omega, \mathbf{R}^3).}$$

Construction of a Parametrix

Symbol of a Pseudo-Differential Operator

$$\begin{aligned} & \mathbf{t}(x', \xi') \\ &= \alpha(x') \mathbf{p}_1(x', \xi') \\ &+ [(1 - \alpha(x')) \mathbf{I} + \alpha(x') \mathbf{p}_0(x', \xi')] \end{aligned}$$

+ . . .

Here:

$$\mathbf{p}_1(x', \xi') \geq \exists c_0 |\xi'| \mathbf{I} \text{ on } T^*(\partial \Omega)$$

Elementary Lemma

$$\boxed{\begin{array}{l} f(x) \in C^2(\mathbf{R}), \\ f(x) \geq 0 \text{ on } \mathbf{R}, \\ \sup_{x \in \mathbf{R}} |f''(x)| \leq \exists c \end{array}}$$

\Rightarrow

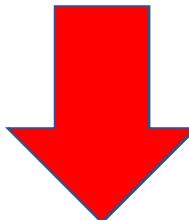
$$|f'(x)| \leq \sqrt{2c} (f(x))^{1/2} \text{ on } \mathbf{R}.$$

Criteria for Parametrices (1)

$$A = p(x, D) \in L_{1,0}^m(\Omega).$$

Assume that:

$$\begin{aligned} |D_\xi^\alpha D_x^\beta p(x, \xi)| &\leq \exists C_{K,\alpha,\beta} |p(x, \xi)| (1 + |\xi|)^{-|\alpha| + (1/2)|\beta|}, \\ |p(x, \xi)|^{-1} &\leq \exists C_K, \forall x \in K \subset \Omega, \forall |\xi| \geq C_K. \end{aligned}$$



Criteria for Parametrices (2)

$\exists B \in L_{1,1/2}^0(\Omega)$ such that

$$AB \equiv I \pmod{L^{-\infty}(\Omega)},$$

$$BA \equiv I \pmod{L^{-\infty}(\Omega)}.$$

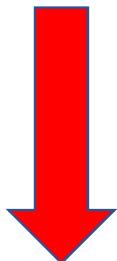
Uniqueness Theorem

Uniqueness Theorem

$$\mathbf{v} \in \mathbf{H}^{s,p}(\Omega, \mathbf{R}^3)$$

$$\mathbf{A} \mathbf{v} = \operatorname{div}(\mathbf{a} \cdot \nabla \mathbf{v}) = \mathbf{0} \quad \text{in } \Omega,$$

$$\mathbf{B}_\alpha \mathbf{v} = \alpha(x') (\mathbf{a} \cdot \nabla \mathbf{v} \cdot \mathbf{n}) + (1 - \alpha(x')) \mathbf{v} = \mathbf{0} \quad \text{on } \partial \Omega,$$



$$(A) \quad 0 \leq \alpha(X) \leq 1 \quad \text{and} \quad \alpha(X) \neq 1 \quad \text{on } \partial \Omega.$$

$$(H) \quad \left[\begin{array}{l} \mathbf{a}_{ij\ell m} = \mathbf{a}_{\ell m ij} = \mathbf{a}_{ji\ell m} \\ \frac{1}{2} \mathbf{e} \bullet \mathbf{a} \bullet \mathbf{e} \geq \exists \eta \|\mathbf{e}\|^2 \end{array} \right]$$

$$\mathbf{v} = \mathbf{0} \quad \text{in } \Omega$$

Korn's Inequality

A non-empty set : $\gamma \subset \partial \Omega$

$$\mathbf{u} \in \mathbf{H}^{1,2}(\Omega, \mathbf{R}^3)$$

$$\mathbf{u} = \mathbf{0} \text{ on } \gamma$$



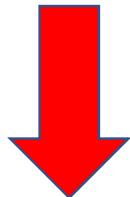
$$\int_{\Omega} \|\mathbf{e}\|^2 dx \geq \exists c(\gamma) \left(\int_{\Omega} \|\mathbf{u}\|^2 dx + \int_{\partial\Omega} \|\nabla \mathbf{u}\|^2 dx \right)$$

$$e_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) : \text{the strain tensor}$$

Existence Theorem

Existence Theorem

$$\mathbf{f} \in \mathbf{H}^{s-2,p}(\Omega, \mathbf{R}^3), \quad \varphi \in \mathbf{B}_{(\alpha)}^{s-1-1/p,p}(\partial\Omega, \mathbf{R}^3),$$



	$(A) \quad 0 \leq \alpha(X) \leq 1 \text{ and } \alpha(X) \neq 1 \text{ on } \partial\Omega.$
	$(H) \quad \boxed{\begin{array}{l} a_{ij\ell m} = a_{\ell m ij} = a_{j i \ell m} \\ \frac{1}{2} \mathbf{e} \bullet \mathbf{a} \bullet \mathbf{e} \geq \exists \eta \ \mathbf{e}\ ^2 \end{array}}$

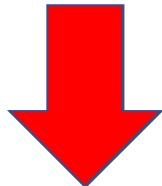
$$\exists \mathbf{v} \in \mathbf{H}^{s,p}(\Omega, \mathbf{R}^3)$$

	$\mathbf{A} \mathbf{v} = \operatorname{div}(\mathbf{a} \bullet \nabla \mathbf{v}) = \mathbf{f} \quad \text{in } \Omega,$
	$\mathbf{B}_\alpha \mathbf{v} = \alpha(x') (\mathbf{a} \bullet \nabla \mathbf{v} \bullet \mathbf{n}) + (1 - \alpha(x')) \mathbf{v} = \varphi \quad \text{on } \partial\Omega$

The Operator A

$$A : L^2(\Omega, \mathbf{R}^3) \rightarrow L^2(\Omega, \mathbf{R}^3)$$

$$\begin{aligned} \text{(a)} \quad D(A) &= \{ \mathbf{u} \in H^2(\Omega, \mathbf{R}^3) : \mathbf{B}_\alpha \mathbf{u} = \mathbf{0} \} . \\ \text{(b)} \quad A \mathbf{u} &= \operatorname{div}(\mathbf{a} \cdot \nabla \mathbf{u}), \quad \forall \mathbf{u} \in D(A). \end{aligned}$$



- A : positive definite, self - adjoint

Function Space (1)

$\mathbf{B} = \sqrt{-A}$: the **square root** of $-A$

$H_A =$ the **domain** $D(\mathbf{B})$ with the inner product $(\mathbf{u}, \mathbf{v})_A = (\mathbf{B}\mathbf{u}, \mathbf{B}\mathbf{v}).$

(A) $0 \leq \alpha(X) \leq 1$ and $\alpha(X) \not\equiv 1$ on $\partial\Omega.$

Function Space (2)

H_A = the **completion** of $D(A)$ with respect to
the inner product

$$(\mathbf{u}, \mathbf{v})_A = -(\mathbf{A}\mathbf{u}, \mathbf{v})$$

$$= \int_{\Omega} \nabla \mathbf{u} \cdot \mathbf{a} \cdot \overline{\nabla \mathbf{v}} \, dx + \int_{\{\alpha \neq 0\}} \frac{1 - \alpha}{\alpha} \mathbf{u} \cdot \overline{\mathbf{v}} \, d\sigma$$

Function Spaces (4)

$$H_A = H_0^1(\Omega, \mathbf{R}^3)$$

if $\alpha(X) \equiv 0$ on $\partial\Omega$ (**Dirichlet**)

$$H_A = H^1(\Omega, \mathbf{R}^3)$$

if $\alpha(X) \equiv 1$ on $\partial\Omega$ (**Neumann**)

Embedding Properties

$$D(A) \subset H_A \subset H^1(\Omega, \mathbf{R}^3)$$

Existence and Uniqueness Theorem

$$\mathbf{u}_0 \in D(\mathbf{A}), \quad \mathbf{u}_1 \in \mathbb{H}_{\mathbf{A}}, \quad \mathbf{f} \in C^1([0, \infty), L^2(\Omega, \mathbf{R}^3)),$$



$$\begin{aligned} & \frac{\partial^2 \mathbf{u}}{\partial t^2} - \operatorname{div} (\mathbf{a} \cdot \nabla \mathbf{u}) = \mathbf{f} \quad \text{in } \Omega \times (0, \infty), \\ & \mathbf{u} \Big|_{t=0} = \mathbf{u}_0 \quad \text{in } \Omega, \\ & \frac{\partial \mathbf{u}}{\partial t} \Big|_{t=0} = \mathbf{u}_1 \quad \text{in } \Omega, \\ & \alpha(x')(\mathbf{a} \cdot \nabla \mathbf{u} \cdot \mathbf{n}) + (1 - \alpha(x'))\mathbf{u} = \mathbf{0} \quad \text{on } \partial\Omega \times (0, \infty) \end{aligned}$$

has a **unique solution**

$$\mathbf{u} \in C^2([0, \infty), L^2(\Omega, \mathbf{R}^3)) \cap C^1([0, \infty), \mathbb{H}_{\mathbf{A}}) \cap C([0, \infty), D(\mathbf{A}))$$

Energy Estimate

$$\begin{aligned} & \left\| \mathbf{u}(t) \right\|_{H^2}^2 + \left\| \mathbf{u}'(t) \right\|_{H_A}^2 + \left\| \mathbf{u}''(t) \right\|_{L^2}^2 \\ & \leq \exists C \left(\left\| \mathbf{u}_0 \right\|_{H^2}^2 + \left\| \mathbf{u}_1 \right\|_{H_A}^2 + \left\| \mathbf{f}(0) \right\|_{L^2}^2 + \int_0^t \left\| \mathbf{f}'(s) \right\|_{L^2}^2 ds \right). \end{aligned}$$

Regularity Theorem

$$\forall \mathbf{u}_0 \in D(\mathbf{A}), \quad \forall \mathbf{u}_1 \in \mathbb{H}_{\mathbf{A}}, \quad \forall \mathbf{f} \in C^1([0, \infty), L^2(\Omega, \mathbf{R}^3))$$

\Rightarrow

$$\exists \mathbf{u} \in C^2([0, \infty), L^2(\Omega, \mathbf{R}^3)) \cap C^1([0, \infty), \mathbb{H}_{\mathbf{A}}) \cap C([0, \infty), D(\mathbf{A}))$$

such that

$$\frac{\partial^2 \mathbf{u}}{\partial t^2} - \operatorname{div}(\mathbf{a} \bullet \nabla \mathbf{u}) = \mathbf{f} \quad \text{in } \Omega \times (0, \infty),$$

$$\mathbf{u}|_{t=0} = \mathbf{u}_0 \quad \text{in } \Omega,$$

$$\left. \frac{\partial \mathbf{u}}{\partial t} \right|_{t=0} = \mathbf{u}_1 \quad \text{in } \Omega,$$

$$\alpha(x')(\mathbf{a} \bullet \nabla \mathbf{u} \bullet \mathbf{n}) + (1 - \alpha(x'))\mathbf{u} = \mathbf{0} \quad \text{on } \partial\Omega \times (0, \infty)$$

Semigroup Approach

Semigroup Approach (1)

$$\frac{\partial^2 \mathbf{u}}{\partial t^2} - \operatorname{div}(\mathbf{a} \cdot \nabla \mathbf{u}) = \mathbf{f} \quad \text{in } \Omega \times (0, \infty),$$

$$\mathbf{u} \Big|_{t=0} = \mathbf{u}_0 \quad \text{in } \Omega,$$

$$\left. \frac{\partial \mathbf{u}}{\partial t} \right|_{t=0} = \mathbf{u}_1 \quad \text{in } \Omega,$$

$$\alpha(x') (\mathbf{a} \cdot \nabla \mathbf{u} \cdot \mathbf{n}) + (1 - \alpha(x')) \mathbf{u} = \mathbf{0} \quad \text{on } \partial\Omega \times (0, \infty).$$

$$\mathbf{u}''(t) = \mathbf{A} \mathbf{u}(t) + \mathbf{f}(t),$$

$$\mathbf{u}(0) = \mathbf{u}_0, \quad \mathbf{u}'(0) = \mathbf{u}_1$$

Semigroup Approach (2)

$$\begin{aligned}\mathbf{u}''(t) &= \mathbf{A} \mathbf{u}(t) + \mathbf{f}(t), \\ \mathbf{u}(0) &= \mathbf{u}_0, \quad \mathbf{u}'(0) = \mathbf{u}_1\end{aligned}$$

$$\begin{aligned}\mathbf{U}(t) &= \begin{pmatrix} \mathbf{u}(t) \\ \mathbf{u}'(t) \end{pmatrix}, \quad \mathbf{U}(0) = \begin{pmatrix} \mathbf{u}_0 \\ \mathbf{u}_1 \end{pmatrix}, \\ \mathbf{F}(t) &= \begin{pmatrix} 0 \\ \mathbf{f}(t) \end{pmatrix}, \quad \mathfrak{A} = \begin{pmatrix} 0 & \mathbf{I} \\ \mathbf{A} & 0 \end{pmatrix}\end{aligned}$$

Semigroup Approach (3)

$$\begin{aligned}\mathbf{U}'(t) &= \mathfrak{A}\mathbf{U}(t) + \mathbf{F}(t), \\ \mathbf{U}(0) &= \mathbf{U}_0.\end{aligned}$$

$$\begin{aligned}X &= \mathbb{H}_{\mathbf{A}} \times L^2(\Omega, \mathbf{R}^3), \\ D(\mathfrak{A}) &= D(\mathbf{A}) \times \mathbb{H}_{\mathbf{A}}.\end{aligned}$$

Semigroup Approach (4)

$$e^{-t\mathfrak{A}} = \cos(t\mathbf{B}) \begin{pmatrix} \mathbf{I} & 0 \\ 0 & \mathbf{I} \end{pmatrix} + \mathbf{B}^{-1} \sin(t\mathbf{B}) \begin{pmatrix} 0 & \mathbf{I} \\ \mathbf{A} & 0 \end{pmatrix}$$

$$\mathbf{B} = \sqrt{-\mathbf{A}}, \quad \mathbf{B}^2 = -\mathbf{A}.$$

Representation Formula

$$\begin{aligned}\mathbf{u}(t) &= \cos(t\mathbf{B})\mathbf{u}_0 + \mathbf{B}^{-1} \sin(t\mathbf{B})\mathbf{u}_1 \\ &\quad + \int_0^t \mathbf{B}^{-1} \sin((t-s)\mathbf{B})\mathbf{f}(s)ds.\end{aligned}$$

$$\mathbf{B} = \sqrt{-\mathbf{A}}, \quad \mathbf{B}^2 = -\mathbf{A}.$$

Open Problems

Open Problems

1. The first problem is to generalize Main Theorem to the case where the domain has **corner singularities**.
2. The second problem is to study the case where the function $\alpha(X)$ is the **characteristic function** of a subset of the boundary.

References (4)

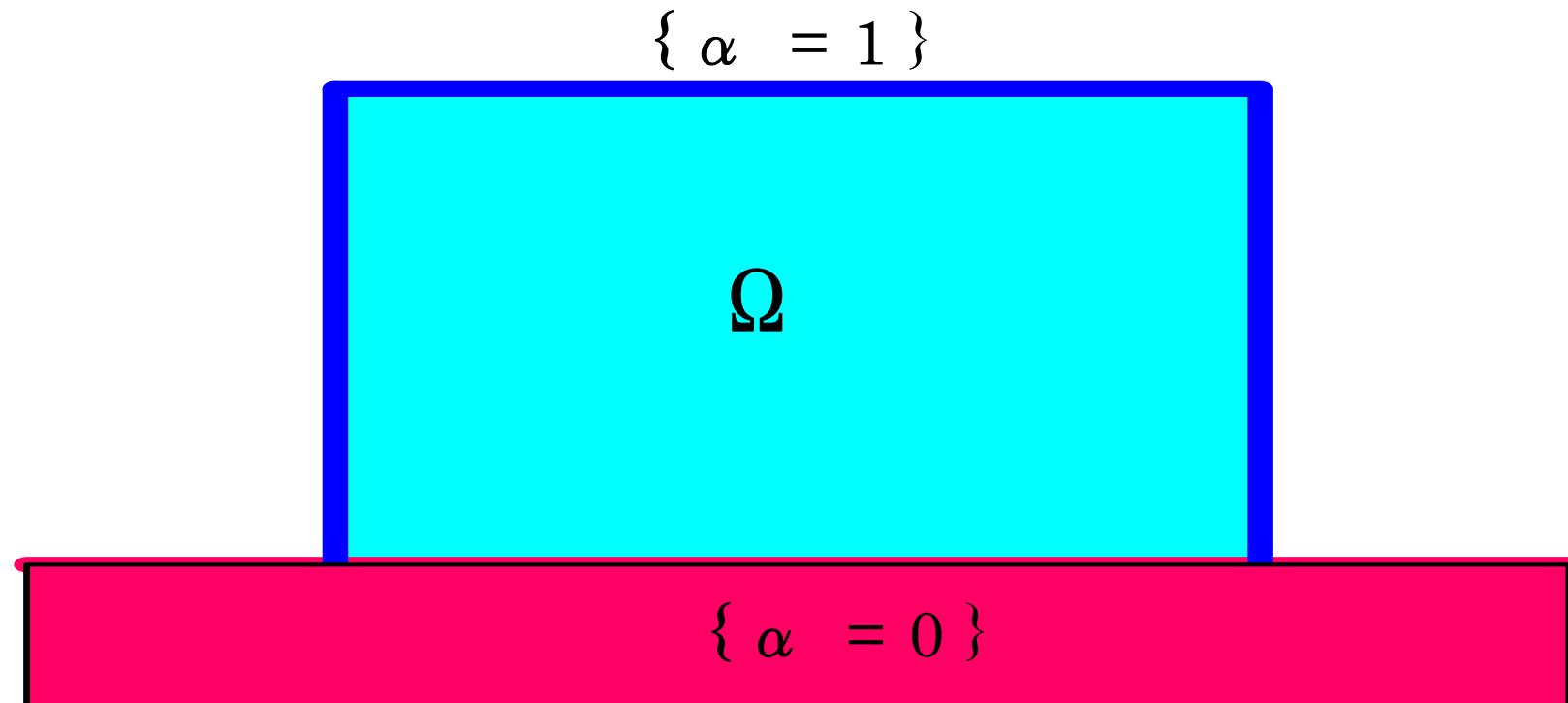
- **Ito**: On a mixed problem of linear elastodynamics with time-dependent discontinuous boundary condition, *Osaka J. Math.* 27 (1990), 667—707.

Nonlinear Boundary Value Problem (Lagrange)

This talk is devoted to an L^p approach to the following equilibrium equations for the unknown deformation $\phi(X)$

$$\begin{aligned} \text{Div } \widehat{\mathbf{P}}(X, D\phi(X)) + \mathbf{B}(X) &= \mathbf{0} \quad \text{in } \Omega, \\ \alpha(X)\widehat{\mathbf{P}}(X, D\phi(X)) \cdot \mathbf{N}(X) + (1 - \alpha(X))\phi(X) \\ &= \tau(X) \quad \text{on } \partial\Omega. \end{aligned}$$

Domain with Corner Singularities

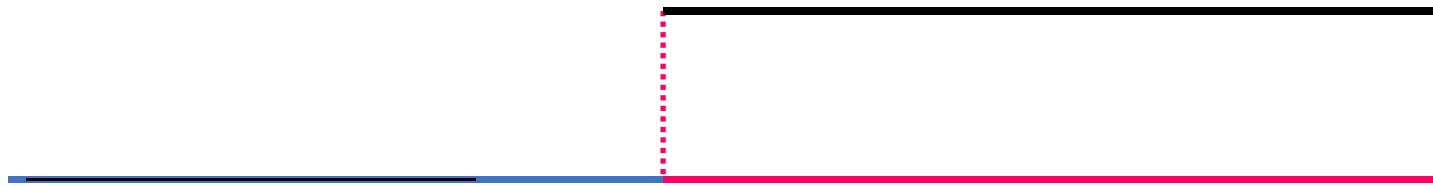


Boundary Condition (1)

$$\alpha(x) = \begin{cases} 1 & \text{on } \partial_N \Omega, \\ 0 & \text{on } \partial_D \Omega \end{cases}$$

Boundary Condition (2)

$$\alpha(x')(\mathbf{a} \cdot \nabla \mathbf{u} \cdot \mathbf{n}) + (1 - \alpha(x')) \mathbf{u} = \mathbf{0} \quad \text{on } \partial\Omega \times (0, \infty)$$



$$\partial_D \Omega = \{\alpha = 0\}$$

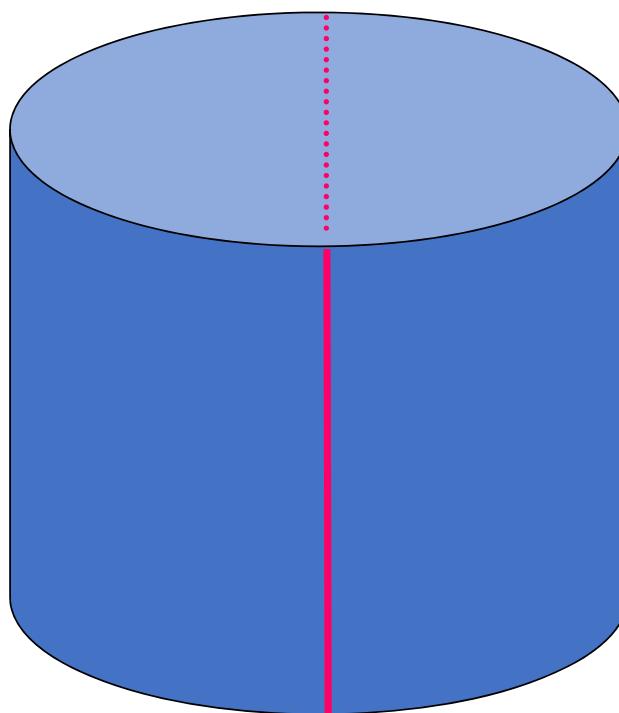
$$\partial_N \Omega = \{\alpha = 1\}$$

Geometry of the Second Problem

$$\Omega \times (0, T)$$

$$\partial_D \Omega \times (0, T)$$

$$\partial_N \Omega \times (0, T)$$



Fundamental Function Space

$$\mathbf{V}_D(0) = \left\{ \mathbf{u} \in H^1(\Omega, \mathbf{R}^3) : \mathbf{u} = \mathbf{0} \text{ on } \partial_D \Omega \right\}$$

Definition of a Weak Solution (1)

$$\mathbf{u}_0 \in \mathbf{V}_D(0), \quad \mathbf{u}_1 \in L^2(\Omega, \mathbf{R}^3), \quad \mathbf{f} \in L^2(\Omega \times (0, T), \mathbf{R}^3)$$

A function

$$\mathbf{u} \in H^1(\Omega \times (0, T), \mathbf{R}^3)$$

is called a **weak solution** if
it satisfies the conditions

$$\mathbf{u}(0, \cdot) = \mathbf{u}_0,$$

$$\mathbf{u} = \mathbf{0} \text{ on } \partial_D \Omega \times (0, T)$$

Definition of a Weak Solution(2)

$$\begin{aligned} & - \int_0^T \left(\frac{\partial \mathbf{u}}{\partial t}, \frac{\partial \eta}{\partial t} \right)_{L^2} dt + \int_0^T (\mathbf{a} \cdot \nabla \mathbf{u}, \nabla \eta)_{L^2} dt \\ &= (\mathbf{u}_1, \eta(0, \cdot))_{L^2} + \int_0^T (\mathbf{f}, \eta)_{L^2} dt \end{aligned}$$

$\forall \eta \in H^1(\Omega \times (0, T), \mathbf{R}^3), \eta = 0$ on $\partial_D \Omega \times (0, T)$, $\eta(T, \cdot) = 0$.

Existence and Uniqueness Theorem

$$\begin{aligned} \frac{\partial^2 \mathbf{u}}{\partial t^2} - \operatorname{div} (\mathbf{a} \cdot \nabla \mathbf{u}) &= \mathbf{f} \quad \text{in } \Omega \times (0, T), \\ \mathbf{u} \Big|_{t=0} &= \mathbf{u}_0 \quad \text{in } \Omega, \\ \frac{\partial \mathbf{u}}{\partial t} \Big|_{t=0} &= \mathbf{u}_1 \quad \text{in } \Omega, \\ \mathbf{u} &= \mathbf{0} \quad \text{on } \partial_D \Omega \times (0, T), \\ \mathbf{a} \cdot \nabla \mathbf{u} \cdot \mathbf{n} &= \mathbf{0} \quad \text{on } \partial_N \Omega \times (0, T) \end{aligned}$$

has a **unique weak solution**

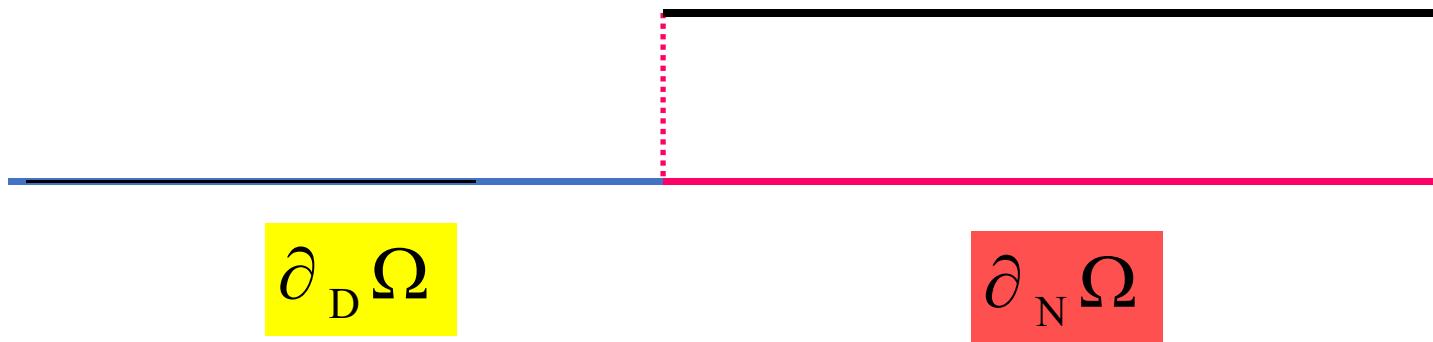
$$\mathbf{u} \in H^1(\Omega \times (0, T), \mathbf{R}^3)$$

Energy Estimate for Weak Solutions

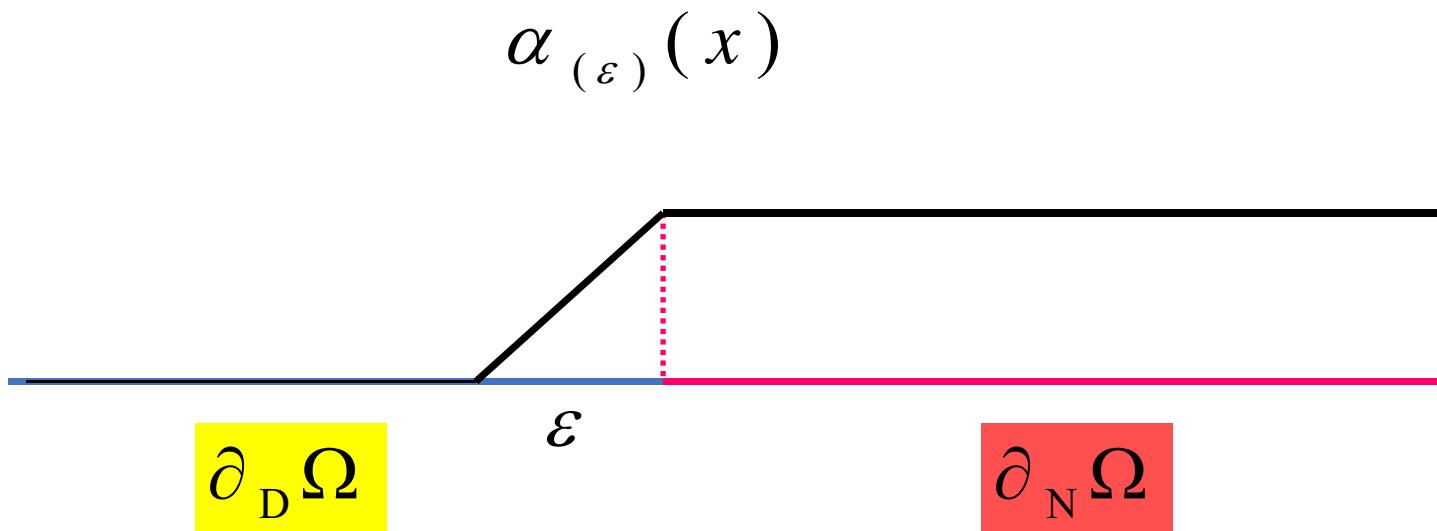
$$\begin{aligned} & \left\| \mathbf{u}(t) \right\|_{H^1}^2 + \left\| \mathbf{u}'(t) \right\|_{L^2}^2 \\ & \leq \exists C(T) \left(\left\| \mathbf{u}_0 \right\|_{H^1}^2 + \left\| \mathbf{u}_1 \right\|_{L^2}^2 + \int_0^t \left\| \mathbf{f}(s) \right\|_{L^2}^2 ds \right). \end{aligned}$$

The Original Boundary Condition

$\alpha (x)$



Approximations to the Boundary Condition



Regularized Mixed Problems

We study the following **regularized** mixed problem:

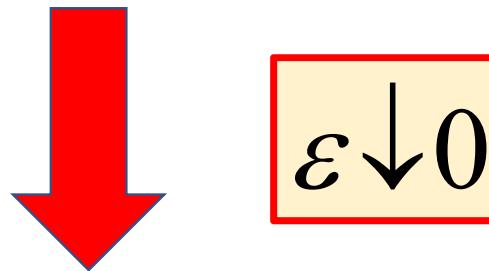
$$\begin{cases} \frac{\partial^2 \mathbf{u}_\varepsilon}{\partial t^2} - \operatorname{div}(\mathbf{a} \cdot \nabla \mathbf{u}_\varepsilon) = \mathbf{f}^{(\varepsilon)} & \text{in } \Omega \times (0, \infty), \\ \mathbf{u}_\varepsilon|_{t=0} = \mathbf{u}_0^{(\varepsilon)} & \text{in } \Omega, \\ \frac{\partial \mathbf{u}_\varepsilon}{\partial t}|_{t=0} = \mathbf{u}_1^{(\varepsilon)} & \text{in } \Omega, \\ \alpha_{(\varepsilon)} (\mathbf{a} \cdot \nabla \mathbf{u}_\varepsilon \cdot \mathbf{n}) + (1 - \alpha_{(\varepsilon)}) \mathbf{u}_\varepsilon = \mathbf{0} & \text{on } \partial\Omega \times (0, \infty). \end{cases}$$

Energy Estimate for Regularized Solutions

$$\forall \varepsilon > 0$$

$$\begin{aligned} & \| \mathbf{u}_\varepsilon(t) \|_{H^1}^2 + \| \mathbf{u}'_\varepsilon(t) \|_{L^2}^2 \\ & \leq \exists C(T) \left(\| \mathbf{u}_0^{(\varepsilon)} \|_{H^1}^2 + \| \mathbf{u}_1^{(\varepsilon)} \|_{L^2}^2 + \int_0^t \| \mathbf{f}^{(\varepsilon)}(s) \|_{L^2}^2 ds \right). \end{aligned}$$

Energy Estimate for Weak Solutions



$$\begin{aligned} & \| \mathbf{u}(t) \|_{H^1}^2 + \| \mathbf{u}'(t) \|_{L^2}^2 \\ & \leq \exists C(T) \left(\| \mathbf{u}_0 \|_{H^1}^2 + \| \mathbf{u}_1 \|_{L^2}^2 + \int_0^t \| \mathbf{f}(s) \|_{L^2}^2 ds \right). \end{aligned}$$

END